

Temperature - Time Curves of Complete Process of Fire Development

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S. E. MAGNUSSON and S. THELANDERSSON

TEMPERATURE-TIME CURVES OF COMPLETE PROCESS OF FIRE DEVELOPMENT

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CIVIL ENGINEERING AND BUILDING CONSTRUCTION SERIES No. 65

Temperature—Time Curves of Complete Process of Fire Development

Theoretical Study of Wood Fuel Fires in Enclosed Spaces

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STOCKHOLM 1970

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Principal notations

\boldsymbol{A}	Area of vertical openings in the en-	
ı	closed space	m^2
A_h	Area of horizontal openings in the en-	
4	closed space	m^2
A_t	Total bounding surface area of the en-	
(4 1/55)	closed space	m ²
$(A \cdot \sqrt[4]{H})$	Air flow factor (Ventilation factor)	$m^{5/2}$
$(A \cdot \sqrt{H}/A_t)$	_ -	$m^{1/2}$
В	Width of an opening in the enclosed	
C	space	m
G	Volume of combustion gases produced	
C	per unit weight of fuel	$m^3 \cdot kg^{-1}$
G_{0}	Volume of combustion gases (expressed	
	in Nm ³) produced per unit weight of	
H	fuel	$\mathrm{Nm^3 \cdot kg^{-1}}$
11	Height of vertical opening in the en-	
H'	closed space	m
11	Height of vertical opening below the	
H"	neutral zone level in the enclosed space	m
11	Height of vertical opening above the	
I	neutral zone level in the enclosed space	m
I_{C}	Enthalpy	kcal ⋅ m ⁻³
¹ C	Heat energy released per unit time	
$I_{\scriptscriptstyle \mathcal{B}}$	during combustion	kcal · h⁻¹
18	Heat energy stored per unit time in the	
	gas volume which is contained in the enclosed space	1 1 1 1
I_L	=	kcal · h⁻¹
*L	Heat energy withdrawn per unit time from the enclosed space owing to the	
	replacement of hot gases by cold air	1 1 1 1 1
I_R	Heat energy withdrawn per unit time	kcal · h⁻¹
	from the enclosed space by radiation	
	. a	1raa1 . 1. = 1
	an owen openings in the enclosed space	kcal · h ^{−1}

I_W	Heat energy withdrawn per unit time	
	from the enclosed space through wall, roof or ceiling, and floor structures	kcal · h ^{−1}
L	Quantity of air consumed per unit weight of fuel during combustion	$Nm^3 \cdot kg^{-1}$
M	Quantity of combustible material	kg
P	Static pressure	kg⋅m ⁻²
$Q_{ m out}$	Rate of flow of the outgoing gases	· ·
≥out	through a vertical opening in the en-	
	closed space	kg⋅h ⁻¹
$Q_{ m in}$	Rate of flow of the incoming air through	
∠ in	a vertical opening in the enclosed space	kg⋅h ⁻¹
Q_h	Rate of flow of the outgoing gases	
z, n	through a horizontal opening in the	
	enclosed space	kg⋅h ⁻¹
Q	Rate of flow of air supplied to the en-	_
2	closed space by means of fans	$m^3 \cdot s^{-1}$
R	Rate of combustion	kg of wood per
		unit time
R_{\max}	Maximum rate of combustion deter-	kg of wood per
	mined by the rate of air supply	unit time
T	Duration of the fire defined as the du-	
	ration of the flame phase	h
W	Heat value of the fuel	kcal ⋅ kg ⁻¹
c	Specific heat	$kcal \cdot m^{-3} \cdot {}^{\circ}C^{-1}$
c_P	Specific heat of the combustion gases	$kcal \cdot m^{-3} \cdot {}^{\circ}C^{-1}$
g	Acceleration of gravity	$m \cdot s^{-2}$
h	Difference in level between the centre of	
	a vertical opening and a horizontal	
	opening	m
h'	Difference in level between the neutral	
	zone and a horizontal opening	m
q	Fire load	Mcal·m ⁻² of bound-
		ing surface area
r	Hydraulic radius	cm
v_y, v_z	Velocity of flow	m ⋅ h ⁻¹
v_h	Velocity of flow through a horizontal	$m \cdot h^{-1}$
	opening in the enclosed space	
t	Time co-ordinate	h
x	Position co-ordinate	m
α_i	Coefficient of heat transfer at a surface	$kcal \cdot m^{-2} \cdot h^{-1} \cdot {}^{\circ}C^{-1}$
	exposed to fire (internal surface)	Keal III II C

C	\mathbf{x}_u	Coefficient of heat transfer at a surface)
		not exposed to fire (external surface)	kcal · m -2 · h -1 · °C-1
2	y	Weight per unit volume	kg·m ⁻³
ε	res	Resultant emissivity for radiation be-	
		tween flames, combustion gases, and a	
		surface exposed to fire (internal surface)	
ε	fl	Emissivity of flames	
3		Emissivity of a surface exposed to fire	
9		Temperature	°C
9		Temperature of the outside air	°C
$\vartheta_{_{_{1}}}$		Temperature of the combustion gases	°C
. 9	i	Temperature of a surface exposed to	-
		fire (internal surface)	°C
ϑ_{ι}	ı	Temperature of a surface not exposed	_
		to fire (external surface)	°C
Δ_{i}	9	Temperature difference between the	-
		combustion gases and the outside air	$^{\circ}\mathrm{C}$
λ		Thermal conductivity	$kcal \cdot m^{-1} \cdot h^{-1} \cdot {}^{\circ}C^{-1}$
μ		Coefficient of contraction	II II
ρ		Density	kg⋅m ⁻³
$ ho_{0}$		Density of the outside air	kg⋅m ⁻³
ρ_g		Density of the combustion gases	kg·m ⁻³
		· ·	

1. Introduction

The efforts made during the past decade in the field of structural fire engineering research have paved the way for differentiated, functionally correct structural fire engineering design carried out on the basis of theoretical calculations. This was rendered possible by investigations which can on the whole be classified in one or several of the main groups enumerated below. At the same time, these groups may be regarded as the essential stages or steps in an appropriate procedure for fire engineering design of load-bearing and separating structures [1].

- (a) Determination of the characteristics of the fire load in an enclosed space under exposure to fire.
- (b) Study of the variations in the development of energy, in the requisite air supply, and in the evolution of gases, with the time in the course of a fire. Determination of the temperature of the combustion gases in the enclosed space as a function of the time.
- (c) Determination of the thermal properties of the materials used for structures in the temperature range which is of interest in connection with fires.
- (d) Determination of the non-stationary temperature fields which are produced in a fire-exposed structure on the assumption that the temperature-time curve for the combustion gases is given, cf. (b).
- (e) Determination of the structural behaviour and the load-bearing capacity of a fire-exposed structure on the basis of the temperature fields defined under (d), and with the help of the available information on those changes in the strength and deformation characteristics of the materials which take place under such conditions.

The object of the present investigation is to make a close study of the stage (b) in order to determine the complete temperature-time curve for the gaseous products of combustion under different conditions, and in particular the temperature-time curve in the cooling phase, 1) for fires of the wood fuel type in enclosed spaces.

¹) The characteristics of the different phases of the process of fire development are represented in Fig. 1. The term "cooling phase" will be used in this publication to designate the smoulder phase and the cooling phase taken together.

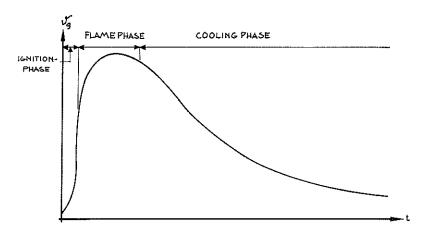


Fig. 1. Phases of the process of fire development as defined in the present publication. The definition in [1] distinguishes between a smoulder phase and a cooling phase, which are here regarded as a single phase designated by the term "cooling phase".

The present state of research in this field is clearly reflected in the sections dealing with fire protection in the Swedish Building Regulations 1967 (abbreviated SBR 67) and in the Draft Specification "Aluminium Structures". In comparison with the relevant regulations which are in force in most other countries, the Swedish rules represent substantial progress on the road to judicious structural fire engineering design. This is primarily due to the fact, that, when the designer has to choose that temperature-time curve which characterises the process of fire development, and which must serve as a basis for all theoretical structural fire engineering design, these rules enables the designer to be guided by all the results, which have been obtained from research in this field during recent years.

On an international plane, it is found that standard temperature-time curves for the process of fire development have been adopted in several countries [1]. If the fuel supply is unlimited, then the agreement between these curves is relatively close, see Fig. 2 a. Under practical conditions, when the fuel supply is limited, the standard specifications used in various countries stipulate that the variation in the temperature with the time shall be in conformity with the standard curve during a certain definite period of time, which is designated by the term "duration of the fire", T, and is defined as the duration of the flame phase. A comparison of the relations between the duration of the fire and the fire load which are employed in various countries is represented in Fig. 2 b [1]. This comparison shows very great differences in an assumption which is fundamental for structural fire engineering design. The wide dispersion between the curves reproduced in Fig. 2 b indicates that there exists no univalued relation between the fire load and the duration of the fire. Concerning

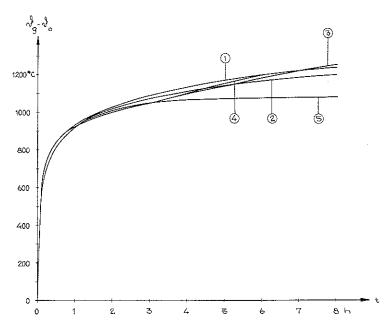


Fig. 2a. Standard curves used in some countries to represent the variation in the temperature, ϑ , in an enclosed space exposed to fire with the time, t. The symbol ϑ_0 denotes the temperature in the enclosed space at the time t=0.

- 1. ISO/TC 92; INSTA 28/2; DIN 4102-62.
- 2. EMPA, Switzerland.
- 3. ASTM 119 (1953), USA.
- 4. V 1076 (1955), Netherlands; BS 476 (1953), United Kingdom.
- 5. A 1304, Japan.

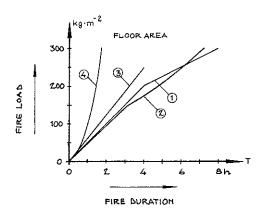


Fig. 2b. Relations between the fire load, in kg of wood per m² of floor area, and the duration of the fire, in h, which are stipulated in the standard specifications used in Sweden (Curve 1), United States of America (Curve 2), and United Kingdom (Curve 3), as well as in a Swiss draft standard specification (Curve 4).

the cooling phase of the process of fire development, it should be noted that it is as a rule completely disregarded.

The relevant Swedish standard specifications provide the designer with three alternative methods of design. Just as in other countries, it is permissible to carry out the design in a roughly simplified and stereotyped manner by using Curve 1 in Fig. 2 a as a point of departure. For the flame phase, this curve gives the temperature of the combustion gases, \mathfrak{F}_g , in the enclosed space in accordance with the equation

$$\theta_g - \theta_0 = 1325 - 430 e^{-0.2t} - 270 e^{-1.7t} - 635 e^{-19t}$$
 (1.1)

where t is the time, in hours, and ϑ_0 is the temperature in the enclosed space at the time t=0. The differences in the combustion characteristics of various fuels, or the fact that the rate of combustion varies within wide limits with the dimensions of the openings in the enclosed space, are not taken into account in this equation. The above-mentioned curve is closely in agreement with that temperature-time curve which is recommended by the ISO for fire tests on building components.

Alternatively, for certain definite types of fire loads and enclosed spaces, the designer may use a method which is simplified, but is nevertheless more differentiated, in comparison with the design procedure outlined in the above. The applicability of this alternative method presupposes that it is possible to comply with the two necessary conditions which are stated in what follows. In the first place, it is required to demonstrate that the characteristics of the fire load in respect of rate of combustion and radiation are approximately in accordance with those which apply in the case of wood fuel. In the second place, it is stipulated that the opening factor of the enclosed space, which is given by the expression $A\sqrt{H}/A_t$, where A is the total opening area of windows and doors, in m², H is a weighted average of the vertical dimensions of these openings, in m, and A_t is the total area of the surfaces bounding the enclosed space, in m², shall be known during all phases of the process of fire development. If these two conditions are satisfied, then it is allowed to carry out the design on the basis of a specific curve representing the variation in the temperature of the combustion gases in the enclosed space, ϑ_g , with the time. For the flame phase, this curve is determined by the opening factor, see Fig. 2 c, in the course of the duration of the fire, T, which is defined by the equation

$$T = qA_t/(25A\sqrt{H}) \quad \min \tag{1.2}$$

where q is the fire load, in Mcal·m⁻² of bounding surface area. The dash-line curve represents the INSTA curve expressed by Eq. (1.1).

Finally, and generally, in the cases where the quantities of combustible materials which constitute the fire load, as well as the rate of combustion, are accurately known, the above-mentioned two Swedish specifications allow

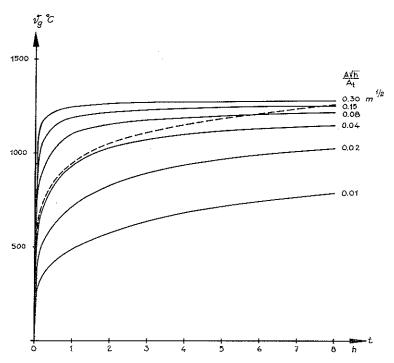


Fig. 2c. Variation in the temperature of the combustion gases, ϑ_g , with the time, t, at different values of the opening factor, $A \cdot \sqrt{H}/A_t$. Curves published in the Swedish Building Regulations 1967 (SBR 67) and in the Swedish draft specification "Aluminium Structures". The dash-line curve is the standard curve calculated by means of Eq. (1.1).

the fire resistance of a building component to be determined on the basis of the variation in the temperature of the combustion gases with the time, calculated with the help of the formula which is known as the equation of heat balance. This equation, which constitutes a fundamental description of the energy balance of the process of fire development, and hence also serves as a basis for Eq. (1.2), will be discussed at some length in Chapter 2.

What has been said up to this point relates only to the ignition and flame phases of the process of fire development. As regards the cooling phase, it is stipulated in a summary manner merely that the time graph of the temperature of the combustion gases shall be chosen so as to be linear, and that the rate of decrease in the gas temperature shall be taken to be 10°C·min⁻¹, unless other assumptions can be demonstrated to be more correct. If the design is carried out by means of the second or third alternative method, each of which is functionally realistic, then this implies that two phases of the same continuous process are represented in such a way that their descriptions are entirely different in the degree of accuracy as well as in the extent to which the actual conditions are taken into account. It is obvious that this gives

rise to a considerable unbalance in the basis for design. To show how important it is that the cooling phase should also be described in a differentiated manner, it may be useful to give two examples of structural members which are characterised by low and high thermal inertia, respectively.

For an enclosed space, where the fire load is $q=12 \text{ Mcal} \cdot \text{m}^{-2}$ of bounding surface area, and the opening factor is $A\sqrt{H}/A_t=0.08 \text{ m}^{0.5}$, Fig. 3 represents a calculated temperature-time curve for a steel column exposed to fire and characterized by the ratio $F_s/V_s=100 \text{ m}^{-1}$ and by $\varepsilon_r=0.5$, where F_s is the total bounding surface area of the column, which is equal to its fire-exposed area, in m^2 , V_s is the steel volume of the column, in m^3 , and ε_r is the resultant emissivity for heat transfer from the flames and the combustion gases to the

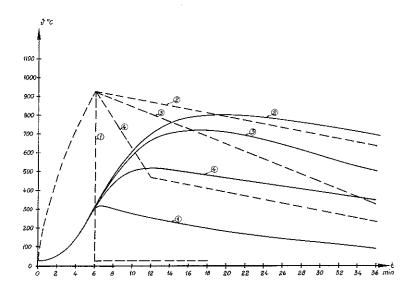


Fig. 3. Variation in the temperature, ϑ_s , of a steel column with the time, t, calculated on the basis of temperature-time curves, which differ in shape during the cooling phase (dash-line curves).

 $F_s/V_s = {
m Ratio}$, in m⁻¹, of the fire-exposed surface area, i.e. the total bounding surface area, of the column to the steel volume of the column.

$$\varepsilon_r$$
 = Resultant emissivity.
 ϑ_s ---- ϑ_g at $A\sqrt{H}/A_t$ =0.08 m^{1/2}.
 q = 12 Mcal·m².

- Instantaneous cooling. Rate of decrease in temperature ∞ °C · h⁻¹.
- 2. Linear rate of decrease in temperature, 600°C · h⁻¹.
- 3. Linear rate of decrease in temperature, 1200°C · h-1.
- 4. Linear rate of decrease in temperature, 6 min $\leq t \leq 12$ min, 4600°C·h⁻¹; $t \geq 12$ min, 600°C·h⁻¹.

$$F_s/V_s = 100 \text{ m}^{-1} \ \varepsilon_r = 0.5.$$

steel column. The characteristics of the flame phase for the temperature-time curve of the enclosed space have been chosen in conformity with the Swedish Building Regulations 1967. In Fig. 3, the full-line curves represent the temperature of the steel, θ_s , and the dash-line curves show the temperature of the combustion gases, θ_g , on the basis of the four alternative assumptions concerning the cooling phase of the process of fire development which are stated in what follows.

- (1) After the duration of the fire $T = qA_t/(25A\sqrt{H}) = 6$ min, the temperature of the combustion gases drops instantaneously to ordinary room temperature.
- (2) The temperature of the combustion gases decreases in accordance with the Swedish Building Regulations 1967 at a linear rate of 10°C · min⁻¹.
- (3) The temperature of the combustion gases decreases at a linear rate of $20^{\circ}\text{C} \cdot \text{min}^{-1}$.
- (4) The temperature of the combustion gases is assumed to vary in a more realistic manner, that is to say, it drops to half their maximum temperature during the first 6 min of the cooling phase, and then decreases at a linear rate of $10^{\circ}\text{C} \cdot \text{min}^{-1}$.

At the end of the flame phase, the temperature of the steel is 303°C. After that, the temperature of the steel continues to increase during the cooling phase of the process of fire development to its respective maximum values corresponding to the four alternative assumptions, viz., 303, 799, 719, and 518°C.

Fig. 4 shows the effects on the load-bearing capacity of a reinforced concrete slab which are produced by different slopes of the linear cooling phase of the time-temperature curves. Each T-value on the horisontal axis corresponds to a specific time-temperature curve. T is the duration of the flame phase, and depends on the fire load. The distance from the centre lines of the reinforcing bars to the fire-exposed surface of the slab, is assumed to be 2 cm. The emissivity of the flames is taken to be 0.7. The temperature in the enclosed space is supposed to vary with the time during the flame phase in accordance with Eq. (1.1). The temperature of the combustion gases is assumed to decrease at linear rates of 5, 10, and $20^{\circ}\text{C} \cdot \text{min}^{-1}$, or to drop instantaneously to ordinary room temperature ($\infty^{\circ}\text{C} \cdot \text{h}^{-1}$). For a static load which causes failure at a temperature of the reinforcing bars $\vartheta_{scr} = 450^{\circ}\text{C}$, we obtain a fire resistance period, t_{fr} , which varies from 0.52 to 0.82 h, and for $\vartheta_{scr} = 500^{\circ}\text{C}$, the corresponding variation in the fire resistance period ranges from 0.72 to 1.01 h.

These examples show that it is necessary to calculate the fire resistance period of a building component so as to take account of that reduction in its load-bearing capacity, or in its separating capacity, which occurs during the cooling phase. Moreover, they indicate that those temperature-time curves for enclosed spaces which are to serve as a basis for such calculations should also be differentiated and as realistic as possible in the cooling phase.

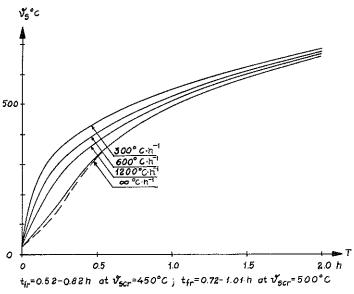


Fig. 4. Relation between the calculated maximum temperature, ϑ_s , of the reinforcing steel in a concrete slab, 18 cm in thickness, exposed to fire on one side, and the duration, T, of the flame phase of the process of fire development. This relation is represented for different values of the fire load and for different slopes of the linear cooling phase of the temperature-time curve.

Cooling phase taken into account.

---- Cooling phase not taken into account.

Concrete cover 2 cm.

 ε_{fl} = Emissivity of the flames=0.7.

 t_{fr} = Fire resistance period.

 $\vartheta_{scr} = \text{Critical temperature of the steel, i.e. the temperature at which the reinforcement fails.}$

 $t_{fr} = 0.52 \text{ to } 0.82 \text{ h at } \vartheta_{scr} = 450 ^{\circ}\text{C}.$

 $t_{fr} = 0.72 \text{ to } 1.01 \text{ h at } \vartheta_{sc_r} = 500^{\circ}\text{C}.$

The object of the present investigation is therefore to evolve a method which shall be applicable to different combinations of the values of the air flow factor (ventilation factor), $A\sqrt{H}$, and the fire load, as well as the type of material of the structures bounding the enclosed space under exposure to fire of the wood fuel type. This shall enable that the temperature-time curves for the enclosed space to be calculated by means of a theoretical procedure so as to cover the whole process of fire development, and shall thus make it possible to carry out judicious structural fire engineering design on the basis of the variation in the temperature with the time during all phases of the fire.

In connection with a general treatment of the equation of heat balance, Chapter 2 deals with the problems which are met with when a theoretical calculation of the temperature-time curve for the combustion gases is extended so as to comprise the cooling phase. Chapter 3 describes the methods which have been used to tackle these problems, and the modifications of the equation of heat balance which have been necessary for this purpose. In Chapter 4, an account is given of the computer programme which has been prepared for the calculations, and which is represented in the form of a flow chart. In Chapter 5, the full-scale tests which have served as a basis for the present investigation are subjected to comparative theoretical analysis. The time graphs of the rate of combustion which have been determined with the help of these theoretical analyses are presented in Chapter 6. Finally, in Chapter 7, these graphs are used as a basis for the calculation of complete temperature-time curves for combustion gases in enclosed spaces which vary in the values of the opening factor and the fire load, as well as in type of material employed in the structures bounding the enclosed space.

2. Equation of heat balance of process of fire development

The papers published by Kawagoe and Sekine [2], as well as by Ödeen [3], in the early 1960ies have made it possible to carry out theoretical calculations of temperature-time curves for combustion gases in the flame phase of the process of fire development to a degree of accuracy that is sufficient for practical purposes. These three authors have studied the energy balance during the process of fire development. The quantity of energy released per unit time, just as the volume of combustion gases evolved, during combustion were assumed to be known. With the help of the calculation of the quantity of energy which was lost per unit time by conduction and radiation from the enclosed space through its bounding structures, it was possible to deduce an equation of heat balance, and to solve it so as to obtain the temperature of the combustion gases. The treatment of this problem was based on the simplified assumptions which are reproduced in what follows.

- (a) The temperature in the interior of the whole enclosed space is uniform at any given instant.
- (b) The coefficient of heat transfer to the interior bounding surfaces of the enclosed space is uniform at every point.
- (c) The heat flow through the bounding structures of the enclosed space is one-dimensional and, except for the window and door openings, if any, uniformly distributed.

Kawagoe and Sekine, as well as Ödeen, confined themselves throughout their papers to a study of the flame phase of the process of fire development, and the equation which they have deduced cannot be applied directly to the cooling phase. Primarily, the calculation of the temperature-time curve for the combustion gases during the cooling period requires an analytical investigation of two fundamental sub-problems, which have been but little studied up to the present time. In the first place, it is necessary to determine the quantity of energy liberated per unit time when this quantity, as is the case in the cooling phase, is no longer determined by the rate of air supply. In the second place, it is required to investigate the thermodynamic conditions which are encountered when the rate of combustion is no longer limited by the dimensions of the openings in the enclosed space. The equation of energy balance which has been deduced by Kawagoe and Sekine and by Ödeen, as

well as that extension of the theory which is required for the calculation of the temperature-time curve in the cooling phase of the process of fire development, will be dealt with in what follows.

The above-mentioned equation expresses for any given instant ,t, the balance between the respective quantities of heat energy generated and lost per unit time in the enclosed space under consideration. In its complete form, this equation is

$$I_C = I_L + I_W + I_R + I_B$$
 where (2.1)

 I_C = the heat energy released per unit time during combustion,

 I_L = the heat energy withdrawn per unit time from the enclosed space owing to the replacement of hot gases by cold air,

 I_W = the heat energy withdrawn per unit time from the enclosed space through wall, roof or ceiling, and floor structures,

 I_R = the heat energy withdrawn per unit time from the enclosed space by radiation through the openings in the enclosed space,

 I_B = the heat energy stored per unit time in the gas volume which is contained in the enclosed space.

The terms entering into the above equation are schematically illustrated in Fig. 5.

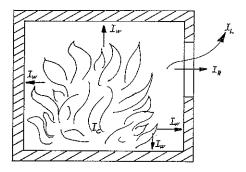


Fig. 5. Schematic illustration of the terms entering into the equation of heat balance.

Term I_B

In comparison with the quantities of energy which are involved in fires, the quantity of energy which can be stored in the gas volume contained in the enclosed space is of minor importance. Therefore, the term I_B can to a close approximation be put equal to zero.

Term I_C

In the case of fires of the wood fuel type, the term I_c , which expresses the quantity of heat released per unit time during combustion, is a factor which is difficult to determine in the equation of heat balance. In order to obtain this term, Kawagoe and Sekine, as well as Ödeen, chose as a point of departure the rate of combustion, expressed in terms of kilogrammes of wood per unit time, which they multiplied by a heat value.

This gives the equation

$$I_C = R \cdot W \quad \text{kcal} \cdot h^{-1}$$
 where

R = the rate of combustion, in kg of wood per h, W = the heat value, in kcal per kg of wood.

With the exception of a small number of calculations which were based on a triangular variation in the rate of combustion, R, with the time, Ödeen assumed that R is constant, and is arbitrarily chosen, and that the quantity of heat, W, released during the flame phase is 4120 kcal per kg of wood. Thus, the heat energy liberated per unit time during combustion, I_C , is supposed to remain constant until the fuel has burnt up, and the cooling phase is characterized by the fact that no additional energy is supplied to the enclosed space. This description of the quantity of energy developed per unit time is applicable to burning of liquid fuels, but does not take account of the real combustion characteristics of wood fuels. As a rule, these fuels give off 30 to 50 per cent of the total quantity of energy after the end of the flame phase.

The investigation made by Kawagoe and Sekine is more differentiated in respect of the characterization of the process of fire development. As *Kawagoe* had shown in an earlier paper [5], during that period of the fire when the rate of combustion reaches a maximum, i.e. during the flame phase, the rate of air supply to the enclosed space, and hence also the maximum rate of combustion, R_{max} , are proportional to the air flow factor, $A \sqrt{H}$.

If the areas are expressed in m^2 , and the maximum rate of combustion, R_{max} , is expressed in kg of wood per min, then we have the approximate relation

$$R_{\text{max}} = 5.5 \cdot A \sqrt{H}$$
 kg of wood per min (2.3)

Furthermore, in the papers published by Kawagoe and Sekine, the quantity of heat liberated during the flame phase, W, is stated to be 2575 kcal per kg of wood. This value was obtained by reducing the nominal heat value of wood so as to take account of the degree of incomplete combustion. The degree of incomplete combustion was estimated with the help of those analyses of the composition of the combustion gases which were carried out during fire

tests [5]. As regards the cooling phase, Kawagoe and Sekine had made a few isolated comparative calculations based on a polygon-shaped time graph of the rate of combustion, and then found that the temperatures obtained when the cooling phase was characterized by a linear decrease in temperature at a rate of 7 or $10^{\circ}\text{C} \cdot \text{min}^{-1}$ were much too high.

Accordingly, if the results of the investigations made by Ödeen, as well as by Kawagoe and Sekine, are to be applied to fires where the fuel is of the wood type, then the calculations have to be confined to the flame phase of the process of fire development.

