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# Laser Action in a Gas Condensation in the Vicinity of a Hot Star<sup>1</sup>

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In the region near 1  $\mu\text{m}$ , we have found laser action in a quantum transition between highly excited states of Fe II, with its higher levels being optically pumped by the intense H Ly $\alpha$  radiation (1215 Å) formed in the ionized HII region of a gas condensation (blob B) in the close vicinity of the central star in  $\eta$  Carinae. © 2002 MAIK “Nauka/Interperiodica”.

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Following the discovery of space microwave masers operating in the OH radicals [1] and H<sub>2</sub>O molecules [2], such masers were found to be operative in more than 100 molecular species [3], as well as in highly excited hydrogen atoms [4] in the submillimeter wavelength region. In the IR region of the spectrum, the effect of stimulated emission of radiation in the CO<sub>2</sub> molecule was discovered in the Martian and Venus atmospheres [5, 6]. In this letter, we report on the laser action in the optical region of the spectrum—a possibility that was discussed as far back as 1972 [7].

The probability of finding a laser effect in the optical region differs considerably from that in the microwave region, as a result of the enormous difference ( $10^{15}$  to  $10^{18}$  times) between the spontaneous emission rates and inverse population production mechanisms in the two wavelength regions. This follows from simple qualitative considerations of the steady-state saturation regime of the isotropic maser/laser action.

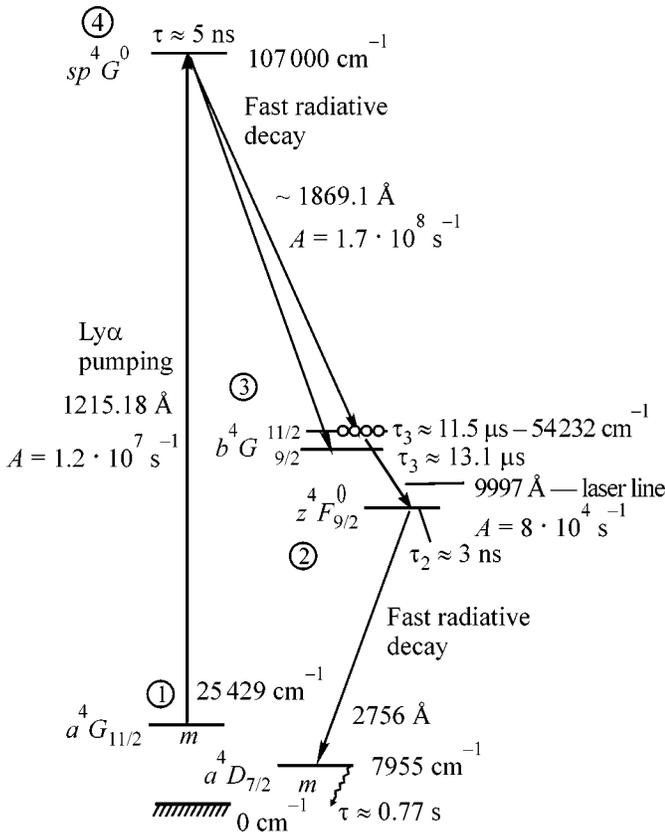
The stimulated emission of radiation in a space maser occurs as a result of pumping the upper maser level, which is not associated with spontaneous emission in microwave radiative transitions. It may therefore have a decay rate much in excess of the spontaneous emission rate, which lies in the region  $A_{mn} \cong 10^{-9}$ – $10^{-7}$  s<sup>-1</sup>. Therefore, the intensity of stimulated radiation can be many orders of magnitude higher than the intensity of spontaneous radiation, and it is only limited by the pumping rate. This is exactly the reason why the brightness temperature of maser microwave lines reaches as high a value as  $10^{10}$ – $10^{15}$  K. The intensity of maser lines is not borrowed from other microwave spectral lines, which are very weak, but from other pumping sources.

