Natural Radiative Lifetimes In the Interacting 5snd 1,3d2 Sequences In Sr

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Published in:
Physical Review A (Atomic, Molecular and Optical Physics)

DOI:
10.1103/PhysRevA.27.947

Published: 1983-01-01

Citation for published version (APA):

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Natural radiative lifetimes in the interacting 5snd \(^{1,3}D_2\) sequences in Sr


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(Received 11 June 1982)

Natural radiative lifetimes have been determined in the 5snd \(^{1,3}D_2\) sequences of strontium in the perturbed region around \(n=15\). Stepwise laser excitations and pulse modulation of a cw dye-laser beam were employed. A plot of the measured lifetime values versus the effective quantum number reveals complex configuration interaction.

I. INTRODUCTION

The alkaline-earth atoms have two electrons outside closed shells and the presence of low-lying, doubly excited states results in strong perturbations of sequences with only one excited electron. These configuration-interaction effects can be studied not only in the basic energy-level structure but are also reflected in, e.g., radiative lifetimes, hyperfine structures, and isotope shifts. For a theoretical description of these effects, the multichannel quantum-defect theory (MQDT) is usually applied.

In strontium, the energy levels of the 5snd \(^{1,3}D_2\) sequences together with perturbing states have been measured by several authors.\(^1\) \(^2\) \(^3\) In Ref. 2 the data were analyzed with the use of MQDT. The \(g_J\) factors, calculated from the obtained wave functions, were found to agree very well with experimental values determined in a Zeeman-effect investigation in the perturbed region around \(n=15\).\(^4\) In the same region, a strong variation in the hyperfine interaction has also been found.\(^5\)

In the present work we have measured radiative lifetimes in the 5snd \(^{1}D_2\), \(n=13-22\), and 5snd \(^{3}D_2\), \(n=14-17\) sequences of strontium using the PUMOLS technique (pulse modulated laser spectroscopy).\(^6\) Stepwise excitations of a collimated atomic beam were employed. By using a single-mode laser for the second step a selective excitation of the isotope \(^{88}\)Sr was performed. Lifetime values for lower-energy members of the \(1D_2\) sequence and some low-lying \(^1S_0\) levels have been measured by Gornik.\(^7\)

For radiative lifetimes, properties of the wave functions other than those for energy levels and \(g_J\) factors are important, and it is interesting to investigate how well different manifestations of perturbation can be reproduced. In recent measurements on barium regarding lifetimes\(^8\) and \(g_J\) factors,\(^9\) published MQDT wave functions\(^10\) were found to describe the radiative properties well but resulted in inadequate predictions of the \(g_J\) factors. This has resulted in a refinement in the MQDT treatment,\(^9\)\(^11\) leading to a much improved \(g_J\) description.

II. EXPERIMENTAL ARRANGEMENT

The experimental setup used in the present measurements was similar to that described in Ref. 8. A blue multimode cw dye laser, operating with the dye Coumarin 1, induced the 4607-A transition from the ground state 5s\(^2\) \(^1\)S\(_0\) to the intermediate 5s 5p \(^1\)P\(_1\) level, which has a lifetime of only 4.8 ns.\(^12\) A single-mode ring dye laser, operating with the dye Stilbene 3, was used for the second step from the \(^1\)P\(_1\) to the \(^{1,3}D_2\) states. Both dye lasers were pumped by ion lasers with uv lines. The output beam of the single-mode dye laser was pulse-modulated by an acousto-optic modulator for delayed-coincidence measurements. The atomic beam of low collimation came from a resistively heated oven containing Sr metal. The two laser beams were made to overlap and crossed the atomic beam at about a right angle. The decay of the investigated level was detected with a cooled photomultiplier tube equipped with suitable filters. A photon, detected by the photomultiplier, was used as a "stop" signal, whereas a signal derived from the prompt pulse was used as a "start" signal for a time-to-amplitude converter. The pulses produced by this unit were fed to a multichannel analyzer (Tracer Northern TN 1710).

![Experimental decay curves for the Sr states 5s\(^1\)D\(_2\) and 5s\(^3\)D\(_2\). Background counts are subtracted.](image)
second step are fairly close, complete suppression of the intense blue resonance fluorescence was sometimes hard to achieve when we detected at the laser wavelength. On some occasions uv detection of the decay to the 5s 5p 3P levels was used resulting in a more efficient suppression. Experimental decay curves for the 5s 16d 1D2 and 5s 17d 1D2 states are shown in Fig. 1. The multichannel analyzer was interfaced to a minicomputer and every decay curve was stored as a data file on a floppy disc. This provided us with the opportunity of calculating the radiative lifetime of one decay curve while another one was being stored in the multichannel analyzer. For each state, measurements were performed several times and particular attention was paid to ensure that there was no influence due to pileup, multiple scattering, collisions, Zeeman quantum beats, and flight-out-of-view effects.