When the treatment of this problem is extended so as to comprise the cooling phase of the process of fire development, Eq. (2.3) is not generally applicable. For the quantity of energy released per unit time during combustion, Eq. (2.3), in combination with Eq. (2.2), gives only the theoretical upper limit, which is determined by the available rate of air supply. During the cooling phase, the energy liberated per unit time will be governed by other factors. For this reason, and since no systematic investigation has so far been made in order to determine the relations between the three quantities which are of interest in this connection, viz., the reduction in the weight of fuel, the quantity of energy developed per unit weight of fuel, and the requisite rate of air supply, the quantity of energy released per unit time during the cooling phase of the process of fire development had to be determined by means of the method described in what follows. The procedure in calculation for the solution of the equation of heat balance was programmed for a CD 3600 computer. A study of the literature was carried out in order to examine the available publications on full-scale tests. A number of these tests were selected in the cases where the reported data were so complete as to enable numerical treatment. After that, the computer was used to calculate the temperature-time curve on the basis of an assumed form of the time graph of the rate of combustion for the complete process of fire development. The time graph of the rate of combustion was then varied until the agreement between the experimental and theoretical temperature-time curves was as close as possible. The only absolute requirement to be fulfilled in this connection was that the total quantity of energy liberated during the whole process of fire development should be equal to the total energy of combustion of the fuel. When an adequate range of variation in the opening factor and in the fire load was considered to have been covered, the time graphs of the rate of combustion obtained in this way were systematized. For a given fire load and a given opening factor, it was then possible to assume that the curve showing the variation in the rate of combustion with the time was known on the basis of this systematization. The investigation referred to in the above is described in Chapter 5, and its results, expressed in terms of time graphs of the rate of combustion in a dimensionless form, are presented in Chapter 6.

In connection with the treatment of the term I_c , it should be pointed out that the fire load must be described in combustion engineering terms in such a way that it may be associated with the equation of heat balance, Eq. (2.1).

In most countries, the fire load is expressed in terms of the quantity of wood that is equivalent to it in heat value per unit floor area. This characterization must be replaced by a parameter which has a physical significance when it is treated in calculations. This has been done in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures", which stipulate that the fire load shall be stated as that total quantity of heat, q, in Mcal·m⁻² of total bounding surface area of the enclosed space exposed to fire, which is liberated on the assumption of complete combustion of all the combustible material contained in the enclosed space. A still more refined description of the fire load, which should take into account the variation with the time in the quantity of energy released by combustion, as well as the emissivities of the flames and the combustion gases, is an urgently recommended subject for research in this field.

Term I_L

The term I_L in the equation of heat balance expresses that quantity of heat which is withdrawn per unit time from the enclosed space owing to the replacement of hot gases by cold air through the openings in the enclosed space. For the determination of I_L , Ödeen, as well as Kawagoe and Sekine, used the equation

$$I_L = R \cdot G_0 \cdot (\vartheta_g - \vartheta_0) \cdot c_p \tag{2.4}$$
 where

R = the rate of combustion, in kg of wood per h,

 G_0 = the volume of combustion gases produced by the fire, in Nm³ · kg⁻¹ of fuel,¹)

 c_p = the specific heat of the combustion gases, in kcal·Nm⁻³·°C⁻¹,

 θ_g = the temperature of the combustion gases, in °C,

 θ_0 = the temperature of the air outside the enclosed space, in °C.

Eq. (2.4) states that the term I_L is put equal to the heat content of the combustion gases produced by the fire with reference to that of the outside air. On account of the difference in density between the cold outside air and the hot gases in the interior of the enclosed space, an exchange of heat by convection takes place in the openings of the enclosed space. The rate of this heat exchange determines the maximum value of the rate of combustion so far as the supply of oxygen is concerned. Therefore, Eq. (2.4) can be used as an expression for I_L when the rate of combustion is determined by the rate

¹⁾ Nm³=normal cubic metre=the quantity of a gas which occupies a volume of 1 cubic metre at 0°C and 760 mm barometric pressure.

of air supply, in spite of the fact that this equation in itself does not describe I_L, but expresses the heat content of the combustion gases produced by the fire. However, if the rate of combustion is limited by factors other than the rate of air supply, e.g. by the available quantity of fuel, or if the combustion is completed, then it is obvious that Eq. (2.4) does not hold good. Accordingly, a theoretical treatment of the cooling phase of the process of fire development requires an expression for I_L which is more generally applicable, and which is based on the rate of air exchange. This problem has been studied in a thesis for degree of Master of Engineering prepared by Ahlquist and Thelandersson [6], in the Division of Structural Mechanics and Concrete Construction, Lund Institute of Technology, Lund, Sweden. These authors based their study on the assumption that the static pressure distribution in the enclosed space varies linearly from the floor to the ceiling or roof, and that there exists a level (the neutral layer) at which the pressure in the enclosed space is equal to the pressure outside the enclosed space. This simplified model has been used by several authors, e.g. by Kawagoe [5] in the theoretical deduction of Eq. (2.3), and by Thomas [7] in studies which dealt with venting in the course of fires. In this connection, Kawagoe showed by means of a number of tests that the assumed pressure was actually applicable to a close approximation. On the assumption that the openings in the enclosed space are vertical only, Kawagoe used Bernoulli's equation to determine the quantities of outgoing combustion gases and incoming cold air as functions of the difference in temperature and the position of the neutral layer. On the basis of the condition that the difference between the quantities of gases flowing into and out of the enclosed space shall be equal to the difference between the quantities of gases produced and consumed by combustion, Ahlquist and Thelandersson calculated the position of the neutral layer as a function of the temperature and the rate of combustion. In this calculation, it was assumed that the rate of combustion may vary from zero to a maximum value, which is dependent on the dimensions of the openings. Furthermore, it was assumed that the liberation of a certain definite quantity of energy is associated with the consumption of the same quantity of air and the production of the same quantity of combustion gases, irrespective of the rate of combustion.

After the position of the neutral layer had been determined in this way, it was possible to obtain an expression for I_L at different values of the temperature and the rate of combustion. In the above-mentioned thesis [6], the treatment was also extended so as to comprise the modifications which are necessary when the enclosed space is provided with a vent. The deduction of the equations which are required for this purpose is reproduced in its main features in Chapter 3 of the present publication. Moreover, this chapter also contains a summary treatment of the case where the roof of the enclosed space is provided with horizontal openings.

Term I_w

The term I_W denotes the quantity of heat which is withdrawn per unit time from the enclosed space through the structures bounding this space. The term I_W is determined by solving the general equation of heat conduction in the one-dimensional case under non-steady flow conditions so as to take into account those thermal properties of the materials which are dependent on the temperature, the evaporation of occluded water, and the possible structural transformations in the materials entering into the bounding structures. This equation is

$$c \cdot \gamma \cdot \frac{\partial \vartheta}{\partial t} = \frac{\partial}{\partial x} \left(\lambda_x \cdot \frac{\partial \vartheta}{\partial x} \right) \tag{2.5}$$

where

c = the specific heat of the wall material,

 γ = the weight per unit volume of the wall material,

 λ_x = the thermal conductivity of the wall material,

9 = the temperature in the interior of the wall material,

t =the time,

x = the position co-ordinate.

The above equation is solved by means of a numerical procedure which had been described by *Odemark* [8], among others, and which has subsequently been further developed by *Ödeen*, and others, [9]. The walls, ceiling, roof, and floor structures which bound the enclosed space are divided into n layers having the thickness Δx_k each, and the equation of heat balance is written for each one of these layers. If the temperature at the centre of the layer k at the time t is denoted by θ_k , and that at the time $t + \Delta t$ is designated by $\theta_k + \Delta \theta_k$, then the application of this procedure gives the relations

$$\varphi_{1} \cdot \frac{\Delta \vartheta_{1}}{\Delta t} = \psi_{1}(\vartheta_{g} - \vartheta_{1}) - \psi_{2}(\vartheta_{1} - \vartheta_{2})$$

$$\vdots$$

$$\varphi_{k} \cdot \frac{\Delta \vartheta_{k}}{\Delta t} = \psi_{k}(\vartheta_{k-1} - \vartheta_{k}) - \psi_{k+1}(\vartheta_{k} - \vartheta_{k+1})$$

$$\vdots$$

$$\vdots$$

$$\varphi_{n} \cdot \frac{\Delta \vartheta_{n}}{\Delta t} = \psi_{n}(\vartheta_{n-1} - \vartheta_{n}) - \psi_{n+1}(\vartheta_{n} - \vartheta_{0})$$

$$(2.6)$$

where

 θ_a = the temperature of the gases in the enclosed space,

 θ_0 = the temperature of the outside air,

$$\varphi_k = \Delta x_k \cdot c(x, \vartheta) \cdot \gamma$$

$$\psi_1 = \frac{1}{\frac{1}{\alpha_i(\vartheta)} + \frac{\Delta x_1}{2 \cdot \lambda(x, \vartheta)}}$$

$$\psi_k = \frac{1}{\frac{\Delta x_{k-1}}{2 \cdot \lambda(x, \vartheta)} + \frac{\Delta x_k}{2 \cdot \lambda(x, \vartheta)}}$$

$$\psi_{n+1} = \frac{1}{\frac{\Delta x_n}{2 \cdot \lambda(x, \vartheta)} + \frac{1}{\alpha_n(\vartheta)}}$$

where

 $\alpha_i(9)$ = the coefficient of heat transfer at the internal surface,

 $\alpha_u(\theta)$ = the coefficient of heat transfer at the external surface,

 $\lambda(x, \theta)$ = the thermal conductivity at the section x,

c(x, 9) = the specific heat at the section x,

y = the weight per unit volume at the section x.

The coefficient of heat transfer, α_i , at the internal surface exposed to fire may be supposed to consist of two components, viz., first, a radiation component, which is markedly predominant at the high temperatures in question, and second, a convection component, which can be chosen with adequate accuracy so as to be constant, and to be equal to 20 kcal·m⁻²·h⁻¹·°C⁻¹ [9]. By applying the Stefan-Boltzmann law, this gives, for α_i , the relation

$$\alpha_{i} = \frac{4.96 \cdot \varepsilon_{res}}{\vartheta_{g} - \vartheta_{i}} \left[\left(\frac{\vartheta_{g} + 273}{100} \right)^{4} - \left(\frac{\vartheta_{i} + 273}{100} \right)^{4} \right] + + 20 \text{ kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$$
(2.7a)

where

 ϑ_i = the temperature of the internal surface,

 ε_{res} = the resultant emissivity for radiation between flames, combustion gases, and the internal surface.

The resultant emissivity, ε_{res} , is determined from the formula

$$\frac{1}{\varepsilon_{res}} = \frac{1}{\varepsilon_{fl}} + \frac{1}{\varepsilon_i} - 1 \tag{2.7b}$$

where

 ε_{fl} = the emissivity of the flames,

 ε_i = the emissivity of the surface exposed to fire.

Properly speaking, Eq. (2.7b) represents the emissivity for radiation between two parallel surfaces, but it was considered to be the best available approximation in the case of radiation between flames and a surface exposed to fire.

According to [9], the coefficient of heat transfer, α_u , at the external surface, which is not exposed to fire, can be represented by the approximate expression

$$\alpha_u = 7.5 + 0.028 \cdot \theta_u \quad \text{kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$$
 where

 ϑ_u = the temperature of the external surface.

The system of differential equations of the first order, Eq. (2.6), is solved numerically (see Chapter 4), and then the term I_W is given by the relation

$$I_W = A_t \cdot \psi_1 \cdot (\theta_a - \theta_1) \tag{2.9}$$

where

 A_t =the total area of the surfaces bounding the enclosed space.

If the structures bounding the enclosed space consist of different materials, or if they differ in thickness, as is usually the case in practice, then the above-mentioned operations are carried out separately for each type of structure, and after that the term I_W is obtained from the expression

$$I_{W} = \sum_{i} I_{W, j} = \sum_{i} A_{j} \cdot \psi_{1, j} \cdot (\vartheta_{g} - \vartheta_{1, j})$$

$$(2.10)$$

The present section, which deals with the term I_W , is based in its entirety on the publications of Kawagoe and Sekine as well as Ödeen referred to in the above.

Term I_R

Kawagoe and Sekine calculated the term I_R from the following formula, which is a generalization of the Stefan-Boltzmann law:

$$I_R = A \cdot (E_q - E_0) \tag{2.11}$$

where

A = the area of the opening,

$$E_g = 4.96 \cdot \left(\frac{9_g + 273}{100}\right)^4$$

$$E_0 = 4.96 \cdot \left(\frac{9_0 + 273}{100}\right)^4$$

This formula is applicable to the whole duration of the process of fire development, and is used in its unchanged form for the calculations in the present publication.

3. Study of combustion gas flow and heat flow through openings in enclosed spaces

As has been shown in the section dealing with the term I_L in Chapter 2, the expression given by Eq. (2.4) is applicable only during the flame phase of the process of fire development. In that section, the term I_L was determined on the basis of a maximum rate of combustion. If this expression is to be extended so as to be valid for the cooling phase of the fire also, then this requires that the rates of gas and air flow through the openings in the enclosed space shall be determined directly, and that these rates shall then be used as a point of departure for determining the quantity of heat which is withdrawn per unit time from the enclosed space. Similar problems, which were defined in thermodynamic terms, and which related to fires in enclosed spaces, have been treated by Kawagoe [5], among others. The primary prerequisites to such a treatment are the assumptions that the pressure distribution in a vertical direction is linear, and that there exists a neutral layer or zone, i.e. a level at which the static pressure in the interior of the enclosed space is equal to the atmospheric pressure outside the enclosed space. From these assumptions, Kawagoe deduced the expression for the maximum rate of combustion, R_{max} , which is given by Eq. (2.3). On the assumption that the position of the neutral zone is the unknown variable, he used the Bernoulli equation to calculate the respective quantities of gases and air which flow out of and into the enclosed space per unit time. After that, he determined the position of the neutral zone from the condition that the rate of flow of the incoming gases shall be equal to the rate of flow of the gases which are consumed by combustion, and that the rate of flow of the outgoing gases shall be equal to the rate of flow of the gases which are produced by combustion. Finally, on the assumption that the quantity of air consumed per unit weight of fuel is known, he calculated the maximum rate of combustion, R_{max} .

In this chapter, a similar analysis will be carried out in what follows. The purpose of this analysis is to determine the term I_L by an expression which is more general than that given by Eq. (2.4). A detailed deduction will be presented for the case where the enclosed space is provided with one or several vertical openings which are equal in height. After that, we shall deal with the modifications which are required in the applications which involve openings of other types.

Flow conditions in vertical rectangular opening in enclosed space

The interchange between the gaseous products of combustion and the combustion air takes place because the density of the hot gases is lower than that of the cold air outside the enclosed space. On the assumption that the temperature in the whole enclosed space is uniform, and that there exists a neutral zone, the velocities of gas and air flow can be determined theoretically. After that, if the dimensions of the opening are known, it is possible to calculate the respective rates of flow, i.e. the masses of the outgoing gases and the incoming air per unit time.

The velocity distribution in a vertical rectangular opening is schematically represented in Fig. 6.

The difference in static pressure between the outside and the inside is equal to zero at the level of the neutral zone. Accordingly, if use is made of the notations given in Fig. 6, the pressure difference, P_y , above the neutral zone i

$$P_{y} = (\rho_0 - \rho_q) \cdot y \tag{3.1a}$$

and the pressure difference, P_z , below the neutral zone is

$$P_z = -(\rho_0 - \rho_a) \cdot z \tag{3.1b}$$

where

 ρ_0 = density of the outside air,

 ρ_a =density of the combustion gases.

The density of the gaseous products of combustion is assumed to be equal to the density of the air at the same temperature [10].

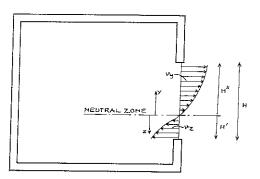


Fig. 6. Gas flow in an enclosed space provided with a vertical opening.

A = Area of the vertical opening in the enclosed space.

H = Height of the vertical opening in the enclosed space,

H'' = Height of the vertical opening above the neutral zone level.

H' = Height of the vertical opening below the neutral zone level.

 v_v = Velocity of gas flow above the neutral zone level.

 v_z = Velocity of gas flow below the neutral zone level.

From Bernoulli's theorem we obtain the following expressions for the variation in the velocity of flow with the distance from the neutral zone

$$v_{y} = \sqrt{2g \cdot y \cdot \frac{\rho_{0} - \rho_{g}}{\rho_{g}}} \tag{3.2a}$$

and

$$v_z = \sqrt{2g \cdot z \cdot \frac{\rho_0 - \rho_g}{\rho_0}} \tag{3.2b}$$

Then the rate of flow of the outgoing gases is

$$Q_{\text{out}} = \mu \cdot B \cdot \rho_g \cdot \int_0^{H''} v_y \cdot dy \tag{3.3a}$$

and the rate of flow of the incoming air is

$$Q_{\rm in} = \mu \cdot B \cdot \rho_0 \cdot \int_0^{H'} v_z \cdot dz \tag{3.3b}$$

where

 μ = the coefficient of contraction,

B = the width of the opening.

By substituting Eqs. (3.2a) and (3.2b) in Eqs. (3.3a) and (3.3b), respectively, we can directly calculate Q_{out} and Q_{in} . We find

$$Q_{\text{out}} = \frac{2}{3} \cdot \mu \cdot B \cdot (H'')^{3/2} \cdot \sqrt{2g \cdot \rho_g \cdot (\rho_0 - \rho_g)}$$
(3.4a)

and

$$Q_{\rm in} = \frac{2}{3} \cdot \mu \cdot B \cdot (H')^{3/2} \cdot \sqrt{2g \cdot \rho_0 \cdot (\rho_0 - \rho_g)}$$
 (3.4b)

The position of the neutral zone is determined by the equation of gas interchange in the enclosed space. This equation states that the difference between the rates of flow of the outgoing gases and the incoming air shall be equal to the difference between the rates of flow of the gases which are produced and consumed by combustion.

The mass of air contained in the enclosed space is assumed to be constant during the whole period of time under consideration. That total error, referred to the whole duration of the process of fire development, which is caused by this assumption in the calculation of heat flow is not greater than the heat content of the volume of air in the enclosed space. In comparison with the quantities of heat which are associated with fires, and in view of the other approximations which have been made in connection with the application of Eq. (2.1), this error may be regarded as negligible.

Deduction of maximum rate of combustion determined by rate of air supply in accordance with [5]

On the assumption that the rate of combustion, R, is determined by the rate of air supply $(R = R_{max})$, we have

$$Q_{\text{out}} = R_{\text{max}} \cdot G_0 \cdot \rho_0 \tag{3.5a}$$

and

$$Q_{\rm in} = R_{\rm max} \cdot L \cdot \rho_0 \tag{3.5b}$$

where

 G_0 = the volume of combustion gases, in Nm³, produced by the combustion of 1 kg of fuel,

L = the volume of air, in Nm³, consumed by the combustion of 1 kg of fuel.

By substituting Eqs. (3.4a) and (3.4b) in Eqs. (3.5a) and (3.5b), respectively, we can calculate the position of the neutral zone, i.e. H' and H''. After that, by substituting H' in Eq. (3.5b), we obtain, for R_{max} , the expression

$$R_{\text{max}} = \kappa(\Delta \theta) \cdot A \cdot \sqrt{H}$$
 (3.6)

where

 $\kappa(\Delta \theta)$ is a coefficient, which depends on $\Delta \theta$

and

$$\Delta \theta = \theta_q - \theta_0$$

The values of $\kappa(\Delta\theta)$ is calculated for two fuels. First, for fires of the wood fuel type, which are most characteristic of actual fires, because wood usually constitutes the predominant fire load. Second, for fires, where the fuel consists of alcohol. The numerical values used in these calculations are given in what follows [10].

$$\mu$$
=0.7
 G_0 =4.86 Nm³ · kg⁻¹ for wood
6.22 Nm³ · kg⁻¹ for alcohol
 L =3.98 Nm³ · kg⁻¹ for wood
5.23 Nm³ · kg⁻¹ for alcohol

The relation between κ and $\Delta \vartheta$ is represented in Fig. 7. This graph shows that the variation in the value of κ with the temperature is very slight in the temperature range which is met with in fires. For practical applications, Kawagoe put the value of κ for wood fires equal to 330 kg·h⁻¹·m^{-5/2}, irrespective of the temperature. The value in question was used in deducing Eq. (2.3). This equation has been verified experimentally by several authors in model tests as well as in full-scale tests, see [5]. Eqs. (3.6) and (2.3) are applicable only to one or several openings of equal height, H. In these equations, A denotes the sum of the areas of the individual openings.

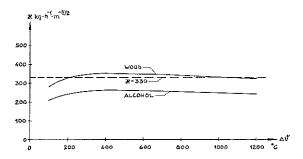


Fig. 7. Relation between the coefficient \varkappa in the equation $R_{\max} = \varkappa A \sqrt{H}$ and the temperature difference, $\Delta \vartheta$, between the combustion gases and the outside air.

Wood fires and alcohol fires.

Determination of quantity of heat, I_L , withdrawn per unit time through openings in enclosed space during the whole process of fire development

In order that an expression may be valid throughout the duration of the process of fire development, it is necessary to presuppose that the rate of combustion, R, may assume all values in the interval extending from zero to the maximum value which is given by Eq. (2.3). Therefore, R can be written

$$R = a \cdot 330 \cdot A \cdot \sqrt{H}$$
 kg · h⁻¹

where

$$0 < a \le 1$$

The balance between the respective rates of flow of the outgoing gases and the incoming air is given by the relation

$$Q_{\text{out}} - Q_{\text{in}} = G_0 \cdot R \cdot \rho_0 - L \cdot R \cdot \rho_0 \tag{3.8}$$

After substitution of Q_{out} and Q_{in} from Eqs. (3.4a) and (3.4b), respectively, and after simplification, we obtain

$$\left(\frac{H''}{H}\right)^{3/2} \cdot \sqrt{\rho_g(\rho_0 - \rho_g)} - \left(1 - \frac{H''}{H}\right)^{3/2} \cdot \sqrt{\rho_0(\rho_0 - \rho_g)} = \frac{(G_0 - L) \cdot \rho_0 \cdot 330a}{\frac{2}{3}\mu \cdot \sqrt{2g}} \tag{3.9}$$

With the help of Gay-Lussac's law of volumes for gases

$$\rho_0 = \rho_g \left(1 + \frac{\Delta \vartheta}{273} \right) \tag{3.10}$$

the ratio H''/H can be determined for different values of a and $\Delta\theta$ from Eq. (3.9).

The quantity of heat which is withdrawn per unit time from the enclosed space can be written

$$I_L = Q_{\text{out}} \cdot c_p \cdot \frac{\Delta \vartheta}{\rho_0} \tag{3.11a}$$

where

 c_p =the specific heat of the outgoing gases, in kcal·m⁻³·°C⁻¹. By substituting Eq. (3.4a) in Eq. (3.11a), we get

$$I_{L} = \varphi(\Delta \vartheta) \cdot c_{p} \cdot \Delta \vartheta \cdot A \cdot \sqrt{H}$$
(3.11b)

where

$$\varphi(\Delta \theta) = \frac{2}{3}\mu \sqrt{2g} \frac{\sqrt{\frac{\Delta \theta}{273}}}{1 + \frac{\Delta \theta}{273}} \cdot \left(\frac{H''}{H}\right)^{3/2}$$
(3.11c)

If the ratio H''/H is determined from Eq. (3.9), then Eq. (3.11c) yields $\varphi(\Delta\theta)$ for different values of a and $\Delta\theta$. For combustion of wood fuel, this relation is represented in Fig. 8, which is based on the values of μ , G_0 , and L given on p. 30. It is seen from this graph that the family of curves in question

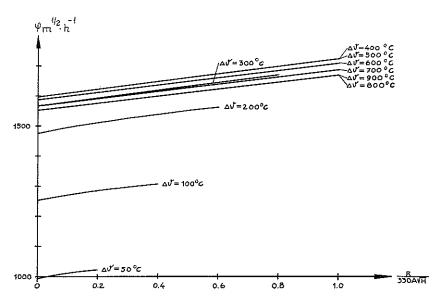


Fig. 8. Relation between the coefficient φ in the equation $I_L = \varphi \cdot c_p \cdot \Delta \vartheta \cdot A \cdot \sqrt{H}$ and the ratio $a = R/330 \cdot A \cdot \sqrt{H}$ for various values of the temperature difference $\Delta \vartheta$.

consists of approximately parallel straight lines. Therefore, we can write

$$\varphi = \varphi_0 + 120 \cdot a \tag{3.12}$$

where

 φ_0 is the value of φ for a=0.

Table 1 gives φ_0 for various values of $\Delta \theta$.

For $\Delta 9 > 300$ °C, the value of φ_0 is nearly independent of the temperature, and may be assumed to range from 1500 to 1600.

If the rate of combustion, and hence the factor a, are known, then I_L can be calculated for any instant, t, during the process of fire development. The assumptions which have been chosen for the deduction of the expression for the term I_L will be briefly discussed in what follows.

In spite of the fact that steady-state conditions have been assumed in the calculation of the flow through the opening, the results obtained from this calculation can also be applied under non-steady-state conditions, because a change in a position of equilibrium almost immediately gives rise to the establishment of a new position of equilibrium [7].

Table 1. Relation between φ_0 and $\Delta\vartheta$.

	Δϑ, °C	$g_0,$ $m^{1/2} \cdot h^{-1}$	Δϑ, °C	$\varphi_0,$ $m^{1/2} \cdot h^{-1}$	
	10	515	500	1597	
	50	991	600	1587	
	100	1254	700	1567	
	200	1476	800	1551	
	300	1567	900	1552	
•	400	1595	1000	1510	

The calculation of I_L requires that the specific heat of the outgoing gases shall be known. The air content of these gases is dependent on the rate of combustion. When a=0, the outgoing gases consist of air alone, and when a=1, they consist of gaseous products of combustion only. However, the difference in the specific heat between the air and the combustion gases is very slight, and the specific heat may therefore to a close approximation be regarded as independent of the rate of combustion.

Furthermore, in the calculation of φ , it is assumed that the values of G_0 and L remain constant during the whole process of fire development, irrespective of the rate of combustion. This assumption has not been verified by any physical considerations, but if we examine the right-hand member of Eq. (3.9), then we find that the effect produced by an error in the difference between G_0 and L on the value of φ is comparable to that of an error in a, that is to say, this effect is very slight.

In his treatment which relates to the flame phase only, Kawagoe has assumed that the temperature is uniform in the whole enclosed space. In the present publication, this assumption has been extended so as to be applicable during the whole process of fire development. The assumption that the variation in temperature in a vertical direction is relatively slight when the intensity of the fire decreases has been confirmed by the full-scale tests which are described in Chapter 5 of this publication. In most of these tests, the dispersion in the temperature measured at different points in the enclosed space during the cooling phase was found to be smaller than during the flame phase.

Modification of treatment in cases where enclosed spaces are provided with several openings which differ in height

The deduction carried out in the preceding section is applicable only in the cases where the air is supplied to the enclosed space through one or several openings which are equal in height, and which have a common neutral zone. If the enclosed space is provided with several openings which differ in height, then a corresponding deduction can be made in each individual case. For the determination of the maximum rate of combustion, R_{max} , Yokoi [13] has described an approximate method. It consists in the determination of a fictitious air flow factor, which is used in the original formula, Eq. (2.3). The fictitious air flow factor is determined from the expression

$$(A \cdot \sqrt{H})_{\text{fict}} = \sum_{i} A_{i} \cdot \sqrt{H}_{i}$$
 (3.13)

For some cases, Kawagoe [14] has compared the values of R_{max} which were obtained from Eq. (3.13) with those which were determined by means of accurate calculations. He found that Eq. (3.13) gives values which are sufficiently accurate for practical uses when the differences in the height and in the vertical position of the openings are not too great. The Swedish Building Regulations 1967 recommended another acceptable approximation, namely, that a weighted average of the heights of the individual openings should be used as a value of H in the calculation of the air flow factor.