In the optical region of the spectrum, the rate of allowed spontaneous transitions is high ( $A_{mn} \cong 10^8$ – $10^9$  s<sup>-1</sup>), and it is precisely spontaneous transitions that provide the pumping of the upper level in an optical space laser at a sufficiently high rate to exceed the rates of collisional pumping mechanisms. This is especially true for the case considered here, where the space laser is indirectly pumped by H Ly $\alpha$  in the vicinity of  $\eta$  Car, one of the most luminous stars of our Galaxy. Therefore, the intensity of the stimulated radiation in the optical region generated by the occurrence of an inversion population and significant amplification cannot exceed in the steady-state regime to any substantial extent the intensity of the pumping spectral lines formed by spontaneous emission in allowed radiative transitions of atoms or ions. This fact presents difficulties in detecting laser action by a large enhancement of the radiation intensity, but it should manifest itself in comparatively moderate changes of the branching ratios of spectral lines, having a common source of pumping.

One exception is the case of quantum transitions having a relatively low spontaneous emission probability ( $A_{mn} \cong 1$ – $10^5$  s<sup>-1</sup>), and, consequently, the spontaneous radiation lines are weak. Once an inverted level population has developed in such a transition with a significant amplification in a properly sized cloud, a stimulated emission channel opens up. The stimulated transition rate cannot be much higher than the spontaneous emission rate, which is limited by the pumping rate of the upper laser level. Thus, the intensity of the laser line should become comparable with the intensity of the lines which are due to the allowed spontaneous emission of radiation and resulting from the optical pumping (direct or indirect) of the upper level. This takes place only in the inverted population volume with size  $L \gg 1/\alpha$ , where  $\alpha$  is the amplification coefficient per unit length. In such a case, a spectral line, expected to be weak, must appear as a spectral line of normal

<sup>1</sup> This work was submitted by the authors in English.

*Photoselective pumping  
of Fe II by HLy $\alpha$  Radiation*



**Fig. 1.** Schematic diagram of the relevant energy levels and radiative transitions of Fe II illustrating the formation of an inverse population of levels 3 and 2 and amplification in the transition 3  $\rightarrow$  2 owing to the photoselective pumping of level 4 by the intense Ly $\alpha$  radiation.

intensity, typical of an allowed transition. This is exactly what we have found to be the case with several spectral lines of Fe II pumped indirectly by the intense H Ly $\alpha$  radiation (1215 Å) in a dense-gas condensation (blob) in the vicinity of  $\eta$  Car.

Using the Space Telescope Imaging Spectrograph (STIS) onboard the Hubble Space Telescope, the authors of [8] recorded the spectra of blob B in the vicinity of  $\eta$  Car with a high angular and spectral resolution. The spectrum contains an intense 9997 Å line, as well as other Fe II lines at 9391, 9617, and 9913 Å, which can be excited by HLy $\alpha$  radiation in a cascade fashion [9], shown schematically in Fig. 1. All these spectral lines have long-lived (from a few microseconds to a few milliseconds) upper states and a short-lived (3.7 ns) lower state; i.e., an inverted population always exists in these transitions. The gain by these inverted-population transitions is governed by the rate of photoexcitation of the upper levels by the HLy $\alpha$  radiation.