The experimentally determined natural radiative lifetimes for the Sr sequences 5s 1D 1D2 (n = 13–22) and 5s 1D 2D2 (n = 14–17) can be seen in Table I and Fig. 2. The states 5s 15s 1S0, 5s 16s 1S0, 5s 7f 1F3, and 5s 10p 1P1 could also be reached in the employed laser wavelength region and the corresponding lifetime values are included in Table I. The last two of these states were excited from the metastable 5s 4d 1D2 level, populated in the cascade decay of the 5s 5p 1P1 level. The given lifetime values are referred to room temperature and are not corrected for expected small effects of black-body radiation (see, e.g., Refs. 13 and 8). The lifetime of the doubly excited perturber state 4d 2P2 was also determined. In order to measure this particularly short lifetime we first recorded the shape of the laser pulse from the acoustooptic modulator by detecting reflected light inside the vacuum system. With uv detection, we recorded the time behavior of the fluorescent light from the state. With a computer program we could generate, from the pulse shape, decay curves for different lifetimes and make the best fit to our experimental curve.

![FIG. 2. Measured lifetimes in the Sr 1D2 sequences plotted vs the effective principal quantum number.](image)

**III. MEASUREMENTS AND RESULTS**

In the measurements the first-step laser was tuned to the 4607-Å transition from the 5s 2S0 to the 5s 5p 1P1 state by observing the intense blue fluorescence light. The wavelengths of the second-step laser were measured with a digital wavelength meter and were found to be in good agreement with Ref. 3. For each individual lifetime value a suitable pulse length and repetition rate was set on the acoustooptic modulator in order to get an optimum duty cycle. Since the wavelengths of the first and the

**TABLE I. Experimentally determined natural radiative lifetimes for Sr states (300 K).**

<table>
<thead>
<tr>
<th>State</th>
<th>Measured τ value (ns)</th>
</tr>
</thead>
<tbody>
<tr>
<td>5s 13d 1D2</td>
<td>410(20)</td>
</tr>
<tr>
<td>5s 14d 1D2</td>
<td>408(12)</td>
</tr>
<tr>
<td>5s 15d 1D2</td>
<td>340(10)</td>
</tr>
<tr>
<td>5s 16d 1D2</td>
<td>365(15)</td>
</tr>
<tr>
<td>5s 17d 1D2</td>
<td>517(16)</td>
</tr>
<tr>
<td>5s 18d 1D2</td>
<td>640(16)</td>
</tr>
<tr>
<td>5s 19d 1D2</td>
<td>738(37)</td>
</tr>
<tr>
<td>5s 20d 1D2</td>
<td>866(40)</td>
</tr>
<tr>
<td>5s 21d 1D2</td>
<td>1023(51)</td>
</tr>
<tr>
<td>5s 22d 1D2</td>
<td>1190(60)</td>
</tr>
<tr>
<td>5s 14d 1D2</td>
<td>247(8)</td>
</tr>
<tr>
<td>5s 15d 1D2</td>
<td>282(11)</td>
</tr>
<tr>
<td>5s 16d 1D2</td>
<td>286(9)</td>
</tr>
<tr>
<td>5s 17d 1D2</td>
<td>319(15)</td>
</tr>
<tr>
<td>4d 2P2</td>
<td>7(2.8)</td>
</tr>
<tr>
<td>5s 15s 1S0</td>
<td>1145(57)</td>
</tr>
<tr>
<td>5s 16s 1S0</td>
<td>1424(72)</td>
</tr>
<tr>
<td>5s 7f 1F3</td>
<td>126(5)</td>
</tr>
<tr>
<td>5s 10p 1P1</td>
<td>92(5)</td>
</tr>
</tbody>
</table>

![FIG. 3. Diagram including previously determined (Ref. 7) and currently measured lifetimes in the Sr 1D2 and 1S0 sequences.](image)
IV. DISCUSSION

In Fig. 2 the lifetimes for the $^1D_2$ sequences are plotted versus the effective principal quantum number. As can be seen, there is a decrease in the lifetime values around $n = 15$ for the $^1D_2$ sequence. In this energy range the $5snd$ $^1D_2$ levels are perturbed by the $4d6s$ $^1D_2$ levels; moreover the perturbers induce a large singlet-triplet mixing. In the figure a line with the slope of three is also included. If the effect of perturbations is neglected, the lifetimes in a Rydberg sequence are expected to scale with a power, close to three, of the effective principal quantum number. However, to be in an overall accordance with this rule, the determined values for the $13^1D_2$ and $14^1D_2$ states are too long. In Fig. 3 our $^1D$- and $^1S$-state lifetime values are plotted together with the data for lower states, investigated by Gornik. A slope-of-three line is also included. In this plot it is evident that complex perturbations are present. MQDT calculations have shown that the $4d6s$ $^1D_2$ perturbation is spread out over a large number of Rydberg levels. It would be very worthwhile to see if the apparent shortening of the $6snd$ $^1D_2$ lifetime values for $n \geq 13$ with respect to Gornik's values can be explained by taking this mixing into account.

ACKNOWLEDGMENTS

This work was supported by the Swedish Natural Science Research Council. An interesting discussion with M. Aymar is gratefully acknowledged.

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