After a fictitious value of the air flow factor has been determined from Eq. (3.13), this value can be used instead of $A \cdot \sqrt{H}$ in Eq. (3.11b) where the factor φ is calculated from Eq. (3.12) as before.

Modification of treatment in cases where enclosed spaces are provided with horizontal openings in roofs

In the preceding two sections, it was assumed that all openings in the enclosed space are vertical. In the present section, we shall expound a more general theory which makes it possible to take account of the presence of

horizontal openings in the roof of the enclosed space. We suppose the enclosed space to be in conformity with Fig. 9.

We assume that there exists a neutral zone in the enclosed space, and that the level of this zone is not higher than the upper edge of the vertical opening, and not lower than its lower edge. A condition which is prerequisite to this assumption will be stated further on in the present section. For the vertical opening, the rate of flow of the outgoing gases, Q_{out} , is obtained from Eq. (3.4a), and the rate of flow of the incoming air, Q_{in} , is computed from Eq. (3.4b). The velocity of the gases which flow out through the horizontal opening is (Bernoulli's theorem)

$$v_h = \sqrt{2gh' \cdot \frac{\rho_0 - \rho_g}{\rho_g}} \tag{3.14}$$

This formula has been verified experimentally in connection with studies of venting fires [7]. The rate flow of the outgoing gases through the horizontal opening is

$$Q_h = \mu \cdot A_h \cdot \nu_h \cdot \rho_g \tag{3.15}$$

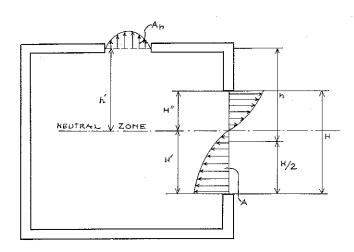


Fig. 9. Gas flow in an enclosed space provided with a vertical opening and a horizontal opening.

A = Area of the vertical opening.

 A_h = Area of the horizontal opening.

H = Height of the vertical opening.

H'' = Height of the vertical opening above the neutral zone level.

H' = Height of the vertical opening below the neutral zone level.

h = Vertical distance from the centre of the vertical opening to the level of the horizontal opening.

h' = Vertical distance from the neutral zone to the level of the horizontal opening.

At a maximum rate of combustion, $R = R_{\text{max}}$, the equation of mass balance of gases requires that

$$\left.\begin{array}{l}
Q_{\text{out}} + Q_h = G_0 \cdot R_{\text{max}} \cdot \rho_0 \\
Q_{\text{in}} = L \cdot R_{\text{max}} \cdot \rho_0
\end{array}\right\} \tag{3.16}$$

These two relations can be used to determine the position of the neutral zone which is modified in view of the presence of the horizontal opening, and then the maximum rate of combustion, R_{max} , can be calculated. R_{max} is a function of the term $\frac{A_h \cdot \sqrt{h'}}{A \cdot \sqrt{H}}$ at a given temperature. The value of R_{max} varies slightly with the temperature. If we write

$$R_{\text{max}} = 330 \cdot (A \cdot \sqrt{H})_{\text{fict}} \tag{3.17}$$

then $(A \cdot \sqrt{H})_{\text{fict}}$ can be determined from the alignment chart in Fig. 10, which is entered at the value of $A_h \cdot \sqrt{h}/A \cdot \sqrt{H}$.

It is to be expected that the value of $(A \cdot \sqrt{H})_{\text{fict}}$ determined in this manner may be used to an adequate degree of accuracy in the same way as the air flow factor, $A \cdot \sqrt{H}$, to characterize a fire. Accordingly, for an enclosed space with horizontal openings in the roof, the opening factor is given by the expression $(A \cdot \sqrt{H})_{\text{fict}}/A_t$, where $(A \cdot \sqrt{H})_{\text{fict}}$ is determined from the alignment chart in Fig. 10. The term I_L in the equation of heat balance, Eq. (2.1), of the process of fire development is obtained by analogy with Eq. (3.11b) from

$$I_{L} = \varphi(\Delta \theta) \cdot c_{p} \cdot \Delta \theta \cdot (A \cdot \sqrt{H})_{\text{fict}}$$
(3.18)

where $\varphi(\Delta \theta)$ is given by Eq. (3.12).

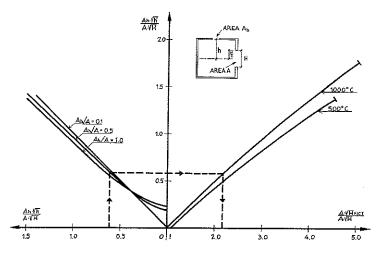


Fig. 10. Alignment chart for the calculation of the value of the modified air flow factor, $(A\sqrt{H})_{\text{fict}}$, on the basis of known geometrical data. For notations, see Fig. 9.

This procedure presupposes that the flow through the horizontal opening in the roof is not predominant. Consequently, the factor $A_h \cdot \sqrt{h'}/A \cdot \sqrt{H}$ has an upper limit at which the above model of the flow conditions ceases to be relevant. This upper limit is

$$\frac{A_h \cdot \sqrt{h'}}{A \cdot \sqrt{H}} = \begin{cases} 1.76 \text{ at } 1000^{\circ}\text{C} \\ 1.37 \text{ at } 500^{\circ}\text{C} \end{cases}$$

At this limit, the neutral zone is on a level with the upper edge of the vertical opening, and then h' is identical with the vertical distance from the level of the horizontal opening in the roof to the upper edge of the vertical opening. Tests [7] have indicated that the model used is relevant up to this upper limit. For values of $A_h \cdot \sqrt{h'}/A \cdot \sqrt{H}$ which are higher than this limit, all gaseous products of combustion will be vented through the horizontal opening in the roof. This is sometimes intentional if it is desired that the spread of fire to adjoining rooms should be prevented by venting the fire [15]. If the air and the combustion gases flow in the main through horizontal openings, then the flow becomes unstable, and it is difficult to represent it by a simple theoretical model [16].

Example showing how to use alignment chart in Fig. 10

Calculate the maximum rate of combustion, R_{max} , during the flame phase at 1000°C in the enclosed space characterized by the following data:

$$A = 2 \text{ m}^2$$
, $H = 1 \text{ m}$, $A_h = 1 \text{ m}^2$, $h = 1.5 \text{ m}$.
 $\frac{A_h \cdot \sqrt[4]{h}}{A \cdot \sqrt[4]{H}} = \frac{1 \cdot \sqrt[4]{1.5}}{2 \cdot \sqrt[4]{1}} = 0.61$; $\frac{A_h}{A} = 0.5$

The dash line in the alignment chart gives

$$(A \cdot \sqrt{H})_{\text{fict}} = 2.18 \cdot A \cdot \sqrt{H} = 4.36 \text{ m}^{5/2}$$

 $R_{\text{max}} = 330 \cdot 4.36 \text{ kg} \cdot \text{h}^{-1} = 1440 \text{ kg} \cdot \text{h}^{-1}$

This alignment chart can also be used for enclosed spaces where the horizontal opening in the roof is replaced by a ventilation duct. In such cases, the height h is replaced by the height of the gas column (the static head), with the reduction of the losses due to friction, which can be expressed in terms of the equivalent loss in static head. In an ordinary flat or office equipped with common ventilators made of non-combustible materials, the effect of the ventilation ducts is usually negligible.

If the enclosed space is ventilated through air inlets and outlets by means of a fan installation, then the corresponding fictitious air flow factor can be calculated in an analogous manner. If the quantity of gases exhausted per unit time is $Q_{\text{out,}}^{\nu}$ in kg·h⁻¹, and the quantity of air supplied per unit time is $Q_{\text{in,}}^{\nu}$ in kg·h⁻¹, then, for $R = R_{\text{max,}}$ we have the conditions

$$\left. \begin{array}{l}
Q_{\text{out}} + Q_{\text{out}}^{\nu} = G_0 \cdot R_{\text{max}} \cdot \rho_0 \\
Q_{\text{in}} + Q_{\text{in}}^{\nu} = L \cdot R_{\text{max}} \cdot \rho_0
\end{array} \right\}$$
(3.19)

from which R_{max} and $(A \cdot \sqrt{H})_{\text{fict}}$ can be determined by analogy with the above.

4. Description of programme for digital computer

The integration of the system of equations given by Eq. (2.6) was carried out by using the Runge-Kutta method in a modified form which has been suggested by Merson [22]. This modified method enables the computer to choose that interval of integration, Δt , which is required in order to ensure a certain definite degree of accuracy. To integrate the above-mentioned system of equations in the time interval from t to $t+\Delta t$, this system was evaluated in five individual operations. A determination of the temperature of the combustion gases is required for each one of these operations. According to Eqs. (2.1), (2.10), (2.11), and (3.11), the equation of heat balance of the process of fire development can be written

$$I_C = I_L + I_W + I_R \tag{4.1}$$

where

$$\begin{split} I_{L} &= \varphi(\Delta \vartheta) \cdot c_{p} \cdot A \sqrt{H}(\vartheta_{g} - \vartheta_{0}) \\ I_{W} &= \sum_{j} I_{W, j} = \sum_{j} A_{j} \psi_{1, j}(\vartheta_{g} - \vartheta_{1, j}) \\ I_{R} &= A \cdot (E_{g} - E_{0}) \end{split}$$

If I_R , φ and c_p are calculated on the basis of that value of the temperature of the combustion gases which has been obtained from the next preceding determination, then Eq. (4.1) can be solved for the temperature of the combustion gases, ϑ_g , in an explicit form

$$\vartheta_{g} = \frac{I_{C} + \varphi(\Delta \vartheta) \cdot A \sqrt{H} \cdot \vartheta_{0} + \sum_{j} A_{j} \cdot \psi_{1, j} \cdot \vartheta_{1, j} - I_{R}}{\varphi(\Delta \vartheta) \cdot c_{p} \cdot A \sqrt{H} + \sum_{j} A_{j} \psi_{1, j}}$$
(4.2)

which can be substituted in the system of equations represented by Eq. (2.6). The programme for the computer has been prepared in such a way as to be applicable to enclosed spaces bounded by structures which were assumed to be of up to three different types. Two of these structures were supposed to be homogeneous, whereas the third might be divided into two or three layers consisting of different materials, e.g. plasterboard panels, mineral wool, and brick.

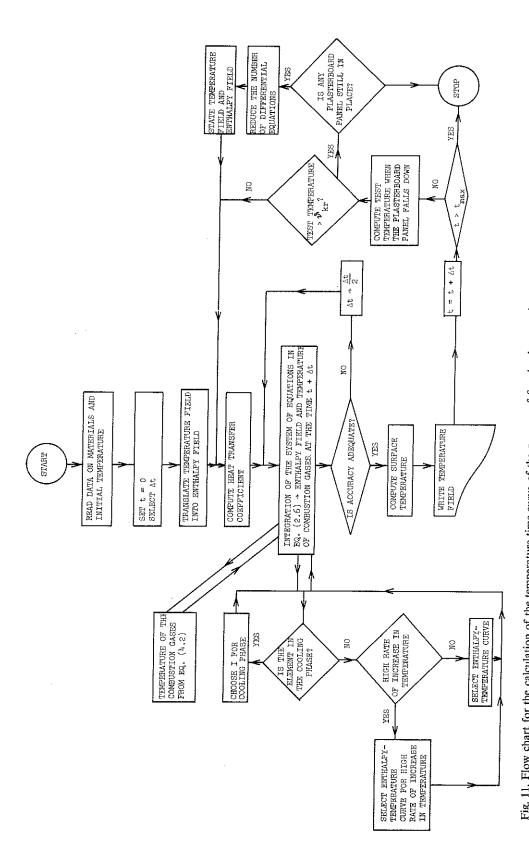


Fig. 11. Flow chart for the calculation of the temperature-time curve of the process of fire development in an enclosed space. The structures which bound the enclosed space are supposed to comprise, first, a roof or/and ceiling and a floor made of homogeneous material, and second, partitions which consist of a load-bearing frame made of steel studs and covered on the inside as well as on the outside with two plasterboard panels, 13 mm in thickness each.

The specific heat of most structural materials, e.g. concrete having a non-negligible moisture content or plasterboard panels, varies discontinuously with the temperature when these materials are subjected to physical or chemical transformations. Therefore the programme used the enthalpy I, in kcal·m⁻³, of the different materials as a dependent variable in the calculation of the temperature fields.

Fig. 11 represents a schematic flow chart which shows the programming procedure in the calculation of the combustion gases in an enclosed space where the floor and the roof or/and ceiling are made of concrete, while all the walls consist of a lightweight frame built of steel which are covered on the inside as well as on the outside with two plasterboard panels, 13 mm in thickness each. This type of wall exhibits two characteristic features, viz., first, experiments have shown that a plasterboard panel exposed to fire disintegrates when the temperature of the panel has reached a certain definite value, and second, the relation between the temperature and the enthalpy is dependent on the rate of temperature rise. Furthermore, it was necessary to choose different enthalpy-temperature curves depending on whether the temperature of the plasterboard panel in question was assumed to be increasing or decreasing. For further particulars, reference is made to the description of the calculations for the Type G enclosed space in Chapter 7.

5. Calculation of time graphs of rate of combustion for some full-scale tests described in literature

As has been mentioned in Chapter 2, the present chapter will deal with the comparative calculations which have led to a determination of the variation in the rate of combustion, expressed in kcal·h⁻¹, with the time. The method of successive approximations employed for this purpose consisted in making calculations which were based on different forms of the time graph of the rate of combustion. These calculations were repeated until they resulted in that curve which corresponded to the closest agreement between experimental and calculated curves representing the variation in the temperature of the combustion gases with the time. This method required a certain systematization of the description which represents the variation in the quantity of energy released by combustion with the time. What can be assumed to be known to a sufficient degree of accuracy in this connection is solely the total quantity of energy that can be liberated during the whole process of fire development, i.e. the fire load. Furthermore, it can be assumed that Eq. (2.3)

$$R_{\text{max}} = 330 A \sqrt{H}$$
 kg of wood per h

expresses the maximum rate of combustion, in kg of wood per h. In order that a theoretical determination of the temperature of the combustion gases may be possible, it is moreover necessary to determine the relation between released energy and weight loss of fuel, W, in

$$I_C = R \cdot W \quad \text{keal} \cdot h^{-1}$$

This determination is rendered difficult by the fact that the combustion of the gases formed by pyrolysis and that of solid wood fuel constitute a complicated process which involves a series of chemical reactions in different phases, see [4]. In respect of energy conditions, some of these reactions are endothermic, others exothermic. So far as the Authors know, no systematic investigation has been made up to now in order to carry out a quantitative analysis of the liberation of energy during the individual phases of the process of fire development.

These considerations have necessitated certain assumptions which concern the form of the curve showing the variation in the rate of combustion with the time. These assumptions are stated in what follows. The quantity of energy

liberated per unit time during the ignition phase was supposed to increase according to a polygonal function of the time to a level which corresponds to the rate of combustion during the flame phase. The determination of this level was based on Eq. (2.3), $R_{\text{max}} = 330 A \sqrt{H}$ kg of wood per h. When the quantity of energy released per kg of wood fuel during the flame phase was assumed to range from 2500 to 2800 kcal, it was found that the calculated temperature-time curves were closely in agreement with the results of the full-scale tests in respect of the maximum temperature and the duration of the flame phase. In order to adapt these assumptions to the temperature-time curves which have been published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures", and which are used in Sweden as a basis for calculations of the fire endurance of structural components in conformity with standard specifications, the quantity of heat evolved during the flame phase was supposed to be 2575 kcal · kg⁻¹ of wood, since this value had been employed for calculating the above-mentioned temperature-time curves. Cf. Chapter 2. The value in question has originally been stated by Kawagoe, who used it as a measure of the quantity of energy that is liberated by incomplete combustion of 1 kg of wood. By analysing the composition of the combustion gases in fire tests, see e.g. [5], it has been found that they contain considerable quantities of carbon monoxide, and this indicates that the combustion is not complete. It seems that the analysis of the combustion gases was performed during the flame phase, and the value 2575 therefore applies to this phase only. During the cooling phase, the weight loss of fuel per liberated energy unit is considerably less than during the flame phase. This means, that, when the whole process of fire is considered, the energy released by the combustion of 1 kg of wood must be higher than 2575 keal even if the combustion is incomplete. As a rule, there was scant basis for an accurate prediction of how the rate of completeness of combustion varied during the different phases. Consequently, in our comparative analysis the nominal heat value of wood, ranging from 3500 to 4500 kcal·kg⁻¹, was used as the energy liberated during the whole course of fire. This means that combustion is assumed to be complete throughout the present publication. The resulting time-temperature curves for the combustion gases therefore as a rule are found to be more in agreement with the maximum temperature curves than with the mean temperature curves obtained in the full scale tests.

Consequently, during the flame phase, the value of the rate of combustion will be constant, $330 \cdot A\sqrt{H} \cdot 2575$ kcal·h⁻¹. The cooling phase is characterized by a rate of combustion which decreases in conformity with a polygonal time graph in such a way that the slope of each individual side of the polygon is dependent on the duration of the flame phase. In every case, the area between the rate of combustion curve and the time axis must equal the fire load, expressed in energy units. The use of a polygonal time graph of the rate of

combustion makes it easier to check the agreement between the total quantity of energy liberated during the process of fire development and the fire load which is given from the outset.

The factors which must be taken into account in a comparative theoretical calculation of the temperature of the combustion gases in fire tests are enumerated in what follows.

- (1) Characteristics of fuel: Quantity, moisture content, porosity factor (hydraulic radius), distribution in the enclosed space.
- (2) Geometric characteristics of the enclosed space: Opening factor and its variation with the time (e.g. in the case where the fire burns through a door), shape of openings, cross-sectional area of ventilation ducts (if any).
- (3) Characteristics of the structures bounding the enclosed space: Structural design, thermal properties and temperatures of disintegration (if any) of the materials entering into the structures, emissivity characteristics of the surfaces.

By studying the available literature on full-scale wood fire tests, it was found, first, that the number of published tests is relatively small, and second, that those data which are so detailed as to render possible a comparative theoretical calculation have been stated only for a few tests out of this number. For this reason, it was necessary to confine the comparative calculations to only four test series reported in the literature or in other sources. These test series are enumerated below.

- (1) Test series A. Tests carried out by Sjölin, and dealt with in a thesis for a L. Techn. degree, [17].
 - (2) Test series B. Tests made by Kawagoe, and published in [5].
- (3) Test series C. Tests performed by Ödeen at the Royal Institute of Technology, Stockholm, and described in his doctoral thesis [18].
- (4) Test series D. Tests directed by *Pettersson* and *Ödeen*, which were carried out in a test house of the Atomic Energy Co., Ltd. (AB Atomenergi) at Studsvik, Sweden.

Test series A. Tests carried out by Sjölin. Calculation of time graphs of rate of combustion

The tests made by Sjölin were undertaken in order to study the spread of fire, and the process of fire development, in rooms and combinations of rooms exposed to ignition at a single point by heat radiation emitted by an explosion of nuclear weapons. The fire tests were carried out in a test house which was provided with concrete floor structures and concrete or lightweight concrete

¹ Not until too late in the publication of this paper did the authors learn about the full scale fire tests carried out at Fire Research Station, Boreham Hood, London [11, 12].

This is the more regrettable as these experiments make an excellent basis for comparative theoretical calculations.

walls. The test house was designed in such a way that various enclosed spaces might be formed by individual rooms, or by several rooms connected together. The available model scales were 1 to 1, 1 to 2, and 1 to 4. The variables recorded in these tests were expressed as functions of the time, and comprised the temperature and the velocity of gas flow at characteristic points, the intensity of radiation, the composition of the gases, and the rate of combustion, which was determined by continuous weighing of the quantity of fuel in the enclosed space.

The fire load in all these tests consisted of authentic furniture. This was an extraordinarily valuable feature of the tests, seeing that all the other full-scale tests which are dealt with in the present publication were made by using fire loads of the wood crib type.

Seven of the fire tests included in this test series were found to be suited for the present theoretical study. In the other tests, the ignition did not cause the fuel to take fire. Table 2 shows the scope of the seven tests under consideration.

Table 2. Test series A.

Test No.	Type of room	Window area, m ²	Opening factor, $\frac{A \cdot \sqrt[4]{H}}{A_t}$ $m^{1/2}$	Fire load, kg·m ⁻² of bounding surface area	Remarks
1	В	1.16	0.0237-0.06	3.5	(The fire burnt through a door, 1.6
2	В	1.16	0.0237-0.06	4.4	m^2 in area, during the time interval from $t=0$ min. to $t=6$ min.
3	L	1.16	0.0160-0.0356	4.9	(The fire burnt through a door, 1.6
4	L	1.88	0.0278-0.0486	5.6	m^2 in area, during the time interval from $t=8$ min. to $t=12$ min.
5	L	1.88	0.0548	5.0	•
6	L	2.95	0.068	5.7	
7	B+L	3.20	0.040	8.1	

B = Furnished two-person bedroom, 10.4 m² in floor area.

The curves representing the variation in the temperature of the combustion gases with the time in the above-mentioned seven tests, as well as the corresponding curves obtained by calculations with the help of automatic computer, are reproduced in Appendix 1. In the test No. 7, the enclosed space consisted of two contiguous rooms, which communicated through an open door. The partition between these two rooms was considered in the theoretical calculations to be an enclosed structure which possessed a heat-absorbing capacity.

L = Furnished living room, 18.8 m² in floor area.

B+L =Combination of B and L.

For this test, the curves representing the temperature of the combustion gases as a function of the time are shown separately for the bedroom and the living-room.

All the fires in these tests were characterized by a protracted process of ignition, which was followed by a rapid transition to the flame phase. For the calculated curves, the time was put equal to zero at the instant when the flame phase began, i.e. when the fuel took fire, in the actual fire tests. The time graphs of the rate of combustion which were finally obtained from the calculations are also reproduced in the respective diagrams. In the tests Nos. 1 to 4, the opening factor was changed during the process of fire development because the fire had burnt through a door. This change was taken into account in the calculations, and constituted the cause of the somewhat unusual shape of the time graphs of the rate of combustion which refer to these four tests. In these cases, the curves were plotted in such a way as to relate the rate of combustion to that value of the opening factor which was obtained after the fire had burnt through the door. In the test No. 3, the fact that the fire has burnt through a door is reflected very clearly in the curve, which shows that the temperature of the combustion gases remained constant at 500 to 600°C, and then rapidly rose to about 800°C when the fire burnt through the door.

Since the calculations were based on the opening factor and on the total energy content of the fuel, it was possible to choose the time graphs of the rate of combustion in such a manner that the agreement between the observed and calculated time graphs of the temperature of the combustion gases was very close in all the tests except the test No. 7. In the test No. 7, the calculated time graph of the combustion gas temperature was compared with the corresponding observed curves for the living-room as well as for the bedroom. The agreement between these curves was relatively close in the second case, but not in the first, where the curve is slightly displaced in time with reference to the curve for the bedroom. Moreover, when use was made of the opening factor which was determined geometrically, the temperatures obtained for the flame phase were found to be somewhat too low. However, it is not correct to regard the above-mentioned two rooms as a single enclosed space, since there existed quite a considerable difference in temperature between these rooms. Nor are the two rooms in question to be regarded as two separate enclosed spaces, since a certain heat exchange took place between these rooms.

Test series B. Tests made by Kawagoe [5]. Calculation of time graphs of rate of combustion

In [5], Kawagoe has described a large number of fire tests which had been carried out in Japan. In this investigation, he primarily studied the relation between the reduction in the weight of fuel per unit time and the dimensions

of the openings in the enclosed space. Among other things, he also deduced the equation which is reproduced in Eq. (2.3) in the present publication. The variables measured in these tests were the reduction in the weight of fuel, the temperature of the combustion gases, the composition of the combustion gases, the gas velocities, the intensity of radiation, and the pressure distribution in the window openings. On account of the above-mentioned main purpose of Kawagoe's investigation, it is only in three of these tests that the results of measurements, the geometric data, and the data on the materials entering into the structures which bounded the enclosed space are presented in such a way as to make it possible to carry out theoretical calculations of the type under consideration. These three tests were performed in a test house which was provided with walls made of hollow concrete blocks and with concrete floor and roof structures. The other test data are given in Table 3.

Table 3. Test series B.

	Fire load,	Bounding	Window	Opening factor,
•	kg of	surface	dimensions,	$A\cdot \sqrt{H}$
Test		area,	width \times height,	$\overline{A_t}$
No.	wood	m²	m×m	m1/2
1	400	48	0.93×1.8	0.0467
2	900	48	0.93×1.8	0.0467
3	1000	48	0.93×1.8	0.0467

Each hollow concrete block used for the walls comprised a single large cavity, without any subdivisions. The volume of the cavity was estimated at 30 to 40 per cent of the total volume of the block. In the calculations, the walls were considered to be composed of two different structures. One of them consisted of concrete alone, while the other comprised three layers, viz., concrete, air-filled cavity, and concrete, respectively. The second structure represented that part of the wall surface which corresponded to the cavities of the hollow concrete blocks, while the first structure was equivalent to the remaining part of the wall surface.

The observed and calculated temperature-time curves, as well as the time graphs of the rate of combustion used in the calculations, are reproduced in Appendix 2. The fire loads in the tests Nos. 2 and 3 were relatively high and the duration of the fires in these tests was therefore long. In order that the calculated values should agree with the values observed in these tests, it was necessary to choose a comparatively flat slope for the ascending branch of the time graph of the rate of combustion.

The data on the cooling phases in the tests Nos. 2 and 3 reported in [5] are not complete, and this is the reason why the curves relating to these tests break off at such an early stage.

Test series C. Tests published by Ödeen [18]. Calculation of time graphs of rate of combustion

The tests described by Ödeen in [18] were carried out in a tunnel building of an approximately semi-circular shape, which had been specially constructed for this purpose. It was provided with a concrete wall, 20 cm in thickness, its total bounding surface area was 75 m², and its total enclosed volume was 46 m³. A fan system made it possible to regulate and to measure the quantity of air which was supplied per unit time of the fire. In addition, a vent, 0.5 m² in cross-sectional area, for conveying the combustion gases to the outside air was provided in the upper part of each end wall of the tunnel.

A series of fire tests using fir wood as fuel has been carried out in this test building. A study was made of the effects produced on the process of fire development by the factors which are enumerated in what follows.

- (1) The volume of air supplied per unit time to the tunnel, Q, in $m^3 \cdot s^{-1}$.
- (2) The quantity of combustible material (fire load), M, in kg of wood.
- (3) The hydraulic radius of the fuel, r, in cm. This factor expresses the ratio of the total volume of the fuel to its total bounding surface area.

The scope of the test series using wood fuel is shown in Table 4, which was extracted from [18].

In order that the results of these tests may be compared by means of calculations with those of fire tests in ordinary enclosed spaces, where the rate

Table 4. Test series C.

Test No.	Fire load, M, kg of wood	Rate of air supply, Q , $m^3 \cdot s^{-1}$	Hydraulic radius, r, cm	Moisture content of fuel, per cent	Energy content of fuel, Mcal
1	270	1.0	_	9	1129
2	675	2.0	1.0	17	2565
3	675	1.0	1.0	17	2565
4	675	0.7	1.0	22	2468
5	675	1.5	1.0	22	2 468
6	675	MIN	1.0	21	2501
7	675	1.0	1.7	21	2501
8	675	0.7	1.7	21	2501
9	675	1.0	0.6	21	2501
10	270	0.7	1.0	21	1000
11	405	0.7	1.0	22	1481
12	405	0.7	2.4	28	1440
13	135	0.7	1.0	28	481
14	945	1.0	1.6	16	3659
15	1350	2.0	1.4	17	5130
16	405	0.7	0.4	17	1539

of air supply is determined by the openings in the enclosed space, the quantity of air supplied per unit time, Q, must be converted into an air flow factor, $A\sqrt{H}$, or into an opening factor, $A\sqrt{H}/A_t$. In this connection, it is necessary to take account of the air flow which may possibly enter into the tunnel through the outlets for combustion gases at low values of Q, and may therefore increase the value of the air flow factor.