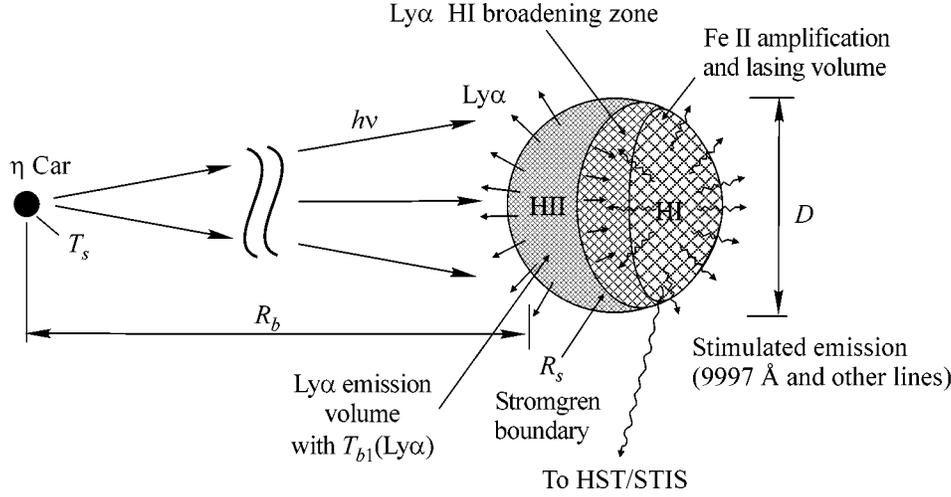
Dense-gas condensations (Weigelt blobs [10]) in the neighborhood of the luminous blue variable (LBV) and massive star  $\eta$  Carinae are capable of absorbing all the photospheric radiation in the wide Lyman continuum region ( $\lambda < \lambda_c = 912$  Å) incident upon them from the central star. This occurs because of the high Lyman-continuum optical density  $\tau(\lambda_c)$  of the neutral hydrogen component in the ionized HII/III region of the blob. Since the gas condensations are very close to the central star [11, 12], the density of the radiative energy deposited in them proves to be so high that the spectral brightness temperature  $T_{br}(\text{Ly}\alpha)$  of the more narrow Ly $\alpha$  line at 1215 Å inside the blobs can be very high. The magnitude of  $T_{br}(\text{Ly}\alpha)$  can become comparable with that of the brightness temperature of the photosphere of  $\eta$  Car,  $T_s = (20\text{--}30) \times 10^3$  K. This is a unique situation for the photoselective excitation of those atoms and ions whose absorption lines from the ground state or low-lying metastable states coincide with Ly $\alpha$ .

This is exactly the situation that occurs for the Fe II ions formed in the HII region of the blob by complete photoionization of Fe I by the radiation from  $\eta$  Car in the spectral region  $I_{\text{Fe}} = 7.6 \text{ eV} < h\nu < I_{\text{H}} = 13.6 \text{ eV}$ . The concentration of iron in typical nebulae is about 0.01% of the hydrogen concentration. When Fe I undergoes photoionization, a substantial portion of the Fe II ions formed occupy low-lying long-lived metastable states with energies  $E < I_{\text{H}} - I_{\text{Fe}} \approx 6 \text{ eV}$  as a result of the decay from numerous autoionization states. There are about 90 such metastable and pseudometastable states in Fe II. Some of them have absorption lines coinciding in wavelength with the extremely bright spectral line Ly $\alpha$ . The large number of low-lying long-lived states and the high density of spectral lines ( $\approx 15$  lines/Å in the Ly $\alpha$  region) result in several such coincidences in Fe II. Several of them are similar to those shown schematically in Fig. 1, which lead to photoselective excitation of energy levels in Fe II, and the subsequent cascading decay gives rise to population inversion.

The Ly $\alpha$  line at  $\lambda = 1215.671$  Å almost coincides with the absorption line of Fe II in the transition  $a^4G_{11/2}$  (level 1 in Fig. 1)  $\rightarrow$   $sp^4G^0$  (level 4), the frequency difference (detuning) between them being  $\Delta\nu = +30 \text{ cm}^{-1}$ . This difference is compensated by the broadening of the line profile when the Ly $\alpha$  radiation is passed through the HII region with an optical density of  $\tau \approx 3 \times 10^5$ . Thus, the transfer broadening of the Ly $\alpha$  radiation makes the photoselective excitation of state 4 in Fig. 1 quite possible.