In an ordinary enclosed space provided with a vertical opening, the rate of flow of the incoming air, Q_{in} , is given according to Eq. (3.4b), by the relation

$$Q_{\rm in} = 2/3 \ \mu B(H')^{3/2} \sqrt{2g \cdot \frac{(\rho_0 - \rho_g)}{\rho_0}} \ \text{m}^3 \cdot \text{s}^{-1} =$$

$$= 2/3 \cdot 0.7 \cdot A\sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left(\frac{\Delta \theta}{273} - \frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 3/3 \cdot 0.7 \cdot A\sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left(\frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 2/3 \cdot 0.7 \cdot A\sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left(\frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 2/3 \cdot 0.7 \cdot A\sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left(\frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 2/3 \cdot 0.7 \cdot A\sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left(\frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 2/3 \cdot 0.7 \cdot A\sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left(\frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 2/3 \cdot 0.7 \cdot A\sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left(\frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

The notations used in Eq. (5.1) have been explained in Chapter 3. With the help of Eqs. (3.9) and (3.10), we obtain, for a=1, the relation between H''/H and $\Delta \theta$ shown in Table 5.

Table 5. Relation between H''/H and $\Delta\vartheta$.

 Δϑ	H''/H	Δθ	H''/H	
400	0.61	800	0.64	
500	0.62	900	0.65	
 600	0.63	1000	0.66	

Since H' = H - H'', we can substitute H'/H in Eq. (5.1). For different values of the temperature $\vartheta_q = \vartheta_0 + \Delta \vartheta$, this yields the simplified expression

$$Q_{\rm in} = K \cdot A \sqrt{H} \tag{5.2a}$$

where K varies with the temperature of the combustion gases, θ_g , in conformity with Table 6, where θ_0 was put equal to zero.

Table 6. Relation between K and ϑ_q .

 ϑ_g	K	ϑ_g	K	
400	0.40	900	0.38	
500	0.40	1000	0.38	
600	0.40	1100	0.37	
700	0.39	1200	0.37	
800	0.39			

In the case under consideration, the rate of flow of the incoming air, $Q_{\rm in}$, was equal to the quantity of air supplied per unit time by the fan system. Hence it follows that this quantity of air can be replaced by a fictitious air flow factor, $(A\sqrt{H})_{\rm fict}$ by means of Eq. (5.2a).

If the constant K is put equal to its value at 600° C, then we obtain

$$Q = 0.40(A\sqrt{H})_{\text{first}} \tag{5.2b}$$

By substituting $A_t = 75 \text{ m}^2$, we get

$$(A\sqrt{H}/A_t)_{\text{fict}} = 0.033Q \tag{5.3}$$

This relation between the fictitious opening factor, $(A\sqrt{H}/A_t)_{\text{fict}}$, and the rate of air flow, Q, shall include the effect of the air which may possibly enter into the tunnel through the paraboliform outlets for combustion gases, 0.5 m² in cross-sectional area each. In the calculations, these outlets were assumed to be replaced by two rectangular openings having a base B=1.13 m and a height H=0.44 m each.

If the rate of flow of the incoming air through these openings is put equal to Q_{in} , then Q_{in} at a maximum rate of combustion is given by the equation

$$Q_{\rm in} + Q\rho_0 = LR_{\rm max}\rho_0$$

For notations in this equation, see Chapter 3. By substituting Q_{in} from Eq. (3.4b), we find

$$2/3 \sqrt{2gH' \frac{\rho_0 - \rho_g}{\rho_0}} \mu B \cdot H' \rho_0 + Q \rho_0 = L R_{\text{max}} \rho_0$$
 (5.4)

where R_{max} is determined from Eq. (3.5a).

$$Q_{\text{out}} = G_0 \cdot \rho_0 \cdot R_{\text{max}}$$

or

$$2/3 \cdot \sqrt{2gH'' \frac{\rho_0 - \rho_g}{\rho_g}} \cdot \mu \cdot B \cdot H'' \cdot \rho_g = G_0 \cdot \rho_0 \cdot R_{\text{max}}$$
 (5.5)

By substituting B=1.13 m, H=0.44 m, $\mu=0.7$, $G_0=4.86$ Nm³·kg⁻¹, L=3.98 Nm³·kg⁻¹ and R_{max} from Eq. (5.5) into Eq. (5.4) we obtain

$$0.83(H'/H)^{3/2} + 0.722Q = 0.38(H''/H)^{3/2}$$
(5.6)

In the calculation of ρ_g , the temperature of the combustion gases was assumed to be 600°C.

For those values of Q which are of interest in this connection, i.e. $Q \ge 0.7$ m³ · s⁻¹ in conformity with Table 4, Eq. (5.6) has no solution in the interval $0 \le H'/H \le 1$. This means that the flow is possible in an outward direction only.

In the test No. 6, the fans were switched off, and the exchange of air took place through the outlets for combustion gases alone. If these outlets are supposed, as before, to be approximately represented by rectangular openings having a base of 1.13 m and a height of 0.44 m each, then we obtain an opening factor $A\sqrt{H}/A_t = 0.0088$ m^{1/2}, which corresponds to a rate of air flow $Q_{\min} = 0.26$ m³ · s⁻¹ according to Eq. (5.3). Cf. the average value of the rate of air flow, 0.19 m³ · s⁻¹, which has been computed by Ödeen for the whole process of fire development.

Thus, in the tests under consideration, the fictitious opening factor, $(A\sqrt{H}/A_t)_{\text{fict}}$, and hence the maximum rate of combustion, are determined approximately in conformity with the above from Eq. (5.3).

For the 16 tests comprised in this series, Table 7 gives, first, the opening factor calculated by means of the above relations, and second, the opening factor, $(A \cdot \sqrt{H}/A_t)_{\rm exp}$, that has proved to give theoretical results which are in agreement with the experimental values. The test results, the theoretical results, and the time graphs of the rate of combustion are shown in Appendix 3. As regards the time graphs of the combustion gas temperature, it is to be noted that the six full-line curves represent the temperatures at different

Table 7. Theoretical and experimental values of the opening factor.

Test series C [18].

Test No.	Fire load, M, kg of wood	Hydraulic radius, r,	Rate of air supply, Q, m³·s-1	Opening factor, theoretical value, $(A\sqrt{H}/A_t)_{\text{theor}}$ calculated from the formula $A\sqrt{H}/A_t = 0.0334 \cdot Q$	Opening factor, experimental value, $(A\sqrt{H}/A_t)_{\rm exp}$	$\frac{(A\sqrt{H}/A_t)_{\rm exp}}{(A\sqrt{H}/A_t)_{\rm theor}}$
		1.0	0.7	0.003	0.035	1.50
4	675	1.0	0.7	0.023	0.035	1.52
8	675	1.7	0.7	0.023	0.015	0.65
10	270	1.0	0.7	0.023	0.015	0.65
11	405	1.0	0.7	0.023	0.020	0.87
12	405	2.4	0.7	0.023	0.012	0.52
13	135	1.0	0.7	0.023	0.005	0.22
16	405	0.4	0.7	0.023	0.060	2.60
1	270	_	1.0	0.033	0.023	0.70
3	675	1.0	1.0	0.033	0.043	1.30
7	675	1.7	1.0	0.033	0.037	1.12
9	675	0.6	1.0	0.033	0.051	1.54
14	945	1.6	1.0	0.033	0.037	1.12
5	675	1.0	1.5	0.050	0.055	1.10
2	675	1.0	2.0	0.067	0.060	0.90
15	1350	1.4	2.0	0.067	0.060	0.90
			about			
6	675	1.0	0.25	0.009	0.010	1,11

points in the enclosed space. The fine dash-line curve summarises the values recorded in the radiation measurements, and the heavy dash-line curve is the calculated curve. The variation in the rate of combustion with the time, I_c , is represented in terms of $330 \cdot (A\sqrt{H})_{exp} \cdot 2575 \text{ kcal} \cdot h^{-1}$ put equal to unity. The rate of flow of the incoming air, Q, conveyed by the fan system was constant during the whole process of fire development. The radiation measurements have provided certain indications for choosing the instant at which the rate of combustion had decreased to zero. The difference between the temperature at the level of the floor surface and the average temperature in the other parts of the enclosed space has been taken into account. In the test series under review, this difference in temperature has probably been increased owing to the fact that air was supplied to the enclosed space by means of fans at the floor surface level. The temperature difference was taken into consideration by assuming that the coefficient of heat transfer at the floor surface was equal to 80 per cent of the corresponding coefficient for the other surfaces. In all cases when the rate of burning was controlled by the fuel bed and not by the ventilation it was taken into account that heat energy was withdrawn from the enclosed space by that part of the incoming air which did not take part in the combustion.

As may be seen from Table 7, the positive as well as negative differences between the value of the opening factor, $(A\sqrt{H}/A_t)_{\rm exp}$, determined from the test results, i.e. the actual maximum values of the rate of combustion, and the corresponding values obtained on the assumption that the rate of combustion is limited by the rate of air supply, were found to be great in some tests. However, the calculated curves were as a rule closely in agreement with the observed values. The agreement between the maximum temperature and the duration of the flame phase indicates that those values of the quantity of energy released per unit time which were used for the theoretical calculations were on the whole correct.

Furthermore, Table 7 shows two other factors among those which, in addition to the air flow factor, determine the rate of combustion. These factors are the hydraulic radius and the amount of fuel (the fire load). The effect of the first-mentioned factor can be demonstrated, for instance, by a comparison between the tests. Nos 16 and 12. The values of the air flow factor, as well as those of the fire load, in these two tests were equal, whereas the respective values of the hydraulic radius were 0.4 and 2.4 cm. In consequence of the difference in the hydraulic radius, the actual maximum rate of combustion in the test No. 16 was about 5 times as high as in the test No. 12. This may roughly be explained by the simplified study of the mechanism of combustion in what follows.

In a wood fuel consisting of comparatively large pieces of wood, pyrolysis takes place in several forms at the same time. In a certain definite inner zone

of the wood, where the temperature is relatively low, say, below about 250° C, the reactions are endothermic, whereas in the outer zones, where the temperature is higher, the reactions are exothermic, and in certain cases, e.g. in the secondary pyrolysis of tar products, markedly exothermic. The combustion of the products of pyrolysis generates heat, which increases the temperature of the fuel by conduction, and hence renders possible an exothermic decomposition in the inner zone. In the test No. 16, the wood fuel consisted of concrete form timber, for which the ratio of the volume to the exposed surface was 0.4 cm. Accordingly, if the width, the length, and the thickness of a piece of wood are denoted by b, l, and t, respectively, then we have

$$r = 0.4 \text{ cm} = \frac{b \cdot l \cdot t}{2l(b+t)}$$

If we put b=t (square cross section), then we obtain a thickness of 1.6 cm, and if we set $b \gg t$, then we get a thickness of 0.8 cm. In view of the small thickness, in combination with the mechanism of heat return to the fuel described in the above, it is probable that a few minutes after the fuel has taken fire the whole quantity of fuel is in a state of active exothermic pyrolysis. Tests [23] have shown that the progression of the charred layer on a wooden beam exposed to fire is about 0.6 mm · min⁻¹. This value is applicable to a firwood beam exposed to fire in conformity with a standard temperature-time curve. The variation in the rate of carbonization with the intensity of the process of fire development is a problem which appears to be wholly unexplored at the present time. But if the above-mentioned value, 0.6 mm · min⁻¹, is assumed to be correct, then this implies that the whole amount of fuel would be charred in the course of 5 to 10 min. Since from one half to two thirds of the total quantity of energy is liberated during the flame phase, this can explain the intense release of energy immediately after the fuel has taken fire.

The test No. 12 shows that when the fire-exposed surface area diminishes below a certain definite limit, the quantity of energy which can be developed per unit time is determined by the rate of progression of the charred layer, and not by the air flow factor. In this test, the hydraulic radius was 2.4 cm. If the above-mentioned value, $0.6 \text{ mm} \cdot \text{min}^{-1}$, which is probably too high in view of the low temperature during the process of fire development in this case, is used as a measure of the progression of the charred layer, then this corresponds to a maximum rate of combustion of $405 \cdot 0.6 \cdot 10^{-2}/0.24 = 10.1 \text{ kg} \cdot \text{min}^{-1}$. The value of the maximum rate of combustion computed from the formula $R = 5.5 A \sqrt{H}$ is 9.7 kg·min⁻¹, and this implies that the quantity of energy released per unit time during a fire is determined by the surface area exposed to fire, and not by the rate of air supply, at least during certain phases of the fire.

Furthermore, a closer study of Table 7 also demonstrates the marked effect produced by the amount of fuel on the rate of combustion. This can likewise be explained by means of the mechanism of return of the heat evolved by combustion which has been outlined in the above. In the tests Nos. 13, 10, 11, and 4, the values of the hydraulic radius, r, were equal, just as those of the rate of air supply Q, whereas the fire load, M, was varied. For the values of the fire load M=135, 270, 405, and 675 kg, which approximately corresponded to 7, 15, 20, and 35 Mcal· m^{-2} of bounding surface area of the enclosed space, the respective values of the maximum rate of combustion calculated from the test results were found to be equal to 0.22, 0.65, 0.87, and 1.5 times the rate of combustion which was computed on the basis of the air flow factor. In [18], it is stated that the flame phase was slightly developed in the tests Nos. 10 and 11, whereas the fuel did not take fire at all in the test No. 13.

Moreover, it is seen from Table 7 that the variation in the rate of combustion with the hydraulic radius seems to decrease as the opening factor or the fire load increases.

Finally, it may be useful to touch on the question to what extent these analytic fire tests, which were carried out under conditions that were idealized so far as possible, and which exhibited characteristics of combustion that in several tests markedly differed from those predicted by the theory, can be utilized as a basis for predicting the behaviour of more conventional fires. An ordinary fire in an enclosed space is governed to a varying degree by two feed-back mechanisms. In the first place, the lower density of the combustion gases forces unconsumed air by natural convection towards the flames, and hence increases the rate of combustion, as well as the evolution of combustion gases. In the second place, part of the heat generated by combustion returns to the fuel, and increases the rate of energy release. So long as these mechanisms are negative, the combustion remains stable. Cf. e.g. the test No. 13. An essential difference between Ödeen's tests described in [18] and an ordinary fire is that the return mechanism which governs the rate of air supply was eliminated in the tests. The effect of this circumstance is difficult to determine, but it may be mentioned for comparison that all the theoretical and experimental results in the test series (A, B, and D) where the exchange of air was self-regulated were found to be closely in agreement if the rate of combustion was assumed to be determined by the rate of air supply. However, a comparison of the time graphs of the rate of combustion for the test series A to D shows that, in the cases where differences were present, the results obtained from the test series C deviated from those of the other test series only in respect of the maximum quantity of energy released per unit time. If the time graphs of the rate of combustion are represented in terms of 330 \cdot $(A)\overline{H})_{\text{exp}}$ \cdot 2575 kcal · h-1 put equal to unity, then the curves for this test series are found

to be closely in agreement with those for the test series A, B, and D when the fire load and the opening factor are given. All the same, since a dispersion in the values of the quantity of energy liberated per unit time has been observed in the test series C even when the tests were identical in respect of the rate of air supply and the wall material, it should be noted that this dispersion indicates the need for determining the factors which, in addition to those mentioned in the above, govern the rate of release of energy.

Test series D. Tests made at the National Swedish Institute for Materials Testing. Calculation of time graphs of rate of combustion

Under the direction of the Fire Engineering Laboratory of the National Swedish Institute for Materials Testing, a test house for model-scale and full-scale fire studies has been erected at the Studsvik Test Station of the Atomic Energy Co., Ltd. The primary object of the investigations carried out in this house was to study the spread of fire and smoke along the exterior walls and along the ventilation ducts in the case of fire in an individual enclosed space in a multi-storeyed building. Extensive measurements of the reduction in the weight of fuel, the temperature, the intensity of radiation, the gas flow, and the composition of combustion gases were made in these tests, and the test results are therefore well suited for theoretical comparisons.

Fig. 12 shows the test house, which was three storeys high, and which consisted of a load-bearing steel frame clad with lightweight concrete elements. The results of the tests carried out up to now have not yet been published, but the test programme and the test equipment are described in [20] and [21]. The Authors of the present publication were afforded an opportunity to acquaint themselves with the results of the first four full-scale tests. The data for these tests are reproduced in Table 8.

The fires were initiated in the lowermost storey, which was connected with the outside air by means of a vertical ventilation duct by-passing the storeys

Table 8. Test series D.

Test No.	Fire load, kg of wood	Window dimensions, width × height m×m	Bounding surface area, m ²	Air flow factor, $A \cdot \sqrt{H}$, $m^{5/2}$	Air flow factor, fictitious value, $(A \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \$	Moisture content of fuel, per cent	Moisture content of lightweight concrete, per cent
1	350	1.3×1.3	75	1.93	2.10	7.7	4.0
2	200	1×1	75	1.00	1.10	8.8	3.1
3	115	0.8×0.8	75	0.57	0.69	7.4	2.1
4	1150	2×2	75	5.63	5.65	9.7	2.3

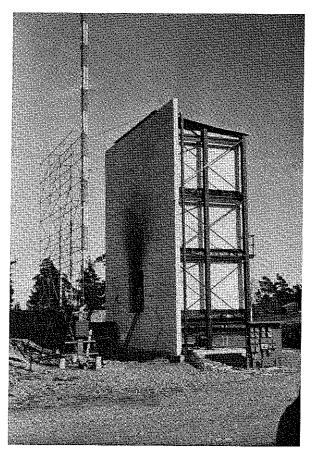


Fig. 12. Photograph of the test house used in the test series D to study the spread of fire and smoke along the external walls and through a vertical ventilation duct in the case of fire in a certain definite storey in a multi-storey building.

situated above. Calculations showed that the effect of the ventilation duct on the air flow factor was not to be disregarded, and a fictitious air flow factor, $(A\sqrt{H})_{\text{fict}}$, was therefore computed in accordance with the principles stated in Chapter 3. The magnitude of the corresponding correction can be seen from Table 8. The time graphs of the temperature of the combustion gases which were obtained from these tests, and which are expressed in terms of the mean value of the temperatures observed during the tests at 21 points in the enclosed space, are reproduced in Appendix 4. In order to give an idea of the dispersion about this mean value, the diagram relating to the test No. 2 shows the temperature-time curves for a point situated 45 cm above the floor and a point located 45 cm below the ceiling. Furthermore, the four diagrams in Appendix 4 also comprise the calculated curves which represent the varia-

tion in the temperature of the combustion gases with the time, as well as the time graphs of the rate of combustion in a dimensionless form which were used in the calculations.

It proved possible to bring the calculated and observed curves into close agreement by choosing the time graphs of the rate of combustion which were similar in shape in all the tests, and were based on a maximum rate of combustion $I_c = 330 \cdot A\sqrt{H} \cdot 2575$ kcal·h⁻¹. The value of the air flow factor used in the calculations was the value $(A\sqrt{H})_{fict}$, which was corrected so as to take account of the effect of the ventilation duct.

Summary

All the comparative calculations dealt with in the present chapter were based on the assumption that the energy conditions during the process of combustion can be characterized by an ignition phase in which the quantity of energy released per unit time increases from a zero value in accordance with a polygonal function of the time to a value that is given by the air flow factor. This phase is followed by a flame phase, during which the rate of combustion was supposed to be constant. After that the rate of combustion decreases to zero as a polygonal function of the time in the course of the cooling phase, during which the slopes of the individual sides of the polygon vary in a marked manner with the fire load. The higher the fire load, the slower the decrease in the rate of combustion. Of course, these assumptions give a simplified picture of the variation in the liberation of energy per unit time during the process of fire development. Thus, for most types of fire loads, it is to be expected that the plane part of the curve which represents the rate of combustion during the flame phase is rather to be regarded as the mean value of the quantity of energy released per unit time. This is illustrated in Fig. 13, which represents the variations in the observed rate of combustion (expressed in terms of the rate of reduction in the weight of fuel) and in the oxygen content of the combustion gases during the test No. 1 in the test series A. On account of technical difficulties in measurements, the values of the rate of reduction in the weight of fuel were somewhat uncertain immediately after the fuel had taken fire, as the rate of combustion was then liable to very wide variations. However, these values were confirmed by the fact that the oxygen content of the combustion gases in the enclosed space exhibited corresponding variations. Even when the curve which represents the variation in the rate of reduction in the weight of fuel with the time is known, our present knowledge of the relation between the rate of reduction in the weight of fuel and the rate of release of energy during the different phases of the process of fire development does not make it possible to determine the quantity of energy liberated per unit time during combustion. For the test No. 3 in the test series D, Fig. 14 shows three theoretical temperature-time curves calculated on the basis

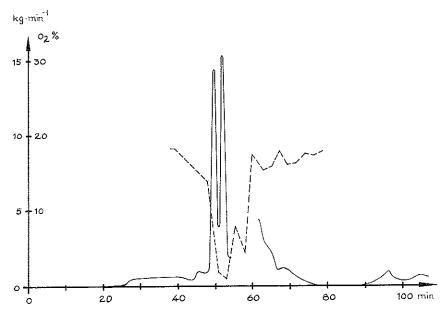


Fig. 13. Variation in the rate of combustion, in kg·min⁻¹, with the time, t, determined by measuring the reduction in the weight of fuel in one of the tests comprised in the test series A, see Chapter 5. Furthermore, this figure also shows the time graph of the oxygen content of the combustion gases in the same test (dash-line curve). During the interval from the 55th to the 60th minute, the weighing of the fuel was disturbed by the fact that parts of the ceiling of the enclosed space fell down.

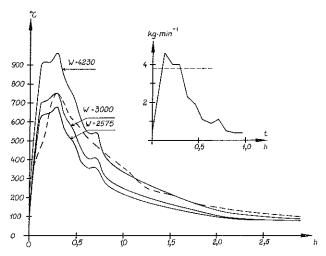


Fig. 14. Time graphs of the temperature of the combustion gases calculated for three different heat values of the fuel, viz., W=2575, 3000, and 4230 kcal·kg⁻¹ of wood, on the basis of an observed time graph of weight loss of fuel (inset). An experimental temperature-time curve is represented by the dash-line curve.

of a measured rate of weight loss. It was assumed that the heat value W corresponding to 1 kg of weight loss was constant during the whole process of fire development and equal to 2575,3000 and 4230 kcal, respectively. The dash line curve represents the values observed in the test. The observed rate of reduction in the weight of fuel is represented in a separate diagram in fig. 14. As seen from this diagram, the maximum rate of reduction in the weight of fuel is closely in agreement with the theoretical value $R = 5.5 \times A\sqrt{H} = 3.80 \text{ kg} \cdot \text{min}^{-1}$ (horisontal dash line). It is also seen from fig. 14 that the assumption of a constant value of W obviously is incorrect. A comparison between the curves indicates that the quantity of energy released per unit weight of fuel, at least during the first part of the cooling phase is greater that that during the flame phase.

Accordingly, the relative values of the rate of combustion obtained from the above comparative calculations cannot be directly expressed in terms of the rate of reduction in the weight of fuel, in kg·min⁻¹. If the quantity of energy liberated per unit time is expressed directly in kcal·min⁻¹, then this obviates the difficulty of determining the heat of combustion of the wood fuel during the various phases of the process of fire development.

A prerequisite in most of the calculations has been that the fire process is controlled by ventilation. This condition is far from being generally realized when it comes to actual fires. As a rule, a combination of small ventilation area and high fire load gives a process of fire development where the rate of burning is proportional to the air flow factor. If the fire load is low and the ventilation area large, the combustion will proceed as if in the open. This means that the rate of burning depends on the fire load density (fire exposed surface) and that an increase in ventilation will not result in a corresponding increase in rate of burning.

It is, however, impossible to say in advance if the process of fire development will be ventilation controlled or not even if the air flow factor and the fire load are known. The orientation and the distribution of the fuel in the enclosed space and the thickness or the porosity of the fuel will be a decisive factor in each particular case. An assumption that the rate of burning is determined by the air flow factor ought to give time-temperature curves which are on the safe side in practically every case. If such an assumption is made and the combustion in spite of this happens as in the open, i.e. is fuel bed controlled, the result will be a fire process of lower maximum temperature and, at least in some cases, of longer duration. The longer duration will not increase the severity of the fire to a corresponding degree. This is due to the fact that part of energy released by the combustion will be withdrawn from the enclosed space by the surplus air. In this way the temperature of the combustion gases will be lower compared to the case when the duration is the same but the process controlled by the ventilation.

To sum up, a comparative theoretical analysis of the results obtained from some thirty full-scale fire tests of the wood fuel type has been carried out in Chapter 5. The calculations made for this purpose covered relatively wide variations in fire load, opening factor, and hydraulic radius, as well as in the thermal properties of the structures bounding the enclosed space. As a result of these calculations, the time graph of the quantity of energy released per unit time may be assumed to be known within this range of variation.

6. Determination of general time graphs of quantity of energy released per unit time during different phases of process of fire development

In order to afford a basis for the calculation of the curve which represents the variation in the temperature of the combustion gases with the time during the process of fire development under varying conditions, it is necessary to systematize the time graphs of the rate of combustion which have been obtained in Chapter 5. A detailed investigation has been made of these graphs in order to find out how they vary with the fire load and with the opening factor. This investigation indicated the possibility of the simplification outlined in what follows.

If the ratio of the fire load, which is given from the outset, to the air flow factor, $A\sqrt{H}$, is constant, that is to say, if the duration of the fire is constant, then the time graph of the rate of combustion, expressed in a relative form in terms of the maximum rate of combustion, $330 \cdot A\sqrt{H} \cdot 2575$ kcal·h⁻¹, put equal to unity, is independent of the opening factor. This implies, for instance, that an enclosed space where the opening factor is $A\sqrt{H}/A_t = 0.01$ m^{1/2} and the fire load is q = 5 Mcal·m⁻² can be characterized by the same graph of the rate of combustion as an enclosed space where the opening factor is 0.04 m^{1/2} and the fire load is 20 Mcal·m⁻². Accordingly, the results of the calculations in Chapter 5, which show how the quantity of energy released per unit time varies with the time, can be represented by a graph which comprises a separate curve for each value of the ratio $qA_t/A\sqrt{H}$. In this connection, it is convenient to introduce the duration of the fire, T, defined as the duration of the flame phase, cf. Eq. (1.2),

$$T = qA_t/(1500A \cdot \sqrt{H})$$
 h

as the variable at which the graph shall be entered. In this formula for calculating the duration of the flame phase, the product of the constant 330 in the expression $R=330 \cdot A\sqrt{H} \text{ kg} \cdot \text{h}^{-1}$ and the heat value of the wood fuel, i.e. 4.5 Mcal·kg⁻¹, has been put equal to 1500. For the values of the duration of the fire defined in this way, T=0.1, 0.2, 0.3, 0.5, 0.75, 1.0, 1.5 and 2.0 h, which correspond to the respective fire loads (150, 300, 450, 750, 1125, 1500, 2250, and 3000) $\cdot A\sqrt{H}/A_t$ Mcal·m⁻² of the total surface area bounding

the enclosed space, Fig. 15 shows the variation in the rate of combustion with the time.

In order to make it easier to check the agreement between the total quantity of energy liberated during the process of fire development and the fire load, which is given from the beginning, and since a more accurate representation would be illusory, considering the character of the available data, the curve form has been assumed to be polygonal, just as in the comparative calculations in Chapter 5. In Fig. 15, $330A\sqrt{H} \cdot 2575$ kcal·h⁻¹ has been put equal to unity. The respective areas between the above-mentioned curves and the axis of time shall therefore be $\frac{1}{330 \times 2575} (150, 300, 450, 750, 1125, 1500, 2250, and 3000)$ area units. In the relevant Swedish regulations, the quantity T, determined by the relation

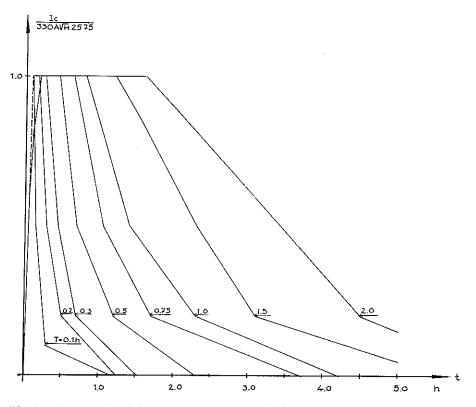


Fig. 15. Time graphs of the energy released per unit time in the process of combustion, I_c , expressed in a relative form by putting $330 \cdot 2575 \cdot A \sqrt{H}$ equal to unity. The eight curves shown in this figure correspond to different values of the duration of the fire defined as the duration of the flame phase by the expression $T = q \cdot A_t / (1500 \cdot A \cdot \sqrt{H})$. The dash-line portion of the curve for the ignition phase belongs to the curves relating to the lowest four values of the duration of the fire.