The photoselective excitation rate of state 4 is defined by the expression

$$W_{\text{exc}}^{14} = A_{41} \left[ \exp\left(\frac{h\nu_{14}}{T_{bl}}\right) - 1 \right]^{-1}, \quad (1)$$



**Fig. 2.** Physical model of the photoionization-resonance photoexcitation pumping of Fe II in the gas condensation by the radiation emitted by  $\eta$ Car with a brightness temperature of  $T_s$ .

where the indices 1 and 4 correspond to the level designations in Fig. 1,  $A_{41} = 1.2 \times 10^7 \text{ s}^{-1}$  is the Einstein coefficient for the spontaneous decay of state 4 to state 1, and  $T_{bl}$  is the brightness temperature of the Ly $\alpha$  radiation inside blob B. For  $kT_{bl} \approx 1.0\text{--}1.5 \text{ eV}$  ( $T_{bl} \approx (12\text{--}18) \times 10^3 \text{ K}$ ), the photoselective excitation rate  $W_{\text{exc}}^{14} \approx 10^3\text{--}10^4 \text{ s}^{-1}$ .

The main decay channel (branching fraction = 0.97) of state 4 is the radiative transitions to two fine-structure levels of  $b^4G$ , one of which ( $b^4G_{11/2}$  is level 3) has a channel of spontaneous radiative decay with a probability of  $A_{32} = 8 \times 10^4 \text{ s}^{-1}$  to level 2 ( $z^4F_{9/2}^0$ ) with a much shorter lifetime ( $\approx 3 \text{ ns}$ ). It is exactly in this  $3 \rightarrow 2$  transition that a steady-state population inversion is reached, its density being

$$\Delta N = N_3 - N_2 = N_3 = (W_{\text{exc}}^{14} \tau_3) N_1, \quad (2)$$

where  $N_2 \ll N_3$  because of the much faster decay of the level 2,  $\tau_3 = 11.5 \mu\text{s}$ , and  $N_1$  is the population of the initial metastable state with a lifetime of  $\tau_1 = 0.77 \text{ s} \gg \tau_3$ .

The linear amplification coefficient for the  $3 \rightarrow 2$  transition at  $9997 \text{ \AA}$  is defined by the standard expression

$$\alpha_{32} = \sigma_{32} \Delta N, \quad (3)$$

where  $\Delta N$  is the inverted-population density defined by Eq. (2). The stimulated emission cross section  $\sigma_{32}$  is given by

$$\sigma_{32} = \frac{\lambda_{32}^2 A_{32}}{2\pi 2\pi \Delta\nu_D}, \quad (4)$$

where  $\Delta\nu_D$  is the Doppler width of the  $3 \rightarrow 2$  transition in Fe II. At a temperature of  $T \approx 100\text{--}1000 \text{ K}$  in the relatively cold HI region,  $\Delta\nu_D \approx (300\text{--}1000) \text{ MHz}$ , i.e.,

$\sigma_{32} = (0.6\text{--}2) \times 10^{-13} \text{ cm}^2$ . Thus, the amplification coefficient may be estimated by the expression

$$\alpha_{32} = \sigma_{32} (W_{\text{exc}}^{14} \tau_3) f N_0, \quad (5)$$

where  $f$  is the fraction of the Fe II ions in their initial state relative to all Fe II ions (all iron atoms in the HI region of the blob are ionized).

The fraction  $f$  is governed by the excitation rate of the level 1 and its lifetime  $\tau_1$ . The rates of collisional excitation mechanisms for the level 1 (i.e., the recombination of the  $\text{Fe}^{2+}$  ions and electron collisional excitation) are negligible in comparison with the decay rate  $1/\tau_1$  of the state 1 because of the low electron concentration in the HI region ( $n_e \approx 10^4\text{--}10^5$  due to the photoionization of iron, as well as other elements with an ionization potential of  $I < 13.6 \text{ eV}$ ). Most important for the laser action observed are the following radiative channels: (1) the radiative decay of the high-lying states of Fe II excited by the Ly $\alpha$  or other intense radiation lines and (2) the decay of autoionization states of Fe I excited by the radiation emitted by  $\eta$  Car in the spectral region  $I_{\text{Fe}} + E_1 = 10.6 \text{ eV} < h\nu < 13.6 \text{ eV}$ , where  $E_1$  is the energy of state 1. These excitation channels can provide an excitation rate  $> 1/\tau_1 \approx 1 \text{ s}^{-1}$  and, hence, sustain the relative population of state 1 at a level, for example, of  $f \approx 10^{-2}$ . This would correspond to an approximately equal distribution of the Fe II ions among their 90 metastable and pseudometastable states, including state 1. Leaving the calculation of the magnitude of the fraction  $f$  for future consideration, we will restrict ourselves here to the qualitative estimate of  $f \approx 10^{-2}$ .

In this approximation, the amplification coefficient for the transition  $3 \rightarrow 2$  may be estimated, according to Eq. (5), at  $\alpha_{32} \approx (3 \times 10^{-18} \text{ to } 10^{-16}) N_0 [\text{cm}^{-1}]$ , where  $N_0 \approx 10^{-4} N_{\text{H}}$  is the density of the Fe atoms,  $N_{\text{H}}$  being the

density of the hydrogen atoms in the blob. According to the data presented in [11] and the results of a calculation of the critical density of hydrogen atoms [13], the density  $N_{\text{H}}$  of the hydrogen atoms in the blob is obviously higher than  $10^8 \text{ cm}^{-3}$ . Thus, for  $\alpha_{32} \geq 3 \times 10^{-14} - 10^{-12} \text{ cm}^{-1}$  and a blob diameter  $D \approx 10^{15} \text{ cm}^{-3}$  [11], which can be regarded as the size  $\alpha L$  of the amplifying region (Fig. 2),  $\alpha L \approx (30-1000)$  at a blob temperature of  $T_{\text{bl}} = (12-18) \times 10^3 \text{ K}$ .

These values correspond to fairly high linear amplification coefficients  $K = \exp(\alpha L)$ . However, at an amplification coefficient of  $K \approx A_{41}/A_{32} \approx 10^3$  for an isotropic radiation, the intensity of the weak line  $\lambda_{32}$  (in photons/cm<sup>2</sup> s) approaches that of the strong line  $\lambda_{43}$ ; i.e., the rate of stimulated transition approaches the pumping rate for level 3. Thereafter, the amplification regime becomes saturated, and the intensity of the laser line  $\lambda_{32}$  becomes equal to that of the pumping line  $\lambda_{43}$ . Under such saturated amplification conditions, the intensities of both these lines grow in proportion to the propagation length  $L$ .

We believe that laser amplification and stimulated emission of radiation is a fairly common and widespread phenomenon, at least, for gas condensations in the vicinity of bright stars. This is due to the occurrence of two types of processes (fast radiative and slow collisional) in a very rarefied gas of the condensations, whereby the populations of electronic levels in atoms (ions) can relax. These relaxation processes occur on highly different time scales, radiative relaxation operating on a time scale of  $10^{-9}-10^{-3}$  (sometimes even within  $10^{-3}$ ), and collisional relaxation, on a time scale of over 100 s (at gas densities  $<10^9 \text{ cm}^{-3}$ ). In the case of photo-selective excitation of some high-lying levels of an atom or ion with a complex energy-level structure, radiative relaxation takes place as a consequence of downward transitions with spontaneous emission of radiation, in the course of which there inevitably develops an inverse population of some pair (pairs) of levels. If the size of a gas cloud is large enough, large amplification

of the inverted-population transition automatically switches on the radiative relaxation channel, which leads to faster stimulated quantum transitions until collisional relaxation becomes important. Thus, the laser action is an intrinsic characteristic of the radiative cooling of gas clouds near bright stars by stimulated emission for inverted transitions along with spontaneous emission for normal noninverted transitions.

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