$T = qA_t/(25A\sqrt{H})$ min

designates the instant which marks the end of the flame phase and the beginning of the linear cooling phase. On the basis of the experiences derived from the comparative theoretical analyses, the instant at which the rate of combustion begins to decrease has been chosen so as to be slightly anterior to the instant defined by T. For the values of the duration of the fire T=6, 12, 18, and 30 min, the time graphs of the rate of combustion during the ignition phase have been given a slightly different shape, which implies that the fuel takes fire within a shorter period of time.

7. Calculation of time graphs of temperature of combustion gases for characteristic types of enclosed spaces varying in opening factor and in fire load

The time graphs of the rate of combustion represented in a relative form in Fig. 15 for fires of the wood fuel type in enclosed spaces are utilized in the present chapter as a basis for the calculation of complete time graphs of the temperature of the combustion gases. This is done for varying values of the opening factor and the fire load in enclosed spaces of the seven types dealt with in what follows, which differ in respect of the bounding structures.

Type A enclosed space

Bounding structures.

All the surfaces which bound the enclosed space are supposed to consist of a material, 20 cm in thickness, whose thermal properties are characterized by the average values given below, which apply to structural materials of such types as concrete, brick, and lightweight concrete.

Thermal conductivity, $\lambda = 0.7 \text{ kcal} \cdot \text{m}^{-1} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$.

Product of the specific heat and the weight per unit volume,

$$c \cdot \gamma = 400 \text{ kcal} \cdot \text{m}^{-3} \cdot {}^{\circ}\text{C}^{-1}$$
.

The same data on the properties of materials had also been used for the calculation of those temperature-time curves for the flame phase of the process of fire development which have been published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures".

Type B enclosed space

Bounding structures.

Concrete, 20 cm in thickness.

Thermal conductivity, $\lambda = 1.4 \cdot e^{-0.001} \cdot \theta \text{ kcal} \cdot \text{m}^{-1} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$ [1]. Enthalpy, *I*, see Fig. 17.

Type C enclosed space

Bounding structures.

Lightweight concrete, 20 cm in thickness. Weight per unit volume, $\gamma = 500 \text{ kg} \cdot \text{m}^{-3}$.

Thermal conductivity, λ , see Fig. 16.

Enthalpy, I, see Fig. 17.

The specific heat and the weight per unit volume of the lightweight concrete are assumed to be independent of the temperature. Consequently, the enthalpy-temperature curve is rectilinear. The variation in the thermal conductivity, λ , with the temperature is based on a determination which has been made in connection with the test series D described in Chapter 5.

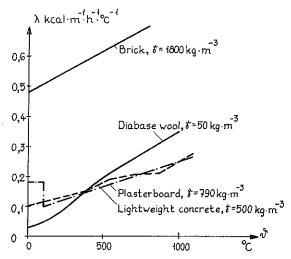


Fig. 16. Relations between the thermal conductivity, λ , and the temperature, ϑ , used in the calculations for brick, diabase wool, plasterboard, and lightweight concrete.

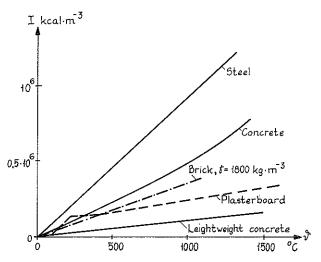


Fig. 17. Relations between the enthalpy, I, and the temperature, ϑ , used in the calculations for steel, concrete, brick, plasterboard, and lightweight concrete.

Type D enclosed space

Bounding structures.

Concrete, 50 per cent of the total bounding surface area.

Lightweight concrete, 50 per cent of the total bounding surface area.

Thicknesses, weight per unit volume, and thermal properties as in the Type B and Type C enclosed spaces, respectively.

Type E enclosed space

Bounding structures.

Lightweight concrete, 50 per cent of the total bounding surface area. Thickness, weight per unit volume, and thermal properties as in the Type C enclosed space.

Concrete, 33 per cent of the total bounding surface area.

Thickness and thermal properties as in the Type B enclosed space. Other structural components, 17 per cent of the total bounding surface area, enumerated in the order from the interior to the exterior:

Plasterboard panel, 13 mm in thickness.

Weight per unit volume, $y = 790 \text{ kg} \cdot \text{m}^{-3}$.

Diabase wool, 10 cm in thickness.

Weight per unit volume, $\gamma = 50 \text{ kg} \cdot \text{m}^{-3}$.

Brickwork, 20 cm in thickness.

Weight per unit volume, $\gamma = 1800 \text{ kg} \cdot \text{m}^{-3}$.

Thermal properties of plasterboard, diabase wool, and brick, see Figs. 16, 17, and 18.

The enthalpy-temperature curve chosen for brick is based on a value of the specific heat which is supposed to be independent of the temperature. In reality, the specific heat of brick slightly varies with the temperature [1]. However, in the present case, the values of the temperature rise in the brickwork, which is most remote from the surface exposed to fire, are so low that

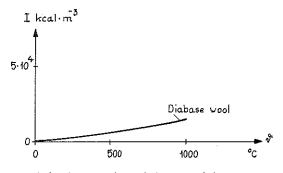


Fig. 18. Relation between the enthalpy, I, and the temperature, ϑ , used in the calculations for diabase wool.

the effect of this variation may be disregarded. The variation in the enthalpy of diabase wool with the temperature is based on that relation between the specific heat and the temperature which was published in [1]. The thermal properties of plasterboard are based on the curves which were published in [19]. It is assumed that the plasterboard panels will not fall down or disintegrate.

The values of the thermal conductivity, λ , of brick and diabase wool were taken from [1].

Type F enclosed space

Bounding structures.

Sheet steel, 2 mm in thickness, 80 per cent of the total bounding surface area.

Concrete, 20 cm in thickness, 20 per cent of the total bounding surface area.

Thermal properties as in the Type B enclosed space.

Curves representing the variations in the enthalpy and in the thermal conductivity of sheet steel with the temperature, see Figs. 17 and 19, respectively.

This type of enclosed space corresponds to a storage space, or the like with a sheet steel roof, sheet steel walls, and a concrete floor.

Type G enclosed space

Bounding structures.

Concrete, 20 per cent of the total bounding surface area.

Thickness and thermal properties as in the Type B enclosed space.

Other structural components, 80 per cent of the total bounding surface area, enumerated in the order from the interior to the exterior:

Two plasterboard panels, 2×13 mm in thickness.

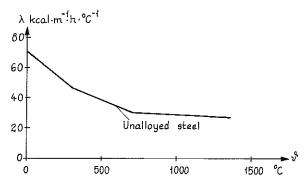


Fig. 19. Relation between the thermal conductivity, λ , and the temperature, ϑ , used in the calculations for unalloyed steel.

Weight per unit volume, $\gamma = 790 \text{ kg} \cdot \text{m}^{-3}$. Cavity, 10 cm in width.

Two plasterboard panels, 2×13 mm in thickness.

Thermal properties of plasterboard, see Figs. 16 and 20.

This structure represents a type of partition which is becoming more and more common, and which consists of two plasterboard panels on each side, supported on steel stud framing. It is assumed that the steel studs have no thermal conductivity and no thermal absorptivity.

The test results published in [19] have shown that plasterboard panels which are not fibre-filled disintegrate when their temperature on the side that is not exposed to fire reaches about 550°C. However, this does not apply to the outermost, i.e. the fourth, plasterboard panel. This panel is in contact with the air, which has a temperature of 20°C. Therefore, this panel never reaches a surface temperature of 550°C. In fact, tests have demonstrated that a plasterboard panel in this position disintegrates when the temperature at its centre rise to about 750°C. These criteria have been used in calculating the time graph of the combustion gas temperature for enclosed spaces of the type in question. The calculations were discontinued when they had been carried out to the instant at which the fire was expected to burn through the wall, that is to say, after all four plasterboard panels had disintegrated.

As may be seen from Fig. 20, which was taken from [19], the variation in the enthalpy of plasterboard with the temperature is dependent on whether

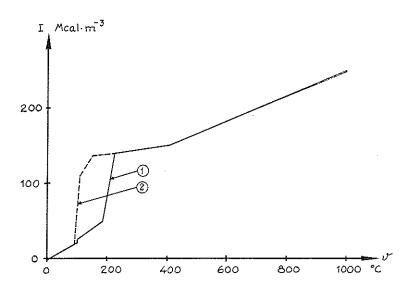


Fig. 20. Relation between the enthalpy, I, and the temperature, θ, used in the calculations for plasterboard. Curve 1: High rate of temperature rise. Curve 2: Low rate of temperature rise.

the rate of temperature rise is high or low. This circumstance was taken into account in the calculations by assuming a high rate of temperature rise for that plasterboard panel which is nearest to the fire and a low rate of temperature rise for all the other plasterboard panels. Since the structural transformations of plasterboard require an additional quantity of heat that cannot be recovered, the enthalpy-temperature curve during the cooling period is not identical with that during the heating period.

In the calculations, it was assumed that the variation in the enthalpy with the temperature during the cooling period is represented by a straight line which corresponds to a constant value of the product $c \cdot \gamma = 200 \text{ kcal} \times \text{m}^{-3} \cdot {}^{\circ}\text{C}^{-1}$.

For the surfaces exposed to fire in the enclosed spaces of all the types dealt with in the present chapter, the coefficient of heat transfer, α_i , was calculated by means of Eq. (2.7a), where it was assumed that $\varepsilon_{fi} = 0.7$ and $\varepsilon_i = 0.8$ for the Type A to the Type E enclosed spaces, and for the Type G enclosed space. These values give $\varepsilon_{res} \sim 0.60$. For the Type F enclosed space, the calculations were based on three values of ε_i , viz., 0.1, 0.4, and 0.8. Hence, for $\varepsilon_{fi} = 0.7$, the respective values of ε_{res} were found to be 0.1, 0.35, and 0.6. For the exterior surfaces of the structures bounding the enclosed space, the coefficient of heat transfer, α_u , for the Type A to Type E enclosed spaces, and for the Type G enclosed space, was supposed to vary with the surface temperature, ϑ_u , in accordance with Eq. (2.8), while its value for the Type F enclosed space was supposed to vary in conformity with the relation

$$\alpha_{u} = 7.5 + \frac{4.96 \cdot \varepsilon_{\text{res}}}{\theta_{u} - \theta_{0}} \left[\left(\frac{\theta_{u} + 273}{100} \right)^{4} - \left(\frac{\theta_{0} + 273}{100} \right)^{4} \right] \text{kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1}. \quad ^{\circ}\text{C}^{-1}$$
(7.1)

where ε_{res} was chosen in the same way as in the above-mentioned calculation of α_i .

The variation in the specific heat, c_p , of the combustion gases with the temperature is shown in Fig. 21 [3]. For the rest, the calculations were based on the principles which have been stated in Chapters 2 and 4.

The input values used in the calculations were only the opening factor $A \cdot \sqrt{H}/A_t$, in $m^{1/2}$, and the fire load, q, in Mcal·m⁻² of bounding surface area. If the radiation term, I_R , is disregarded, then Eq. (2.1) becomes independent of dimensions. In other words, if the values of $A \cdot \sqrt{H}/A_t$ and q are given, then the result will be independent of the terms A, H, and A_t entering into the opening factor. However, I_R is proportional to the total opening area, A, and, in order that I_R may be taken into account, it is necessary to specify the above-mentioned terms. This has been done in the calculations dealing with test results in Chapter 5. For the determination of the total

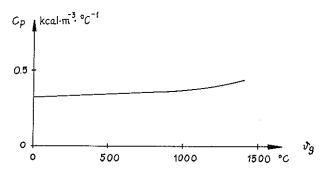


Fig. 21. Relation between the specific heat of the combustion gases, c_p , and their temperature, ϑ_q , used in the calculations.

opening area, A, in the calculations described in the present chapter, it was supposed that the dimensions of the enclosed space were the same as those which had been assumed as a basis for the curves published in the Swedish Building Regulations 1967, viz., a total bounding surface area $A_t = 10,000 \text{ m}^2$ and a square opening. Consequently, the ratio of the total opening area to the total bounding surface area, A/A_t , was in all cases lower than those values which can be expected to be met with in ordinary buildings, and all the results will therefore be on the safe side. The value $A_t = 10,000 \text{ m}^2$ was used only to determine the value of the ratio A/A_t , which was then substituted in the term I_R in the equation of heat balance. For a value of the opening factor $A \cdot \sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$, Table 9 gives the respective values of the ratio A/A_t which correspond to $A_t = 10,000 \text{ m}^2$ and $A_t = 1 \text{ m}^2$ on the assumption that the opening is square, or that it has a height H = 1 m.

Table 9. Values of the ratio A/A_t .

_	Square opening	H=1 m
$A_t = 1 \text{ m}^2$	0.075	0.04
$A_t = 10,000 \text{ m}^2$	0.012	0.04

In order to illustrate the consequence of this variation in the ratio A/A_t , Fig. 22 shows the temperature-time curve for an enclosed space characterized by an opening factor $A \cdot \sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$, a fire load $q = 30 \text{ Mcal} \cdot \text{m}^{-2}$, as well as by the two extreme values of the ratio A/A_t , i.e. 0.012 and 0.075. As is seen from this graph, the effect of the difference in the value of the ratio A/A_t is practically negligible.

Temperature-time curves have been calculated for each one of the seven enclosed spaces, Types A to G, for varying values of the duration of the fire and the opening factor. As regards the relation between the duration of the fire and the fire load, reference is made to Chapter 6. For the Type A to the

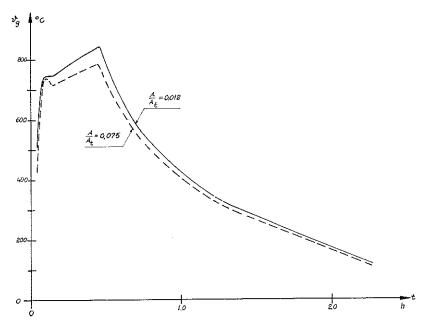


Fig. 22. Calculated temperature-time curves for the Type A enclosed space. Opening factor $A\sqrt[4]{H}/A_t = 0.04 \text{ m}^{1/2}$. Fire load $q = 30 \text{ Mcal} \cdot \text{m}^{-2}$ of bounding surface area. Ratio A/A_t set equal to 0.012 and 0.075, respectively.

Type E enclosed space, as well as for the Type G enclosed space, the temperature-time curves were computed on the basis of 6 different values of the opening factor, viz., $A \cdot \sqrt{H}/A_t = 0.01$, 0.02, 0.04, 0.06, 0.08, and 0.12 m^{1/2}. For each of these values the curves were computed for 8 different values of the duration of the flame phase of the fire for which the time graphs of the rate of combustion have been constructed in Chapter 6. This means that 48 temperature-time curves have been obtained for each type of enclosed space. In the case of the Type F enclosed space, the calculations were carried out for 3 different values of the resultant emissivity. Furthermore, for each one of these values, use was made of 5 different values of the opening factor, viz., 0.01, 0.02, 0.04, 0.08, and 0.12 $m^{1/2}$. For each one of these combinations of values, the curves were calculated on the basis of 5 different values of the fire load corresponding to 5 different values of the duration of the flame phase, viz., 0.1, 0.3, 0.5, 1.0, and 2.0 h, if computed by means of Eq. (2.1). All these curves are shown in Appendix 5. The curves are denoted by the symbols Al to G6, where the letter A refers to the Type A enclosed space, etc. All these curves are represented in an approximate form after smoothing out the irregularities which were caused by the polygonal shape of the time graphs of the rate of combustion. Furthermore, in order to render the graphs in Appendix 5 more readily legible, that part of each one of the curves which represents the ignition phase was based on that ascending branch of the time graph of the rate of combustion which corresponds to the lower four values of the fire load. The exact results of the calculations are reproduced in tabular form in Appendix 6. In the case of the Type G enclosed space, where the calculation of the temperature-time curves was carried out with reference to the instants when the plasterboard panels fell down, each one of these instants is marked with a circle on the corresponding curve. When a plasterboard panel falls down, this corresponds in the calculations to an instantaneous temperature drop in the enclosed space, as may be seen from the relevant curves.

In practical design, it should be possible to proceed in three steps, viz., first, to choose that type of enclosed space which is most closely similar in respect of the thermal properties of the bounding structures to the case under consideration; second, to determine the opening factor and the fire load; and third, to interpolate linearly between the values given in the tables in Appendix 6. If, instead of using this procedure, the designer chooses a curve which is determined without interpolation so as to be on the safe side, that is to say, if he chooses the next higher values of the opening factor and the fire load, then this will probably not involve errors which are too great. In order to afford a basis for the choice of the type of enclosed space, the temperature-time curves which correspond to an opening factor $A \cdot \sqrt{H}/A_t = 0.04$ $m^{1/2}$ and to a fire load of 60 Mcal \cdot m^{-2} of bounding surface area are represented in Fig. 23 for the Type A to the Type F enclosed spaces. For comparison, this graph also reproduces the standard ISO temperature-time curve and the curve for an opening factor of 0.04 m^{1/2} published in the Swedish Building Regulations 1967.

A comparison between the temperature-time curves which correspond to the different types of enclosed spaces in Fig. 23 shows that the maximum difference in the maximum temperature amounts to about 400°C. The Type C enclosed space, which is provided with lightweight concrete bounding structures, exhibits markedly higher temperatures than the other types of enclosed spaces comprised in the present calculations. The lowest maximum temperature was obtained in the case of the Type F enclosed space ($\varepsilon_{res} = 0.60$), which is equipped with bounding structures made of sheet steel, 2 mm in thickness. However, it is seen from Fig. 23 that the resultant emissivity for radiation between the flames and a sheet steel surface produces a substantial effect on the magnitude of the maximum temperature. If $\varepsilon_{\rm res}$ is supposed to change from 0.6 to 0.1, then the corresponding difference in the maximum temperature is slightly over 200°C, other conditions being equal. Therefore, it is important that the heat transfer conditions, and particularly the resultant emissivity, should be accurately determined in the calculation of the temperature-time curve for an enclosed space of this type.

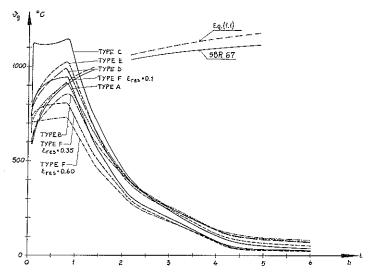


Fig. 23. Temperature-time curves for the Type A to the Type F enclosed spaces. Opening factor $A\sqrt[4]{H}/A_t=0.04~\mathrm{m}^{1/2}$. Fire load $q=60~\mathrm{Mcal}\cdot\mathrm{m}^{-2}$ of bounding surface area. Furthermore, this graph also shows the respective temperature-time curves calculated by means of Eq. (1.1) and determined in conformity with the Swedish Building Regulations 1967 (SBR 67) for an opening factor $A\sqrt[4]{H}/A_t=0.04~\mathrm{m}^{1/2}$.

With the exception of the Type F, it is the Type B enclosed space, which is bounded by concrete walls, that exhibits the lowest maximum temperature. This is due to the relatively high thermal conductivity and the great heat capacity of the concrete. On the other hand, since the large quantity of heat that is stored in the concrete is partly transferred back to the enclosed space during the cooling period, comparatively high temperatures are obtained in the course of the cooling phase.

The temperature-time curves published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures" relate to "enclosed spaces bounded by wall, roof or ceiling, and floor structures which are made of brickwork, concrete, or lightweight concrete as a material that is predominant in thermal respects". As has previously been mentioned, these curves had been calculated on the basis of those characteristics of the bounding structures which were used to describe the Type A enclosed space. Moreover, for guidance, the comments on the Draft Specification "Aluminium Structures" also comprise temperature-time curves for an enclosed space which is bounded by walls made of mineral wool. In addition, it is shown how the temperature-time curve is influenced by a concrete wall, 20 cm in thickness, which is situated in the interior of the enclosed space. By examining Fig. 23, it will readily be understood that a further differentiation of the above-mentioned Swedish standard temperature-time curves according to the thermal characteristics of the bounding structures would be desirable.

In order that the temperature-time curves for the cooling phase which have been determined in the present publication might be compared with the corresponding Swedish standard curves, Fig. 24 shows the temperature-time curves for the Type A enclosed space calculated on the basis of an opening factor $A \cdot \sqrt{H}/A_t = 0.04 \, \text{m}^1/^2$, together with the Swedish standard temperature-time curves for the cooling phase determined on the assumption that the rate of temperature decrease is $10\,^{\circ}\text{C} \cdot \text{min}^{-1}$. The ascending branches of the curves are identical because the curves for the Type A enclosed space are based on the same assumptions as the standard curves. The linear temperature-time curves for the cooling phase, which start from the ascending branch at the time T, calculated by means of Eq. (1.2), are represented by dash lines in Fig. 24. As is seen from this graph, the calculated curves result in a rate of cooling which is higher or lower than the standard rate of temperature decrease, $10\,^{\circ}\text{C} \cdot \text{min}^{-1}$, according as the duration of the process of fire development is shorter or longer, respectively.

Thus, if the duration of the fire is comparatively short, then an application of the temperature-time curves which have been computed in the present publication gives considerably more favourable results, i.e. lower temperatures, than the standard rules which are at present in force in Sweden. For instance, the temperature-time curve for the cooling phase in fires of short duration is a decisive factor in determining the temperatures of unprotected steel structures, as has already been shown in an example which was adduced under

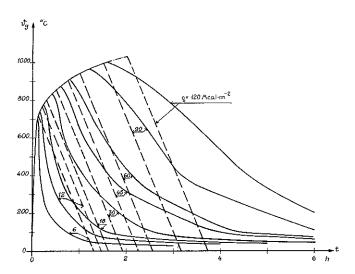


Fig. 24. Temperature-time curves for the Type A enclosed space. Opening factor $A\sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$. (Full-line curves.) Curves for the cooling phase corresponding to a rate of decrease in temperature of $10^{\circ}\text{C} \cdot \text{min}^{-1}$ stipulated in the relevant Swedish regulations. (Dash-line curves.)

the heading "Introduction". However, if the duration of the fire is relatively long, then the calculated temperature-time curves are less favourable, i.e. they give higher temperatures, than the curves stipulated in the Swedish regulations. All the same, since a given structure that is exposed to fire has already been subjected to the action of high temperatures for a long time before the beginning of the cooling phase, the difference between these curves in the latter case will have a comparatively slight effect on practical design.

8. Summary

The point of departure of the present investigation was the fact that the results obtained in recent years from research in the field of structural fire engineering had made it possible to carry out reliable calculations of the load bearing and separating capacity in the design of structural components exposed to fire. Such design calculations must be based on the knowledge of the temperature-time curve which covers the whole process of fire development. However, further progress towards realistic structural fire engineering design was impeded by the circumstance that it was not possible to make a theoretical determination of the temperature-time curve for the cooling phase of the process of fire development under known external conditions. Up to now, the research in this field has evolved methods for calculating the variation in the temperature with the time during the flame phase of the process of fire development on the basis of known external conditions, whereas the cooling phase has not been dealt with in this connection.

In consequence of this gap in our knowledge, the methods for determining the temperature-time curves for the flame phase and the cooling phase stipulated in the Swedish Building Regulations 1967 are widely different in degree of accuracy. The determination of the temperature-time curve for the flame phase shall be based on the fire load which characterizes the case under consideration, as well as on the shape and the dimensions of the openings in the enclosed space. For the cooling phase, on the other hand, it is stipulated only that the rate of temperature decrease shall be set equal to $10^{\circ}\text{C} \cdot \text{min}^{-1}$, irrespective of the actual conditions which characterize the case in question. This undifferentiated characterization of the cooling phase is particularly unfavourable to structures which possess a low thermal inertia, e.g. noninsulated or slightly insulated load-bearing steel structures. It was therefore considered to be urgently required to undertake an investigation in order to find out whether a theoretical determination of the temperature-time curve for the cooling phase would be possible.

The theoretical calculations in the present publication are founded on a basic equation of heat balance in an enclosed space which has been deduced by Kawagoe and Sekine, as well as by Ödeen. This equation states that the quantity of heat, I_c , which is released per unit time during the process of combustion is at any instant equal to the sum of the quantities of heat which

are withdrawn per unit time in different ways from the enclosed space. Heat is ordinarily abstracted from the enclosed space by heat transfer through the structures which bound the enclosed space (term I_W in the equation of heat balance), by radiation through the openings in the enclosed space (term I_R), and by the replacement of combustion gases by cold air (term I_L).

In order to extend the range of application of the equation of heat balance so that it might cover the whole process of fire development, it was necessary to solve two fundamental problems. In the first place, the quantity of heat released per unit time had to be determined as a function of the time for the entire process of fire development. In the second place, the expression for I_L which had been deduced previously, and which was applicable to the flame phase only, had to be extended and supplemented.

The study of the last-mentioned problem resulted in an expression for I_L which was based on the magnitude of the heat transfer by convection through the openings in the enclosed space. The rates of gas and air flow involved in this process were calculated in two steps, viz., first by determining the velocity distribution of gas flow in a vertical opening by which two masses of gas differing in density are separated from each other, and second, by satisfying the condition that the net exchange of gases between the enclosed space and its surroundings shall be equal to the difference between the quantity of gas produced and the quantity of air consumed in the process of combustion. After that, it was possible to determine I_L directly as the difference in heat content between the outgoing gases and the incoming air. It was found that I_L was approximately proportional to the temperature of the combustion gases and to the air flow factor $A \cdot \sqrt{H}$.

Since no physical basis is available which could enable the quantity of energy liberated per unit time during the process of fire development to be determined as a function of the time, a study of the literature was carried out with a view to an analysis of full-scale fire tests. For the tests where the external conditions were stated in a sufficiently precise manner, comparative calculations of temperature-time curves were made by means of a computer.

A tentatively chosen time graph of the quantity of energy liberated per unit time was used for this purpose. The time graph in question was varied until the agreement between the observed and calculated temperature-time curves became as close as possible. The only requirement to be fulfilled in this connection was that the total quantity of energy released during the whole process of fire development should be equal to that which was available in the fuel from the outset. When all those tests which were suited for this study had been examined, it was possible to systematize the results of the study in such a way that the time graph of the quantity of energy released per unit time during the process of fire development might generally be assumed to be known.

This procedure was primarily justified by the consideration that an error, if any, could only be involved in the time graph of the quantity of energy liberated per unit time, since the total magnitude of this quantity is determined by the fire load, i.e. by the quantity of energy which is available from the beginning.

The computer programme which was used for the calculation of the temperature-time curves has a far-advanced general validity. One of the features of this programme is that it affords a possibility of taking into account various factors, viz., first, those thermal properties of the materials entering into the structures bounding the enclosed space which are dependent on the temperature, second, the variations in the dimensions of the openings during the process of fire development, third, the moisture content of the bounding structures, and fourth, the effects of heat-absorbing structures in the interior of the enclosed space. This programme can be used for enclosed spaces which are bounded by structures of up to three different types at the same time, and one of these structures may be built of up to three different materials.

Moreover, a modified programme has been prepared for enclosed spaces provided with plasterboard panel walls, which are assumed under certain definite conditions to disintegrate during the fire.

The time graphs of the quantity of energy liberated per unit time which had been obtained by means of the method outlined in the above were used to calculate the time graphs of the temperature of the combustion gases during the process of fire development. The latter time graphs were computed on the assumption of different values of the fire load and the opening factor for seven types of enclosed spaces which differ in respect of the bounding structures. The results of these calculations are represented in graphs as well as in tables.

In carrying out the comparative theoretical analyses, it was possible to discuss to a limited extent the effects produced on the temperature-time curve of the process of fire development by some quantities which do not directly enter into the equation of heat balance, with the result that their effects on this process must be determined in each individual case. In addition to the size and the shape of the openings, the factors which may be expected to be of importance in this connection comprise, among others, the porosity of the fuel and its distribution in the enclosed space, the moisture content of the fuel, and the magnitude of the fire load. For future research in this field, it may be urgently recommended to make a study of the effects produced by these and other parameters on the quantity of energy released per unit time in the process of combustion, and hence also on the temperature-time curve, which is dependent on this quantity.

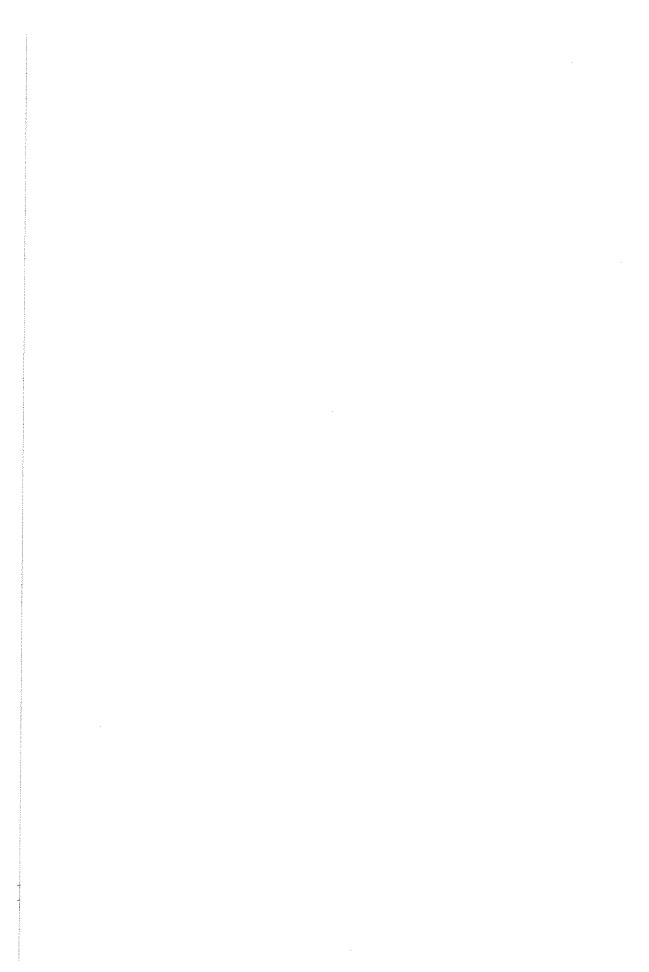
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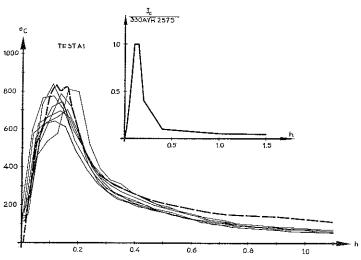
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Appendix

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series A. See chapter 5.



Test A1

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 34.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 42.2 per cent.

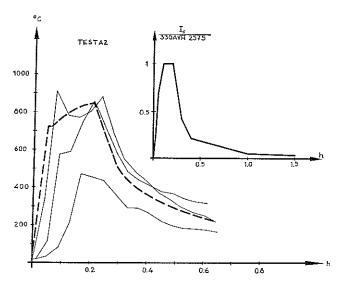
Concrete, 3 cm in thickness+lightweight concrete, 10 cm in thickness, 18.3 per cent.

Window area 4.7 per cent.

Opening factor 0.06 m^{1/2} (t > 0.1 h).

Duration of the fire 0.17 h.

Fire load 15.1 Mcal·m⁻² of bounding surface area.



Test A2

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 34.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 42.2 per cent.

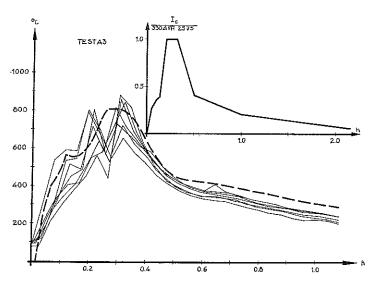
Concrete, 3 cm in thickness+lightweight concrete, 10 cm in thickness, 18.3 per cent.

Window area 4.7 per cent.

Opening factor 0.06 m^{1/2} (t > 0.1 h).

Duration of the fire 0.21 h.

Fire load 19 Mcal·m⁻² of bounding surface area.



Test A3

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 38.6 per cent.

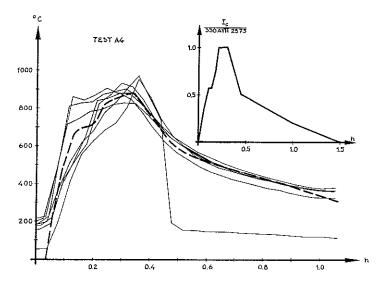
Lightweight concrete, 12.5 cm in thickness, 60.0 per cent.

Window area 1.4 per cent.

Opening factor 0.0356 m^{1/2} (t > 0.2 h).

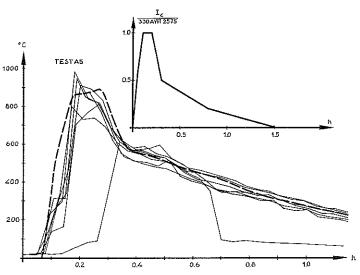
Duration of the fire 0.36 h.

Fire load 19.6 Mcal·m⁻² of bounding surface area.

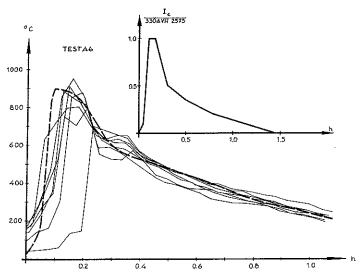


Test A4

Percentages of the total bounding surface area:
Concrete, 20 cm in thickness, 38.3 per cent.
Lightweight concrete, 12.5 cm in thickness, 59.4 per cent.
Window area 2.3 per cent.
Opening factor 0.0486 $m^{1/2}$ (t > 0.2 h).
Duration of the fire 0.325 h.
Fire load 22.4 Mcal· m^{-2} of bounding surface area.



Test A5
Percentages of the total bounding surface area:
Concrete, 20 cm in thickness, 38.3 per cent.
Lightweight concrete, 12.5 cm in thickness, 59.4 per cent.
Window area 2.3 per cent.
Opening factor $0.0548~\text{m}^{1/2}$.
Duration of the fire 0.24~h.
Fire load 20 Mcal·m⁻² of bounding surface area.



Test A6

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 37.8 per cent.

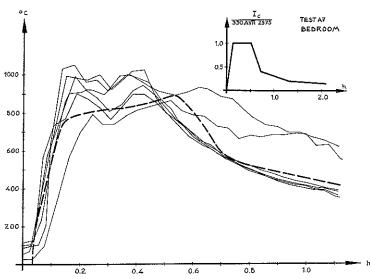
Lightweight concrete, 12.5 cm in thickness, 58.6 per cent.

Window area 3.6 per cent.

Opening factor 0.068 m^{1/2}.

Duration of the fire 0.23 h.

Fire load 23 Mcal·m⁻² of bounding surface area.



Test A7 a

Living-room.

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 40.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 47.3 per cent.

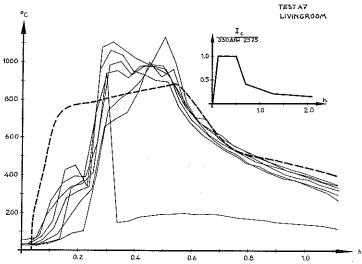
Concrete, 3 cm in thickness+lightweight concrete 10 cm in thickness, 8.3 per cent.

Window area 3.6 per cent.

Opening factor 0.04 m^{1/2}.

Duration of the fire 0.32 h.

Fire load 32 Mcal·m⁻² of bounding surface area.



Test A7b

Bedroom.

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 40.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 47.3 per cent.

Concrete, 3 cm in thickness+lightweight concrete 10 cm in thickness, 8.3 per cent.

Window area 3.6 per cent.

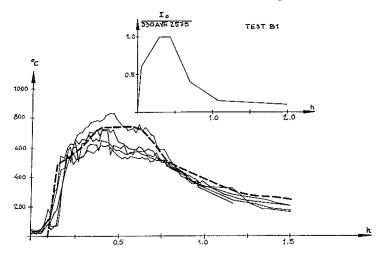
Opening factor 0.04 m^{1/2}.

Duration of the fire 0.32 h.

Fire load 32 Mcal·m⁻² of bounding surface area.

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series B. See chapter 5.

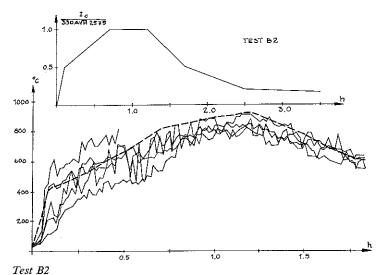


Test B1

Opening factor 0.0467 m^{1/2}.

Duration of the fire 0.48 h.

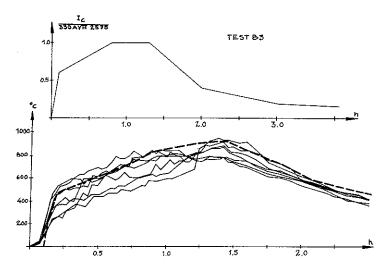
Fire load 33.3 Mcal·m⁻² of bounding surface area.



Opening factor 0.0467 m^{1/2}.

Duration of the fire 1.07 h.

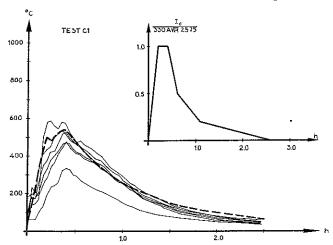
Fire load 75 Mcal·m⁻² of bounding surface area.



Test B3 Opening factor 0.0467 $m^{1/2}$. Duration of the fire 1.18 h. Fire load 83.5 Mcal· m^{-2} of bounding surface area.

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series C. See chapter 5.



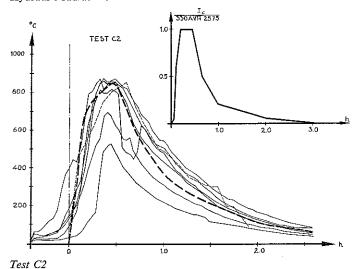
Test CI

Quantity of combustible material 270 kg.

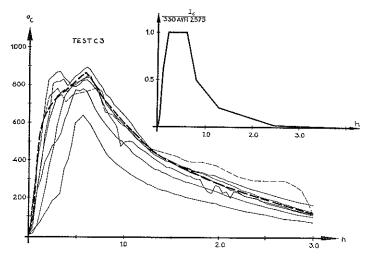
Rate of air supply by fan 1.0 m³/s.

Opening factor 0.0234 m¹ /2.

Hydraulic radius —.



Quantity of combustible material 675 kg. Rate of air supply by fan 2.0 m³/s. Opening factor 0.0601 m¹/². Hydraulic radius 1.0 cm.



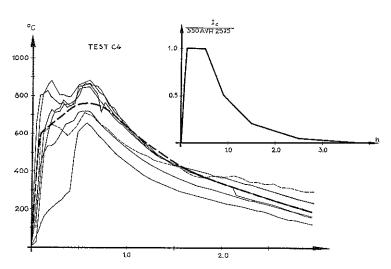
Test C3

Quantity of combustible material 675 kg.

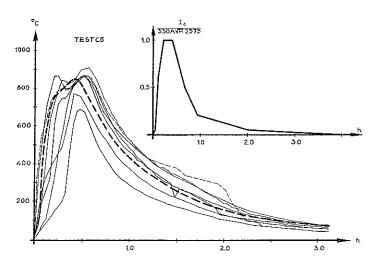
Rate of air supply by fan 1.0 m³/s.

Opening factor 0.0434 m¹/².

Hydraulic radius 1.0 cm.



Test C4 Quantity of combustible material 675 kg. Rate of air supply by fan 0.7 m 3 /s. Opening factor 0.0351 m $^{1/2}$. Hydraulic radius 1.0 cm.



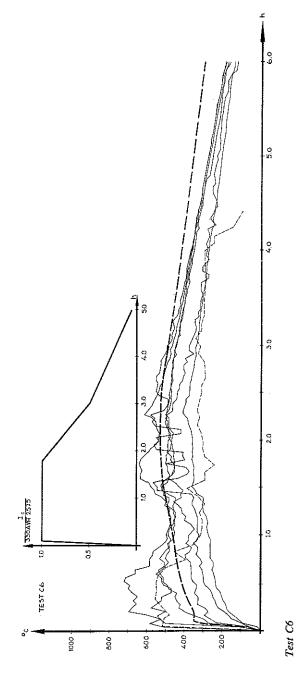
Test C5

Quantity of combustible material 675 kg.

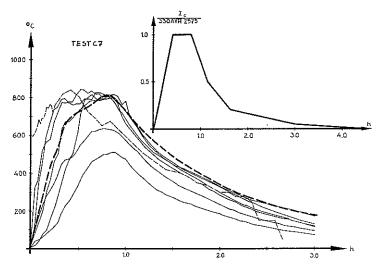
Rate of air supply by fan 1.5 m³/s.

Opening factor 0.0551 m¹/².

Hydraulic radius 1.0 cm.



Quantity of combustible material 675 kg. Rate of air supply by fan 0.25 m³/s. Opening factor 0.01 m^{1/2}. Hydraulic radius 1.0 cm.



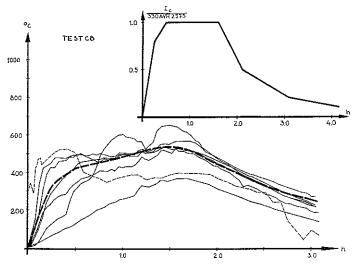
Test C7

Quantity of combustible material 675 kg.

Rate of air supply by fan 1.0 m³/s.

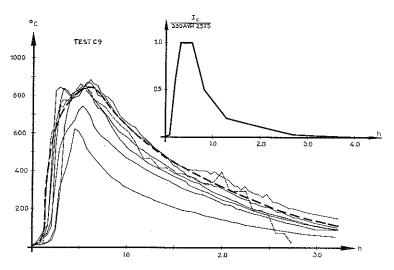
Opening factor 0.0367 m¹/².

Hydraulic radius 1.7 cm.



Test C8

Quantity of combustible material 675 kg. Rate of air supply by fan 0.7 m³/s. Opening factor 0.015 m¹/². Hydraulic radius 1.7 cm.



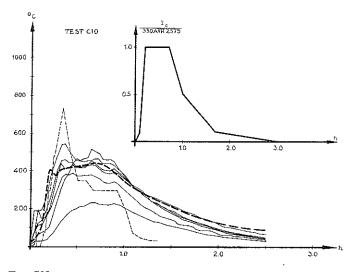
Test C9

Quantity of combustible material 675 kg.

Rate of air supply by fan 1.0 m³/s.

Opening factor 0.051 m¹/².

Hydraulic radius 0.6 cm.



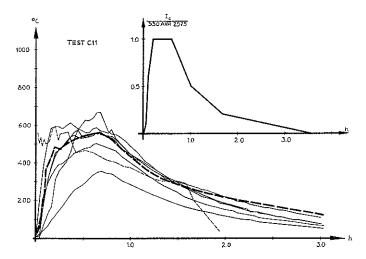
Test C10

Quantity of combustible material 270 kg.

Rate of air supply by fan 0.7 m³/s.

Opening factor 0.015 m¹ 1².

Hydraulic radius 1.0 cm.



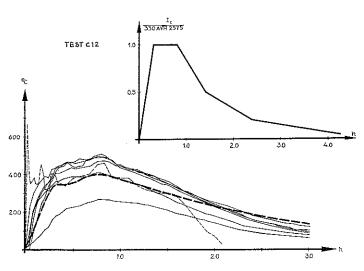
Test C11

Quantity of combustible material 405 kg.

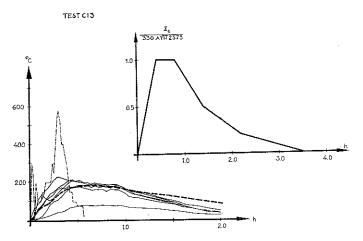
Rate of air supply by fan 0.7 m³/s.

Opening factor 0.02 m¹/².

Hydraulic radius 1.0 cm.



Test C12 Quantity of combustible material 405 kg. Rate of air supply by fan 0.7 m³/s. Opening factor 0.012 m^{1/2}. Hydraulic radius 2.4 cm.



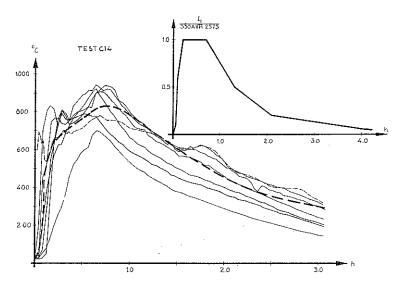
Test C13

Quantity of combustible material 135 kg.

Rate of air supply by fan 0.7 m³/s.

Opening factor 0.005 m¹/².

Hydraulic radius 1.0 cm.



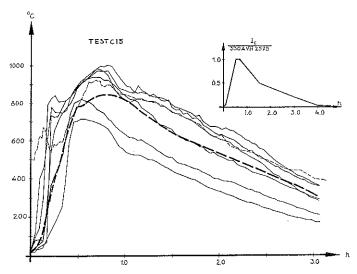
Test C14

Quantity of combustible material 945 kg.

Rate of air supply by fau 1.0 m³/s.

Opening factor 0.0367 m¹/².

Hydraulic radius 1.6 cm.



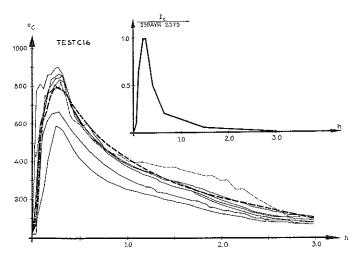
Test C15

Quantity of combustible material 1350 kg.

Rate of air supply by fan 2.0 m³/s.

Opening factor 0.0601 m¹/².

Hydraulic radius 1.4 cm.



Test C16

Quantity of combustible material 405 kg.

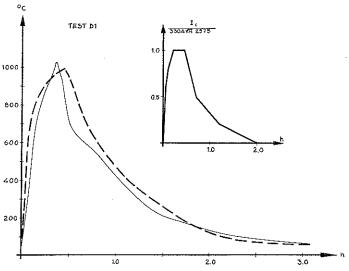
Rate of air supply by fan 0.7 m³/s.

Opening factor 0.06 m¹/².

Hydraulic radius 0.4 cm.

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series D. See chapter 5.



Test D1

Percentages of the total bounding surface area:

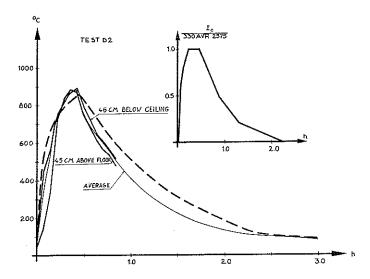
Lightweight concrete, 20 cm in thickness, 97.8 per cent.

Window area 2.2 per cent.

Opening factor 0.028 m^{1/2}.

Duration of the fire 0.5 h.

Fire load 20.2 Mcal·m⁻² of bounding surface area.



Test D2

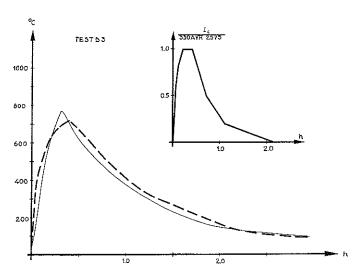
Percentages of the total bounding surface area:
Lightweight concrete, 20 cm in thickness, 98.7 per cent.

Window area 1.3 per cent.

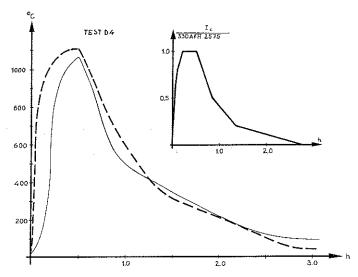
Opening factor $0.0147 \text{ m}^{1/2}$.

Duration of the fire 0.51 h.

Fire load $11.2 \text{ Mcal} \cdot \text{m}^{-2}$ of bounding surface area.



Test D3 Percentages of the total bounding surface area: Lightweight concrete, 20 cm in thickness, 99.15 per cent. Window area 0.85 per cent. Opening factor $0.0092 \, \mathrm{m}^{1/2}$. Duration of the fire 0.47 h. Fire load 6.5 Mcal \cdot m⁻² of bounding surface area.



Test D4

Percentages of the total bounding surface area:
Lightweight concrete, 94.7 per cent.

Window area 5.3 per cent.

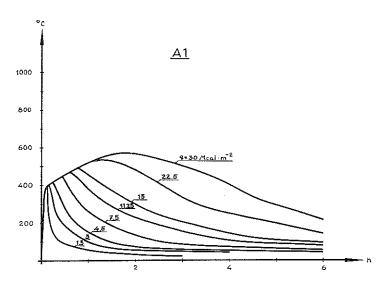
Opening factor 0.075 $\mathrm{m}^{1/2}$.

Duration of the fire 0.58 h.

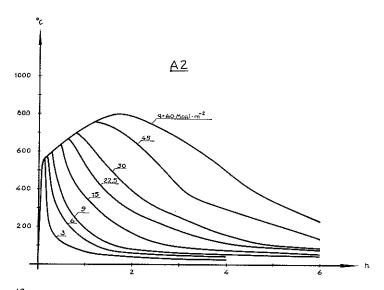
Fire load 65 Mcal· m^{-2} of bounding surface area.

Calculated time graphs of temperature of combustion gases for seven types of enclosed spaces differing in opening factor and in bounding structures.

See chapter 7.



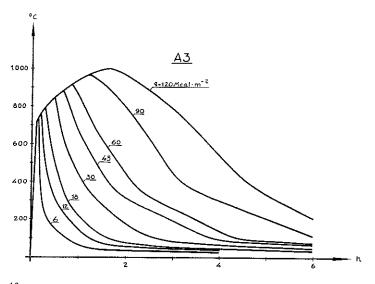
AI Type A enclosed space. Opening factor 0.01 $\rm m^{1/2}$.



A2

Type A enclosed space.

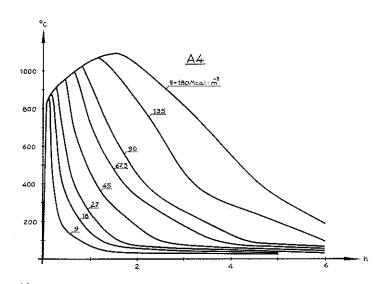
Opening factor 0.02 m^{1/2}.



A3

Type A enclosed space.

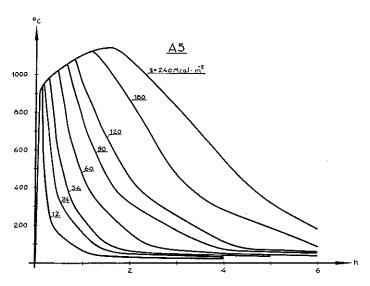
Opening factor 0.04 m^{1/2}.



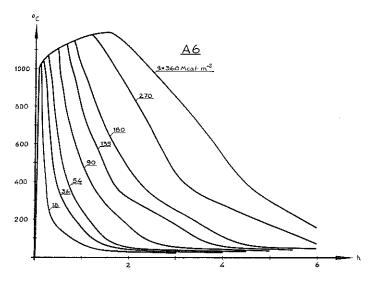
A4

Type A enclosed space.

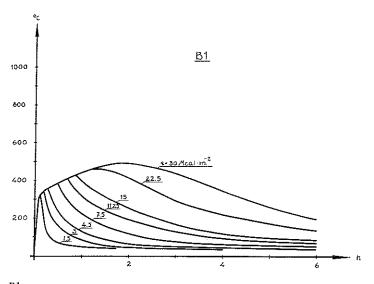
Opening factor 0.06 m^{1/2}.



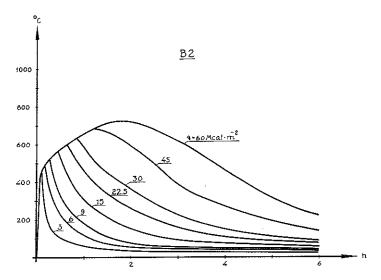
A5 Type A enclosed space. Opening factor 0.08 $m^{1/2}$.



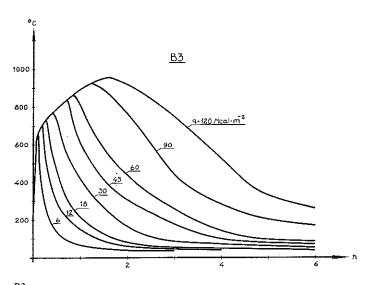
A6
Type A enclosed space.
Opening factor 0.12 m¹/².



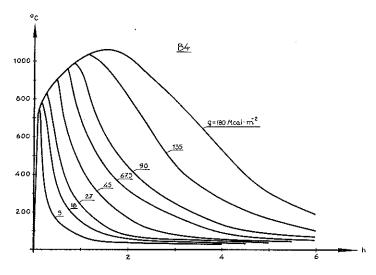
B1 Type B enclosed space. Opening factor 0.01 $m^{1/2}$.



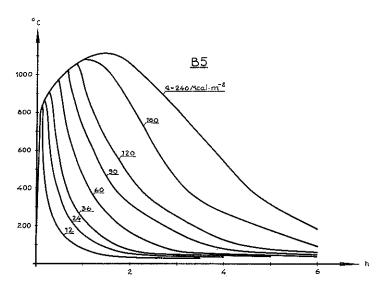
B2 Type B enclosed space. Opening factor 0.02 m^{1/2}.



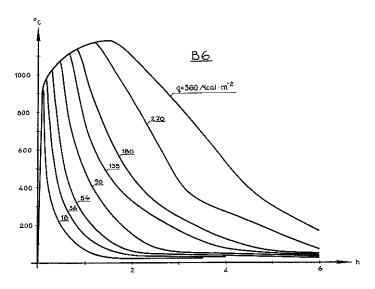
B3 Type B enclosed space. Opening factor 0.04 m $^{1/2}$.



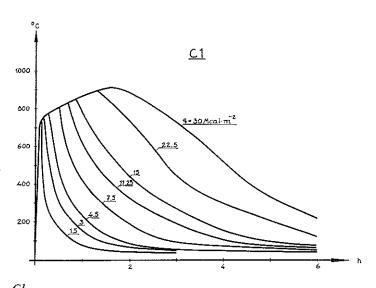
B4 Type B enclosed space. Opening factor 0.06 $m^{1/2}$.



B5 Type B enclosed space. Opening factor $0.08~{\rm m}^{1/2}$.

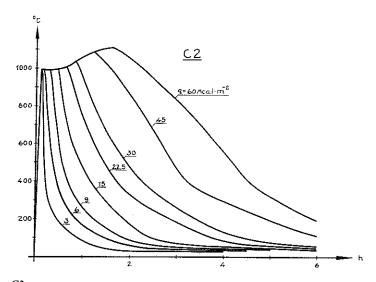


B6 Type B enclosed space. Opening factor 0.12 $m^{1/2}$.

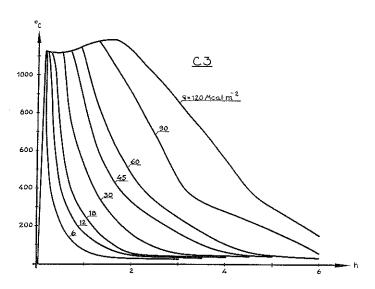


Type C enclosed space.

Opening factor 0.01 m^{1/2}.

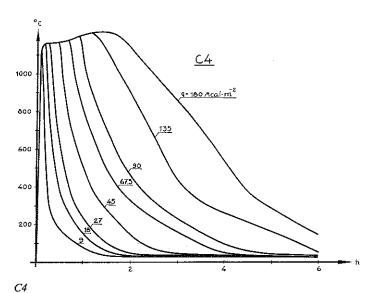


C2 Type C enclosed space. Opening factor 0.02 $\mathrm{m}^{1/2}$.

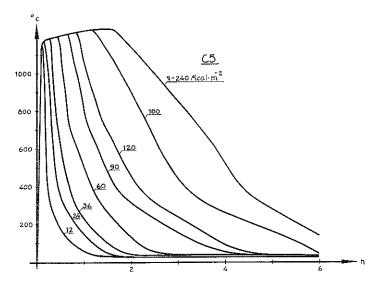


Type C enclosed space.

Opening factor 0.04 m^{1/2}.



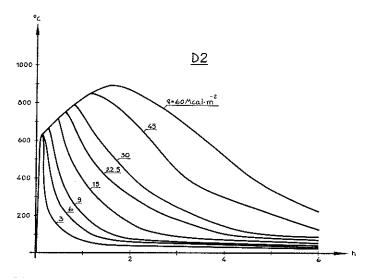
Type C enclosed space.
Opening factor 0.06 m^{1/2}.



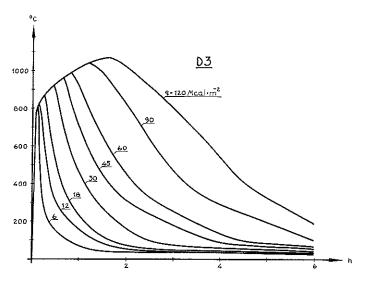
C5

Type C enclosed space.

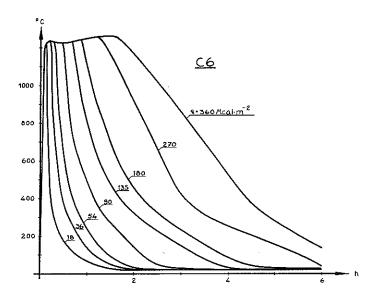
Opening factor $0.08 \text{ m}^{1/2}$.



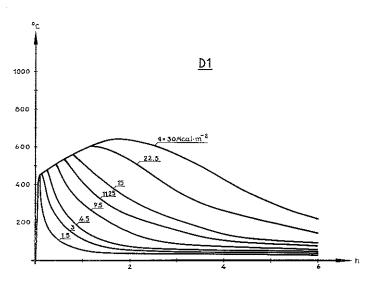
D2 Type D enclosed space. Opening factor 0.02 $\mathrm{m}^{1/2}$.



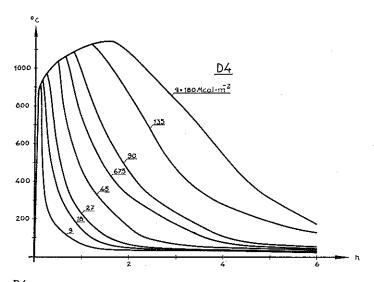
D3 Type D enclosed space. Opening factor 0.04 $\mathrm{m}^{1/2}$.



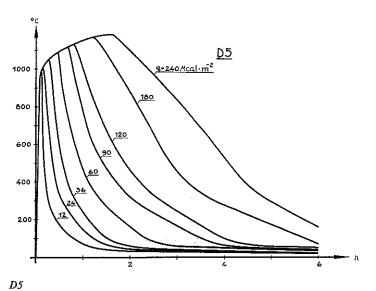
C6 Type C enclosed space. Opening factor $0.12 \ m^{1/2}$.



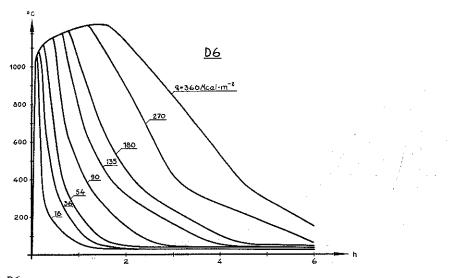
D1 Type D enclosed space. Opening factor 0.01 $m^{1/2}$.



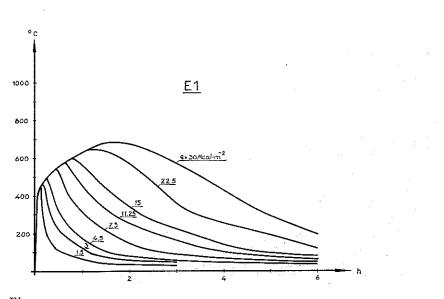
D4 Type D enclosed space. Opening factor $0.06~\text{m}^{1/2}$.



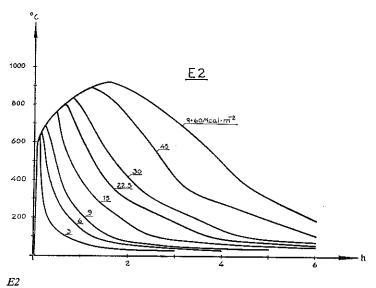
Type D enclosed space.
Opening factor 0.08 m^{1/2}.



D6 Type D enclosed space. Opening factor 0.12 m^{1/2}.

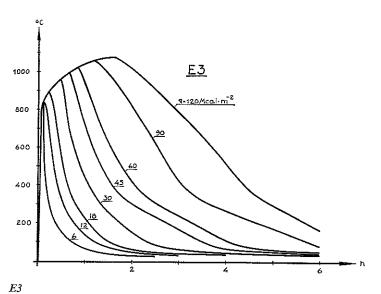


EI Type E enclosed space. Opening factor 0.01 $\mathrm{m}^{1/2}$.



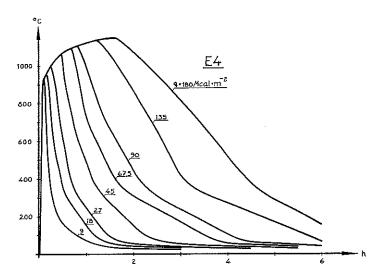
Type E enclosed space.

Opening factor 0.02 m^{1/2}.

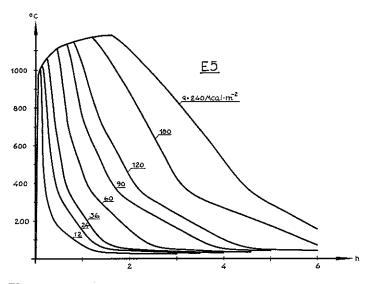


Type E enclosed space.

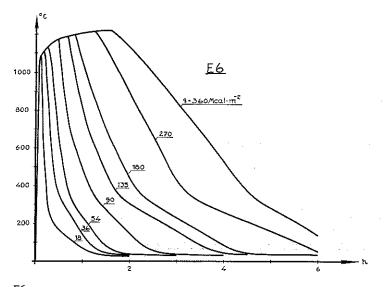
Opening factor 0.04 m^{1/2}.



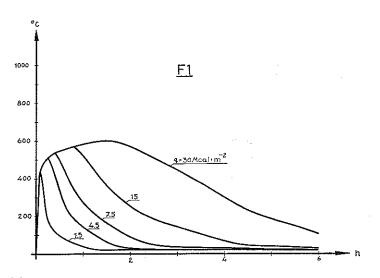
E4 Type E enclosed space. Opening factor 0.06 m^{1/2}.



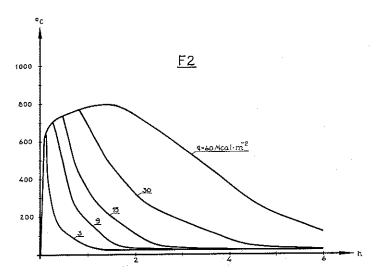
E5 Type E enclosed space. Opening factor 0.08 $\mathrm{m}^{1/2}$.



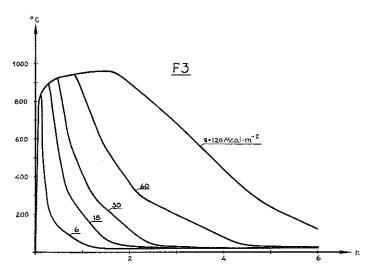
E6
Type E enclosed space.
Opening factor 0.12 m¹/₂.



FI
Type F enclosed space.
Resultant emissivity 0.1.
Opening factor 0.01 $\,\mathrm{m}^{1/2}$.



F2
Type F enclosed space.
Resultant emissivity 0.1.
Opening factor 0.02 m^{1/2}.

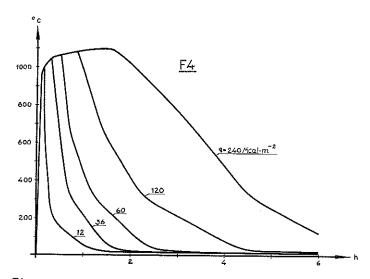


F3

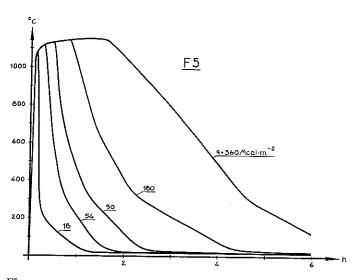
Type F enclosed space.

Resultant emissivity 0.1.

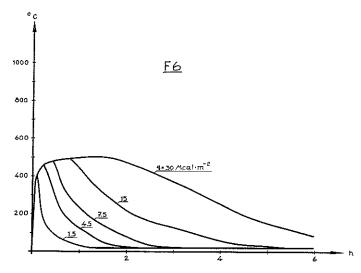
Opening factor 0.04 $m^{1/2}$.



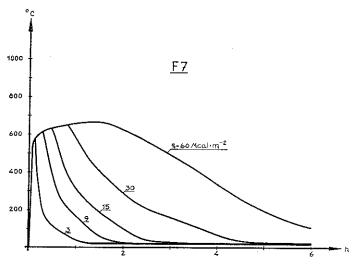
F4
Type F enclosed space.
Resultant emissivity 0.1.
Opening factor 0.08 m^{1/2}.



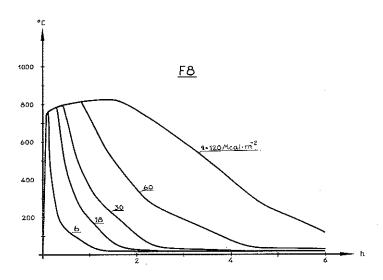
F5 Type F. enclosed space. Resultant emissivity 0.1. Opening factor 0.12 $\mathrm{m}^{1/2}$.



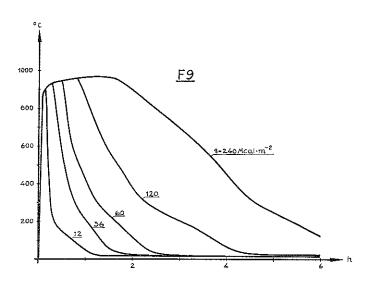
F6
Type F enclosed space.
Resultant emissivity 0.35.
Opening factor 0.01 m^{1/2}.



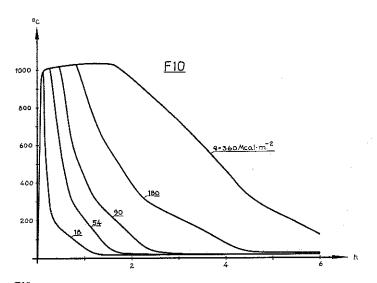
F7
Type F enclosed space.
Resultant emissivity 0.35.
Opening factor 0.02 m^{1/2}.



F8 Type F enclosed space. Resultant emissivity 0.35. Opening factor 0.04 $m^{1/2}$.



F9
Type F enclosed space.
Resultant emissivity 0.35.
Opening factor 0.08 m¹/₂.

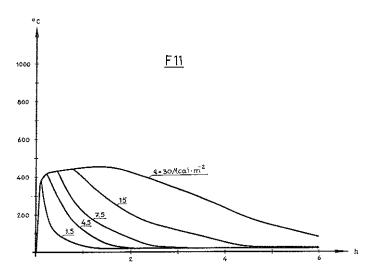


F10

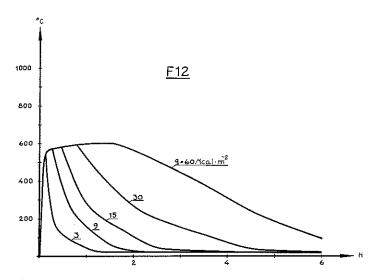
Type F enclosed space.

Resultant emissivity 0.35.

Opening factor 0.12 m^{1/2}.



F11 Type F enclosed space. Resultant emissivity 0.60. Opening factor 0.01 $m^{1/2}$.

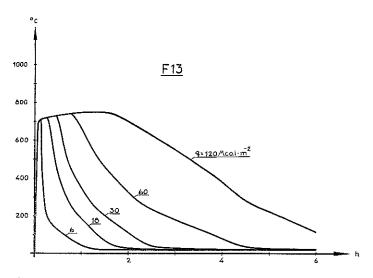


F12

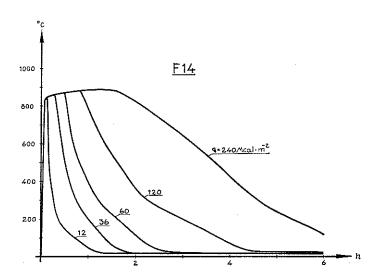
Type F enclosed space.

Resultant emissivity 0.60.

Opening factor 0.02 m^{1/2}.



F13 Type F enclosed space. Resultant emissivity 0.60. Opening factor 0.04 $m^{1/2}$.

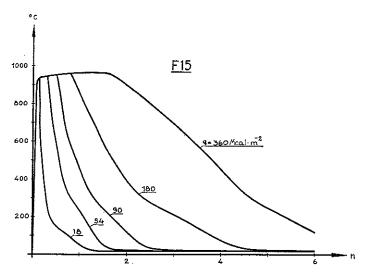


F14

Type F enclosed space.

Resultant emissivity 0.60.

Opening factor 0.08 m^{1/2}.

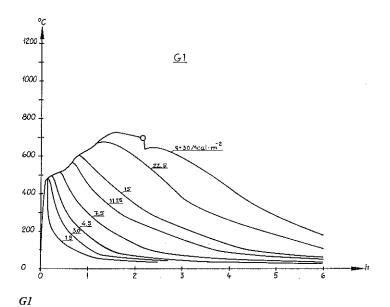


F15

Type F enclosed space.

Resultant emissivity 0.60.

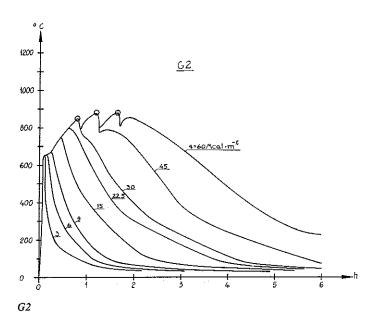
Opening factor 0.12 $m^{1/2}$.



Type G enclosed space.

Opening factor 0.01 m^{1/2}.

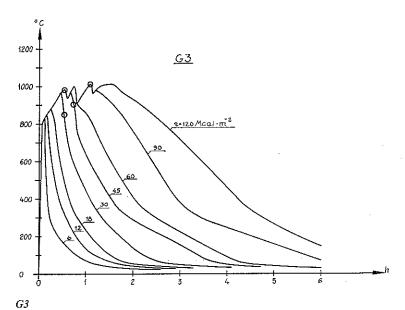
O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.02 m¹/².

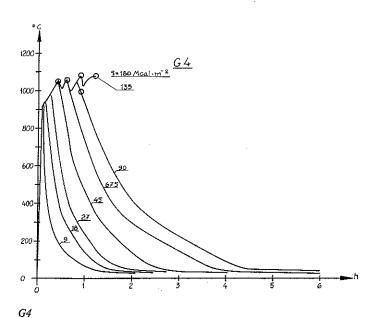
O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.04 m^{1/2}.

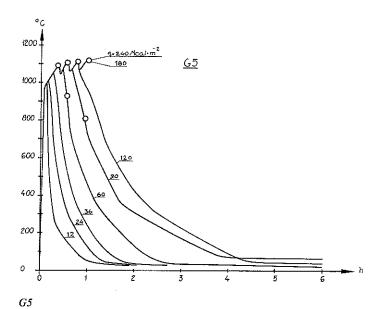
O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.06 m^{1/2}.

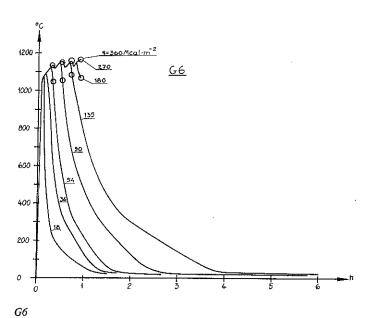
O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.08 m¹/².

O=Plasterboard panel falls down.



Type G enclosed space. Opening factor $0.12~\text{m}^{1/2}$. O=Plasterboard panel falls down.

APPENDIX 6

Calculated time graphs of temperature of combustion gases represented in tabular form for seven types of enclosed spaces differing in opening factor and in bounding structures. See chapter 7.



 $\frac{\textbf{A1}}{\text{Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space.}}$ Opening Factor A. $\sqrt{\text{H}}/\text{A}_{\text{t}}$ = 0.01 m^{1/2}

0.05		L								
Time h 0.05	T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
10.05	q.	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0	
0.10 395 395 395 395 395 328 328 328 328 328 0.15 228 390 390 360 360 360 360 360 360 0.20 196 368 h01 h01 h06 h06 h06 h06 h06 h06 0.25 150 313 h09 h10 h05 h05 h05 h05 h05 h05 h05 h05 h05 h0		T	e m	p	е	r a	ı t	u	r	e
6.00 34 41 57 77 92 140 216	0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60	395 395 196 99 99 88 79 59 14 40 39 37 36 35 44 44 40 39 38 37 33 33 33 33 33 33 33 33	395 396 397 398 399 398 397 397 397 397 397 397 397 397 397 397	39901 9501 9501 9501 9501 9501 9501 9501	395 390 401 421 429 437 441 370 355 314 302 283 263 242 218 209 190 150 150 150 177 83 80 77 75 86 66 65 66 66 65 66 66 66 66 66 66 66 66	328 360 405 405 4125 449 460 451 463 463 463 463 463 463 463 463 463 463	328 3606 4055 4497 4497 4497 4497 4566 4319 4497 4566 4319 4319 4319 4319 4319 4319 4319 4319	320 320 320 320 320 320 320 320 320 320	3606555429306801111110979792102294701196153727272 360655542930680111111097555555555498611961537272727	

<u>A2</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.02 m^{1/2}

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0
Time h	Ŧ	e m	p	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.60 2.80 3.40 3.50 3.40 3.60 3.40 3.60 3.40 3.60 3.40 3.60 3.80 4.00 4.60 4.80 5.00	396 568 322 277 210 134 131 126 121 107 102 98 88 77 66 54 45 44 40 38 36 35 31 33 32 31 30 30	396 556 559 361 361 361 361 361 361 361 361 361 361	3968 5575 5775 5775 5776 5776 5776 5776 57	396 568 556 572 587 6015 632 485 449 359 328 275 2614 221 199 184 81 110 89 87 77 71 86 64 61 61 61 61 61 61 61 61 61 61 61 61 61	397 511 581 578 561 562 663 664 664 665 677 677 677 677 677 677 677 677 677	397 511 581 578 593 612 662 645 645 645 645 645 645 645 645 645 645	397 5111 581 5783 5793 6422 642 643 644 7746 648 648 648 648 328 221 245 241 245 245 245 245 245 245 245 245 245 245	398 467 581 578 5915 662 662 671 687 7456 687 775 778 778 788 789 663 663 741 741 741 741 741 741 741 741 741 741

A3

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{\rm H}/\rm A_t$ = 0.04 m $^{1/2}$

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time h	T	e m	P	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.40 4.60 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	504 745 360 868 164 165 168 1128 1100 87 14 14 15 14 16 16 16 16 16 16 16 16 16 16 16 16 16	504 745 747 696 587 438 338 259 432 264 85 27 40 85 40 40 39 38 37 40 38 38 37 40 40 39 38 38 38 38 38 38 38 38 38 38 38 38 38	545774744453330260556883688773665554475442109887736655544372222883688877366655520887736	504 747 767 767 767 767 767 767 824 768 824 768 824 825 826 827 828 831 831 841 841 841 841 841 841 841 84	141	504 681 776 793 882 862 7768 882 894 862 7768 882 894 862 7768 863 87768 863 87768 8	504 681 7776 7938 8894 9188 9199 9199 9199 9199 8131 655 7188 2104 2104 1688 1531 111	504 681 777 778 802 838 874 882 918 918 910 917 917 910 913 913 913 914 915 917 917 918 918 918 918 918 918 918 918 918 918

<u>A4</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.06 m^{1/2}

	•							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	p	е	r a	t	u	r e
h 0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20	T 575 858 4934 496 159 1516 120 120 48 120 48 136 35 332 31 30 29	575 851 802 938 490 303 272 215 118 776 652 964 442 40 976 555 444 40 3376 334 333 333 333 333 333 333	575 858 861 879 898 838 7669 572 433 400 272 243 300 272 243 184 155 386 81 74 71 661 57 54 44 44 44 41 40	575 858 861 879 891 928 954 7694 620 574 526 469 327 303 259 621 211 21 21 21 21 21 21 21 21 21 21 21	575 704 784 882 889 903 949 961 982 939 8795 505 463 348 332 289 263 210 185 159 131 103 81 77	575 7044 889 908 901 901 901 901 901 901 901 901 901 901	575 704 784 882 889 908 923 949 961 1018 1054 1064 1029 1013 966 830 756 830 756 830 756 830 756 831 355 331 308 288 269	575 704 784 882 890 908 923 949 961 982 10018 1032 1044 1064 1072 1087 1087 1087 1087 1087 1089 943 943 943 956 756 756 756 756 759 754 754
4.40 4.60		32 31	38 37	53 51	73 70	90 85	249 230	423 377
4.80 5.00		31 30	36 35	49 47	67 64	81 77	211 193	348 319
5.20 5.40 5.60			35 34 33	46 44 43	62 59 57	74 71 69	174 155 135	292 265 238
5.80		•	32 32	42 41	55 54	6 ₁ ,	114 94	210 185

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.08 m $^{1/2}$

	U							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	T	e m	ď	е	r a	. t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.60 5.60 5.80 6.00	622 935 532 432 314 180 171 163 155 139 131 126 89 70 47 44 42 40 39 37 36 33 31 30 29 28 27	622 935 937 869 734 575 521 454 386 259 269 259 151 70 662 596 551 49 46 38 33 32 31 33 32 31 33 32 31 33 31 34 35 36 37 37 38 38 38 38 38 38 38 38 38 38 38 38 38	6235 9375 9375 938 936 936 943 936 943 936 943 955 955 955 955 955 955 955 955 955 95	622 935 937 975 975 975 1013 1024 937 1021 937 937 937 937 937 937 937 937 937 937	626 766 767 768 769 769 769 769 769 769 769 769	622 766 853 965 908 1050 1055 1066 1085 836 695 618 753 655 454 454 877 871 885 665 755 877 877 877 877 877 877 877 877 87	622 766 853 959 965 981 1050 1058 1066 1081 1092 11119 1077 1058 1038 1049 932 853 775 6114 373 302 282 222 203 166 125 104 82	622 767 853 959 965 982 995 1008 1050 1058 1066 1081 1109 1119 1126 1132 1138 1105 1074 1058 1016 971 923 873 820 769 714 658 599 539 478 415 838 338 309 225 472 199 172

<u>A6</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.12 m^{1/2}

	U							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time h	T	e m	P	e	r s	ı t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.60 5.80	670 1027 581 465 333 186 185 176 168 159 142 133 124 106 88 67 41 39 37 36 31 32 31 30 29	670 1027 1033 951 799 620 556 480 404 324 292 275 221 186 149 107 63 50 47 45 43 39 37 35 33 32	670 1027 1033 1049 1063 1981 1063 1982 174 659 341 108 67 145 108 69 754 40 40 33 34 33 32 34 33 33 33 33 33 33 33 33 33 33 33 34 34	670 1027 1033 1049 1063 1076 1088 1098 1007 1004 856 492 556 492 253 277 253 288 208 100 429 100 443 46 49 39 736 49 337 337 337 49 40 40 40 40 40 40 40 40 40 40 40 40 40	670 847 933 1057 1071 1083 1094 1103 11127 1133 1060 971 873 7650 628 575 469 414 353 2257 200 173 144 46 47 41 41 41 42 41 42 41 42 41 42 41 42 41 42 41 42 41 42 41 42 42 43 44 44 44 44 44 44 44 44 44 44 44 44	670 847 933 1057 1071 1083 1109 1103 11139 1150 1062 1001 937 868 713 672 9585 452 360 272 245 665 162 162 162 162 163 164 165 166 166 166 166 166 166 166 166 166	670 847 933 1051 1057 1071 1083 1194 1103 11127 1133 1150 1159 1166 1173 1178 1106 1082 1043 1003 962 919 874 789 704 613 519 423 361 336 331 3293 273 273 273 273 273 273 273 273 273 27	670 847 933 1051 1057 1071 1083 1199 1159 1159 1166 1173 1188 1192 1195 1151 1195 1195 1195 1195 1195
6.00				31	38	45	67	158

<u>B1</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.01 m $^{1/2}$

q 1.5 3.0 4.5 7.5 11.25 15.0 22.5 30.0 Time T e m p e r a t u r h 0.05 237 352 352 <td< th=""><th>op</th><th>t</th><th></th><th></th><th></th><th></th><th></th><th></th><th></th></td<>	op	t							
Time	T	0.1	0.2					1.5	2.0
0.05	<u>q</u>	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
0.05	=	T	e m	P	е	r a	t	u	r e
2.60 43 54 82 158 201 338 459 2.80 41 52 78 147 190 311 447 3.00 40 51 75 136 179 281 432 3.20 39 49 72 125 168 259 417 3.40 38 48 70 114 156 248 401 3.60 38 47 68 103 145 237 384 3.80 37 46 66 95 134 227 366 4.00 36 45 64 91 122 217 346 4.20 36 44 62 88 110 208 325 4.40 35 43 61 85 105 199 302 4.50 35 42 59 83 102 190 284 4.80 34 42 58 80 98 181 271	h 0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.60 5.80	237 342 201 172 132 86 85 76 77 65 59 53 46 38 37 36 35 35 34 33 33	237 342 342 343 321 321 321 321 321 321 321 321 321 32	234398655895388112111111111111111111111111111111111	237 342 339 348 3560 378 357 360 245 260 245 260 250 260 260 260 260 260 260 260 260 260 26	2376422031352366933880443555336936728895044722096938884043897388950447221143591888538887765532	237 236 314 352 369 369 372 388 400 400 403 3761 3300 400 403 3761 3300 400 400 400 400 400 400 400 400 40	286 314 352 352 360 377 388 400 4120 439 445 451 447 407 385 331 259 2190 1164 515 1164 515 1164 515 1164 515 1164 515 1164 515 1164 515 1164 515 515 515 515 515 515 515 515 515 51	284 315520 315560 3172 3880 400 400 400 400 400 400 400 400 400

<u>B2</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A. $\sqrt{E}/A_{\rm t}$ = 0.02 m $^{1/2}$

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
ď	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0	
Time h	T	е ш	p	е	r a	. t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	3538 498 249 1118 1109 507 658 454 42 42 42 43 33 33 33 33 33 33 33 33 33 33 33 33	3538 4875 3223 3073 2240 407 507 507 507 507 507 507 507 507 507 5	354 487 551 443 3336 3336 3336 352 213 164 68 80 77 77 66 66 67 67 67 67 67 67 67 67 67	3548 487 26 6 6 5 3 1 9 9 8 5 2 8 5 2 9 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	354 415116008675555555555555555555555555555555555	354 45116 5555 5555 566 5575 5526 463 4210 333 303 457 598 599 5555 566 5755 566 5755 566 5755 566 5755 566 5755 566 5755 566 567 5756 5	355 411 516 528 559 561 563 663 663 663 663 663 663 663 663 663	354 4511 55020 555 5643 6637 7127 7102 6637 7129 6639 7129 6639 7129 7129 7129 7129 7129 7129 7129 712	

<u>B3</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A. $\sqrt{\rm E}/\rm A_t$ = 0.04 m $^{1/2}$

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time h	T	e m	р	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.10 2.10 2.60 2.80 3.00 3.10 3.10 3.10 3.10 3.10 3.10 3.1	481 6754 331 152 152 132 112 106 41 40 41 40 33 33 33 33 33 33 33 33 33	481 675 658 615 537 404 361 403 314 322 217 140 147 67 65 60 60 58 41 42 41 42 41 43 38 37 36 36 37 46 37 47 47 47 47 48 48 49 49 49 49 49 49 49 49 49 49 49 49 49	481 6758 6758 6759 6759 6759 6759 6759 6759 6759 6759	487586696677255411068875554110688755541106887555411099881774196644208755422018683199881774196644208755542	482 557 608 693 709 725 746 802 733 675 733 675 529 454 417 380 341 327 330 276 263 276 116 116 116 116 116 116 117 70	482 5576 693 688 705 706 802 707 708 802 803 705 705 803 803 705 705 803 805 705 705 805 805 805 805 805 805 805 8	457 558 606 693 688 709 778 806 822 855 876 893 907 919 889 878 856 833 807 778 889 429 378 429 378 353 310 273 255 238 203 187 170 153 135	457866388952977668888899999999998888897766888899999999
5.80 6.00		33 32	39 38	52 51	70 68	82 80	135 117	231 206

 $$\underline{B4}$$ Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space Opening Factor A. \sqrt{H}/A_t = 0.06 m $^{1/2}$

Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	р	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.80 3.00 3.20 3.40 3.60 3.80 4.20 4.40 4.60 4.80 5.00 5.40 5.60 5.60 6.00	528 761 440 378 281 170 152 145 130 122 115 100 85 69 47 44 43 41 39 38 37 36 35 31 30 30	528 761 761 760 761 760 460 460 460 460 460 460 460 460 460 4	528 761 766 791 817 817 817 818 818 819 819 819 819 819 819 819 819	528 766 798 818 827 838 877 850 833 877 833 877 877 877 877 877 877 877	528 632 703 805 805 805 905 869 905 869 905 860 905 860 905 861 905 805 805 805 805 805 805 805 805 805 8	528 532 533 533 530 530 530 530 530 530	528 632 703 806 805 805 935 949 909 909 1032 985 995 909 1032 985 991 878 808 737 762 768 768 768 768 768 768 768 768 768 768	528 632 703 8065 8059 907 908 908 807 909 908 8065 909 908 909 1005 1005 908 908 908 908 908 908 908 908 908 908

<u>B5</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A $\sqrt{\rm H}/\rm A_t$ = 0.08 m^{1/2}

	v							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	T	e m	p	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.60 3.60 3.60 3.60 3.60 3.60 4.20 4.40 4.60 5.20 5.60 5.80 6.00	576 836 489 301 177 159 151 137 69 47 45 40 38 36 35 33 33	576 836 4798 437 437 437 437 437 437 437 437 437 437	76644177742433301111088317307288419653433333333333333333333333333333333333	5736441774444 871774444 871774444 8717777777777	552 505 457 408 359 343 327	576 693 773 886 9135 9918 1027 9918 1035 9918 1058 816 606 5528 448 3304 222 1146 117 88 87 77 768 663 663	576 694 773 871 886 910 955 973 1018 1027 1036 1037 1056 1038 1019 985 1019 985 1038 1043 1056 1038 1043 1056 1056 1057 1058 1058 1059 1058 1059 1058 1059 1058 1059 1058 1059 1059 1059 1059 1059 1059 1059 1059	576 694 773 871 886 910 935 973 990 1018 1056 1077 1087 1109 1112 1112 1042 1058 1073 1058 1073 1058 1075 1112 1058 1075 1112 1075 1112 1075 1112 1075 1112 1075 1112 1075 1112 1075 107

<u>B6</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor $A\cdot\sqrt{H}/A_{\rm t}$ = 0.12 m^{1/2}

	ı							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time	T	e m	р	е	r a	ı t	u	r e
h								
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	624 940 548 447 324 185 183 174 166 157 143 122 105 87 33 33 33 31 31	624 940 952 892 7620 541 397 320 289 272 255 149 66 63 556 63 556 44 42 337 36 35 33 32 31 30	624 940 950 1940 1940 1940 1940 1940 1940 1940 194	640 9580 9580 100518 367128 31794 31	624 777 867 985 1040 1076 1060 1076 1060 1076 1060 1076 1060 1076 1076	624 777 980 995 1040 1079 1060 1079 1120 1120 1120 1120 1120 1120 1120 112	624 777 980 995 1019 1060 1076 11089 11103 1159 1165 1190 1193 1159 1166 1090 1031 1090 1031 1090 1031 1090 1031 1090 1031 1090 1031 1090 1091 1091	624 777 867 980 995 10140 1060 1076 11133 1152 1159 1165 1171 1185 1185 1181 1105 1086 988 827 773 565 469 404 2189 2189 2189 2189 2189 2189 2189 2189

<u>C1</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor A+ $\sqrt{H}/A_{\rm t}$ = 0.01 m^{1/2}

Ţ	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	T	е т	Þ	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 6.00 5.20 5.40 6.00	502 773 347 2556 1548 2156 113 108 760 755 44 44 42 43 33 33 33 33 33	502 7769 708 708 418 3197 221 1513 207 1523 207 1523 207 1523 207 1523 207 1523 207 1523 207 1533 207 207 207 207 207 207 207 207 207 207	502 7769 7779 5698 4579 2775 2861 28776 4598 4579 2778 2778 2778 2778 2778 2778 2778 27	5078 5078 5079 5079 5079 5079 5079 5079 5079 5079	502 630 696 7776 808 821 808 825 826 827 8604 827 8604 827 8604 821 821 821 821 821 831 831 831 831 831 831 831 831 831 83	502 630 647 776 788 795 808 821 828 828 829 643 597 685 546 431 826 827 828 829 829 821 829 821 821 821 821 821 821 821 821 822 823 824 824 825 826 827 827 827 828 829 829 829 829 829 829 829 829 829	502 630 677762 808 825 834 854 854 854 854 855 865 865 865 865 865 865 865 865 865	502 630 696 777 788 808 821 788 828 828 836 845 856 869 890 914 897 897 897 897 720 684 528 836 845 856 875 877 877 887 888 888 889 890 890 890 890 890 890 890

<u>C2</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor $\Lambda \cdot \sqrt{H}/A_{\rm t}$ = 0.02 $\rm m^{1/2}$

	τ								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
đ	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0	
Time h	T	e m	р	е	r s	. t	u	r	е
0.05 0.15 0.20 0.20 0.35 0.45 0.67	634 992 538 425 303 177 160 153 130 122 107 45 42 43 37 36 35 33 31 30 28	634 992 977 889 734 563 437 265 2188 73 666 576 45 45 45 45 45 45 45 45 45 45 45 45 45	639 977 9801 9801 9801 9801 9801 9801 9801 9801	639 977 989 999 999 999 999 999 999 999 99	6305 884 988 989 999 999 998 999 998 803 7640 770 877 877 877 877 877 877 877 877 87	634 805 988 989 999 999 10010 1032 998 1010 1032 1032 1032 1032 1032 1032 1032	634 805 807 808 908 908 909 909 1010 10	634 805 884 986 985 995 997 999 1016 1045 1057 1069 1078 1078 1079 1079 1079 1079 1079 1079 1079 1079	

	, - t								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0	
Time	T	е	m p	е	r	a, t	u	r	е
h			_						
0.05	726	726	726	726	726	726	726	726	
0.10	1156	1126	1126	1126	916	916	916	916	
0.15	618	1110	1110	1110	1002	1002	1002	1002	
0.20 0.25	478 334	1005 829	1111 1112	1112 1111	1122 1111	1122 1111	1122 1111	1122 1111	
0.30	183	631	1013	1114	1113	1113	1111	1113	
0.35	182	559	902	1115	1115	1115	1115	1115	
0.40	174	480	784	1116	1118	1118	1118	1118	
0.45	166	401	651	1115	1117	1117	1117	1117	
0.50	158	320	593	1005	1115	1115	1115	1115	
0.60	141	290	469	852	1126	1126	1126	1126	
0.65	132	273	404	770	1131	1132	1132	1132	
0.70	123	256	338	679	1058	1137	1137	1137	
0.80	106	221	308	621	970	1148	1148	1148	
0.90 1.00	88 67	186	276	557 491	874 363	1061 1002	1158 1166	1158 1166	
1.10	41	149 107	244 211	424	767 684	938	1172	1172	
1.20	39	63	178	353	631	930 870	1178	1178	
1.30	37	60	144	328	579	797		1182	
1.40	35	55	107	302	526	717	1105	1184	
1.50	34	52	68	278	473	676	1081	1187	
1.60	32	48	66	254	418	631	1043	1189	
1.70	31	46	62	229	361	589	1003	1146	
1.80	30	43	58	204	345	545	962	1128	
1.90	29	41	54	180	328	500	919	1110	
2.00 2.20	28	39 36	51 46	155	312	456 363	875 7 90	1091 1045	
2.40		36 34	40 42	100 69	283 255	363 331	703	996	
2.60		32	39	61	227	300	614	945	
2.80		30	37	55	199	272	521	890	
3.00		29	35	51	172	244	425	834	
3.20		·	33	47	143	216	362	780	
3.40			31	74.74	112	189	336	722	
3.60			30	41	81	161	312	661	
3.80			29	39	61	131	291	599	
4.00 4.20			28	37	56	100	271	535	
4.40				36 34	51 48	66 62	251 231	471 404	
4.60				33	45	57	211	355	
4.80				32	43	53	193	324	
5.00				31	41	50	173	295	
5.20				-	39	47	153	267	
5.40					38	45	132	239	
5.60					36	43	110	210	
5.80					35	41	89	184	
6.00					34	39	64	155	

T 0.1 0.2 0.3 0.5 0.75	1.0 1	.5 2.0
		0
q 9.0 18.0 27.0 45.0 67.5	90.0 135	.0 180.0
Tîme Tempera	t u	r e
h		
0.20 500 1053 1165 1165 1178 0.25 346 867 1166 1166 1166 0.30 185 659 1060 1168 1167 0.35 185 581 941 1167 1167 0.40 176 498 815 1166 1167 0.45 168 411 674 1166 1166 0.50 159 323 611 1048 1166 0.60 142 292 477 884 1177 0.65 132 275 408 795 1181 0.70 123 257 336 696 1100 0.80 105 220 307 634 1003 0.90 85 183 273 564 898	766	1052 1178 1166 1167 1167 1167 1167 1167 1167 1167 1167 1168 1181 1181 1181 1181 1181 1181 1194 1194 1194 1194 1207 1211 1211 1211 1211 1211 1211 1211 1211 1214 1222 1174 123 1134 1134 1134 1134 1104 1101 1211 1212 1218 1222 123 124 125 126 127 129 129 129 129 129 129 129 129

<u>C5</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$ = 0.08 m^{1/2}

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
g	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0	
Time h	T	e m	p	е	r s	t	u	r	e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.40 2.40 2.40 2.40 2.50 3.60 3.80 4.00 4.00 4.60 4.80 5.00 5.60 5.60 5.80 6.00 5.80 6.00	788 1215 668 512 352 186 186 177 168 160 142 132 104 61 32 28 27 26 25 24 23	788 1215 1196 1079 889 675 593 504 416 325 293 275 219 181 97 46 42 39 37 333 32 31 30	788 1215 1196 1196 1086 963 832 687 410 337 305 239 205 171 135 40 38 33 31 29 27	788 1215 1196 1196 11996 11997 11996 11997 1073 901 807 705 640 5699 496 422 219 169 2194 47 42 386 333 32 298 288	788 987 1079 1210 1196 1197 1197 1197 1206 1209 1123 1021 910 790 697 639 582 525 467 408 334 318 303 275 218 191 163 132 100 66 42 38 36 34 31 30	788 987 1079 1210 1197 1197 1197 1206 1219 1219 1219 1219 1219 1219 1219 121	788 987 1079 1210 1196 1197 1197 1296 1297 1291 1219 1225 1229 1231 1233 1174 1120 1076 1032 987 940 1032 987 940 1032 1076 1032 1032 1032 1032 1032 1032 1032 1032	788 987 1079 1196 1197 1197 1197 1197 1209 1213 1229 1231 1233 1235 1238 1240 1169 1149 1128 1076 1022 9659 124 169 1149 1149 1149 1149 1149 1149 1149	

<u>C6</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor $A \cdot \sqrt{H}/A_{\rm t} = 0.12~{\rm m}^{1/2}$

	U								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
ā	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0	
Ti me h	Ţ	е	m p	е	ŗ	a t	u	r	е
0.05 0.10 0.15 0.20 0.30 0.35 0.40 0.50 0.60 0.70 0.80 0.90 1.00 1.30 1.50 1.60 1.70 1.80 2.20 2.40 2.40 2.40 2.40 2.40 2.40 2.4	812 1250 688 525 359 186 186 177 169 160 141 132 103 82 58 28 27 26 25	812 1250 1230 1130 912 691 606 513 421 326 293 275 256 217 179 138 93 38 35 33 30 29 28 27 26	812 1250 1230 1230 1230 1115 987 851 701 632 485 411 335 303 270 236 202 167 130 90 40 38 35 33 32 31	812 1250 1230 1230 1230 1231 1230 1229 1100 918 820 713 645 571 496 420 342 316 2266 241 215 190 164 137 78 38 35 32 30	812 1014 1109 1245 1230 1230 1231 1237 1239 1148 1039 922 796 700 641 582 523 464 403 329 215 188 159 128 215 188 159 22 215 188 159 22 215 215 215 215 215 215 215 215 215	812 1014 1109 1245 1230 1230 1230 1231 1239 1247 1136 1063 986 905 820 728 683 635 739 491 442 314 286 258 202 174 145 113 80 37 31 32 31 29	812 1014 1109 1245 1230 1230 1230 1231 1253 1254 1255 1256 11964 1089 1043 105 106 107 107 107 108 108 108 108 108 108 108 108 108 108	812 1014 1109 1245 1230 1230 1230 1231 1237 1239 1247 1250 1257 1258 1257 1258 1257 1260 1260 1185 1142 1087 1031 972 849 724 658 590 521 452 381 332 303 275 247 219 163 132	

<u>D1</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A. $\sqrt{\rm H}/\rm A_t$ = 0.01 m $^{1/2}$

T 0.1 0.2 0.3 0.5 0.75 1.0	1.5 2.0
	00.5
q 1.5 3.0 4.5 7.5 11.25 15.0 2	22.5 30.0
Time Temperat h	u r e
0.15	315 315 376 411 411 456 456 460 460 472 472 484 484 494 502 502 509 525 532 532 532 532 532 533 566 566 579 579 591 603 594 624 591 631 591 63

<u>D2</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.02 m $^{1/2}$

<u>D3</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.04 m $^{1/2}$

_	U								
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0	
Time	T	e m	р	e	r a	t	u	r e	:
h								•	
0.05	547	547	547	547	547	547	547	547	
0.10	797	797	797	797	660	660	660	660	
0.15	456	802	802 825	802 825	734 826	734 826	734 826	734 826	
0.20 0.25	387 286	752 638	851	851	837	837	837	837	
0.30	172	510	799	870	862	862	862	862	
0.35	169	470	729	889	881	881	881	881	
0.40	161	416	643	906	898	898	898	898	
0.45	154	358	552	921	915	915	915	915	
0.50	147	295	515	851	930	930	930	930	
0.60	132	270	420	740	955	955	955	955	
0.65	124	255	370	674	966	966	966	966	
0.70	117	241	319	603	91 ¹ 4	974	974	974	
0.80	102	210	293	559	849	990	990	990	
0.90	87	181	266	507	774 688	930	1006	1006	
1.00	70 51	150 116	238 210	455 400	622	887 839	1019 1031	1019 1031	
1.20	9± 48	80	182	343	580	785	1042	1031	
1.30	45	77	154	321	539	725	1009	1051	
1.40	43	72	123	298	495	661	994	1060	
1.50	42	69	91	277	452	627	978	1068	
1.60	40	65	87	255	406	592	949	1075	
1.70	39	62	.83	234	360	557	918	1045	
1.80	37	60	79	211	344	521	886	1032	
1.90	36	57	75	190	328	484	851	1020	
2.00	35 34	55 51	72 67	169 123	313 286	446 368	815 743	1006 969	
2.40	32	48	62	95	260	337	672	929	
2.60	31	46	58	87	234	308	597	887	
2.80	30	43	. 55	81	209	281	519	841	
3.00	29	42	52	76	185	255	437	793	
3.20		40	50	72	. 159	229	380	745	
3.40		39	48	69	133	204	354	696	
3.60		37	46	65	105	179	330	644	
3.80		36	44 1. 2	62	89	153	308	591	
4.00 4.20		35 34	43 42	60 57	84 80	126 98	288 269	537 481	
4.40		34 34	40	55	76	93	250	401	
4.60		33	39	53	73	88	231	378	
4.80		32	38	51	70	84	212	348	
5.00		32	37	50	68	81	194	320	
5.20			37	48	65	77	176	293	
5.40			36	47	63	75	157	267	
5,60			35	46	61	72	138	240	
5.80			34 31	44 5-2	59 57	70	118	214	
6.00			314	43	57	68	98	188	

<u>D4</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A* $\sqrt{H/A_t}$ = 0.06 m $^{1/2}$

<u>D5</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$ = 0.08 m^{1/2}

	u							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
ď	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	T	e m	Þ	е	r a	t	u	r e
	T 642 975 560 454 327 185 123 125 67 43 33 32 31 30	e m 642 975 983 916 776 607 546 473 221 108 63 55 52 0 149 41 38 65 54 41 38 65 55 55 60 60 60 60 60 60 60 60 60 60 60 60 60	975 983 1008 9866 1030 9866 403 9866 403 9375 243 211 146 63 596 48 49 337 56 48 49 337 56 48 49 337 56 49 337 56 49 337 56 49 337 56 49 337 56 49 49 49 49 49 49 49 49 49 49 49 49 49	e 6425 9838 10399 1067 10899 10999 8432 6550 4150 27538 4150 2755 444 442 443 338	642 803 888 1020 1041 1058 1075 1087 1095 1109 1110 958 757 621 7519 463 342 577 146 578 115 115 115 115 115 116 116 116 116 116	t 2386500 105557 5055 5055 318 501 105557 5055 318 501 105557 5055 318 501 1055 1055 1055 1055 1055 1055 105	042 803 888 1005 1020 1041 1058 1075 1095 1109 1116 1122 1134 1153 1160 1166 1117 1095 1071 1032 993 952 910 866 782 698 609 518 423 362 337 314 293 273 273 214 196	r e 642 803 888 1005 1020 1041 1058 1075 1087 1095 1109 1116 1122 1134 1153 1160 1166 1172 1176 1180 1184 1086 1039 990 938 888 774 657 595 533 469 457 327
5.00 5.20 5.40				37 36 35	49 47 45	58 56 53	177 157 137	298 271 243
5.60 5.80 6.00				34 33	44 42 41	51 49 48	115 83 71	215 188 161

<u>D6</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A+ $\sqrt{H}/A_{\rm t}$ = 0.12 m^{1/2}

Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
Q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time h	T	e m	p	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.60 5.80 6.00 5.80 6.00 5.20 5.40 5.60 5.80 6.00 5.20 5.40 5.60 5.60 5.60 5.80 6.00 5.60 5.60 5.60 5.60 5.60 5.80 6.00 5.60	697 1056 605 483 343 189 160 142 133 105 86 65 37 35 34 32 29 28 27	697 1056 1070 991 834 646 5794 412 251 446 251 146 39 37 33 39 29	697 1056 1079 1079 1113 1021 9805 118 410 9805 410 9805 410 9805 410 9805 410 9805 410 9805 410 9805 410 9805 410 9805 410 9805 410 9805 4105 4105 4105 4105 4105 4105 4105 41	697 1056 1070 1194 11128 11156	697 874 965 1093 1124 1135 1154 1159 1176 1095 8776 632 7526 460 3337 321 168 107 540 47 537 40 337 40 337 40 337 40 337 40 337 40 337 40 337 40 337 40 337 40 40 40 40 40 40 40 40 40 40 40 40 40	6974 8765 11091 11159 1115	697 874 965 1093 1104 1159 1171 1176 1180 1197 1208 1212 1155 1130 1061 1018 974 1061 1018 979 881 793 328 609 513 328 266 266 266 266 27 188 189 189 189 189 189 189 189 189 189	697 874 965 1093 1104 1121 1135 1148 1159 1171 1190 1197 1203 1218 1215 1218 1221 1215 1218 1215 1218 1215 1218 1215 1218 1223 1174 1134 1114 1063 1010 955 838 781 721 658 394 394 316 288 261 233 205 177 149

<u>E1</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A $\sqrt{\rm H}/\rm A_t$ = 0.01 m $^{1/2}$

-	· t							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
р	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	Ţ	e m	p	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	312 472 280 242 189 130 127 112 103 95 86 77 68 56 52 46 44 42 40 39 38 37 35 31 30	312 472 465 443 321 3075 210 188 163 148 80 163 148 80 163 148 80 164 128 80 165 149 149 149 149 149 149 149 149 149 149	22552475730826204826912272885266185520875443200988766185520887766618754432009887766618752087758877666187544320098877666187520887776661875208776661875443200988777666187520877666187520877666187520877666187520877666187520877666187520877666187520877666187520877666187666187666060000000000	312252446788.66507133322649633074180025088776853544307513322649633071888877685555555544307133226496330744180025088776855555555549	312 392 486 488 502 514 572 580 573 573 573 573 574 575 575 575 575 575 577 577 577 577	312 312 312 312 312 312 312 312 312 312	312 3928 488 488 5523 472 588 602 603 603 603 603 603 603 603 603 603 603	312 392 486 488 5016 517 5180 620 532 534 620 632 643 656 665 666 667 668 668 669 669 669 669 669 669 669 669

<u>E2</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor $A \cdot \sqrt{E}/A_{\rm t} = 0.02~{\rm m}^{1/2}$

	ı								
T	0.1	0.2		0.5	0.7	75 1.0	1.5	2.	0
<u>ď</u>	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.	0
Time h	T	е	m p	e	r	a t	u	r	e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.20 4.40 4.60 4.80 5.00 5.60 5.80 6.00	441 6378 328 249 163 146 153 101 87 40 56 44 40 38 37 33 33 30	441 637 642 642 643 643 643 643 643 643 643 643 643 643	637 642 665 683 643 607 553	417 645 665 70138 7715 588 475 4381 222 222 107 88 707 63 66 68 57 57 57 57 57 57 57 57 57 57 57 57 57	436 598 671 731 792 807 7730 608 550 481 339 3308 2708 140 988 771 663 59 220 140 110 988 771 663 59 67 778 663 778 663 778 663 778 778 778 778 778 778 778 778 778 77	536 598 671 674 693 712	446 538 671 7748 815 861 877 861 877 860 877 860 877 860 877 860 877 860 877 860 877 860 877 860 877 860 877 860 877 860 877 877 877 877 877 877 877 877 877 87	443 536 598 671 674 671 731 748 815 869 881 901 919 909 901 882 882 853 882 871 671 630 588 497 749 400 360 333 849 215 219 229 238 249 249 259 269 269 279 279 279 279 279 279 279 279 279 27	

<u>E3</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$ = 0.04 m $^{1/2}$

<u>E4</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A+ $\sqrt{\pi}/\Lambda_{\rm t}$ = 0.06 m^{1/2}

	u							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	13.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	ā	е	r a	t	u	r e
3.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.20 3.40 3.20 3.40 3.20 3.40 3.60 3.80 4.90 5.90	615 929 532 445 323 199 170 161 145 126 108 91 149 45 40 396 333 331 30	615 929 937 882 7598 406 409 274 252 288 114 769 460 553 461 38 34 32	615 929 937 9358 9358 9353 9358 9353 9358 9358 9358	615 9237 9587 90132 9013	615 766 854 959 1002 1058 1070 1085 1091 1058 1091 1058 1091 1058 1091 1058 1091 1058 1091 1058 1091 1058 1091 1058 1091 1058 1091 1091 1091 1091 1091 1091 1091 109	615 615 615 615 615 616 617 617 617 617 617 617 617 617 617	615 766 854 976 1024 1058 1070 1085 1091 1107 1112 1129 1134 1092 1070 1017 1011 973 893 850 684 1070 1070 1070 1070 1070 1070 1070 107	615 766 854 960 976 1002 1058 1070 1085 1097 11129 1134 1139 1143 1147 1151 1160 1015 864 805 768 864 768 864 768 864 768 864 875 864 875 876 877 877 878 878 878 878 878 878 878

<u>E5</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor $\text{A}\cdot\sqrt{\text{H}}/\text{A}_{\text{t}}$ = 0.08 m^{1/2}

	ū							
T	0.1	0,2	0.3	0.5	0.75	1.0	1.5	2.0
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time	T	e m	p	е	r a	t	u	r e
h								
Time h 0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60				e 663 999 566 578 310 26 578 310 27 311 32 32 32 32 32 32 32 32 32 32 32 32 32	r a 663 831 913 1027 1043 1067 1086 1101 1113 1122 1133 1138 1068 979 767 623 571 517 463 408 352 335 319 304 277 250 222 195 169 140 109 79 58 54 49 46 43	663 831 913 1027 1043 1067 1086 1101 1123 1138 1151 1069 587 798 444 353 266 238 210 457 127 588 210 457 257 268 257 268 278 278 278 278 278 278 278 278 278 27	10663 831 913 1027 1043 1067 1086 1101 1113 1138 1158 1164 1169 1173 1102 1076 1037 954 1168 954 1168 954 1168 954 1168 954 1168 1173 1173 1173 1173 1173 1173 1173 117	e 663 831 913 1027 1043 1067 1086 1101 1113 1122 1133 1138 1143 1151 1158 1164 1169 1173 1177 1180 1183 1186 1143 1124 1105 1086 1038 987 934 879 821 765 706 645 582 519 455 389 341
4.60 4.80 5.00 5.20 5.40 5.60 5.80					43 41 39 37 36 35 34	55 50 47 44 42 40 39	225 206 187 168 148 128 106 85 60	389 341 312 285 258 231 203 177 148
6.00					33	37	00	T-40

<u>E6</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A $\sqrt{\rm H}/\rm A_t$ = 0.12 m $^{1/2}$

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
Q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time h	T	e m	p	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.60 4.80 5.00 5.60 5.80 6.00	709 1086 610 494 194 192 183 164 145 166 87 666 38 31 29 28 27 25	709 1086 1087 1013 860 670 594 336 297 278 260 222 184 145 102 55 50 46 43 41 39 37 36 34	709 1086 1087 1108 1130 944 827 692 488 417 343 307 441 207 174 138 100 58 42 41 42 39 35 33 30 41 42 39 31 30 41 41 42 31 31 30 41 41 41 41 41 41 41 41 41 41 41 41 41	709 1086 1087 1108 1130 1148 1161 1169 1175 1069 903 811 709 637 5494 420 245 220 245 220 195 145 89 440 37 335 332 31	709 899 985 1106 1121 1156 1178 1187 1190 1111 1013 4688 576 520 463 576 405 3315 274 40 37 331 103 40 37 40 37 40 37 333 333 333 333 333 333 333 333 333	709 899 985 1106 11178 11566 11178 11187 1	709 899 985 1106 1122 1141 1155 1166 1178 1190 1205 1209 1215 1160 1134 1064 1021 1064 1021 1064 1087 1088 109	709 899 985 1106 1122 1141 1155 1166 1178 1190 1205 1215 1218 1220 1215 1215 1221 1215 1221 1221 1175 1131 1062 1007 1007 1007 1007 1007 1007 1007 100

<u>F1</u>

Time Graphs of	Temperature	of	Combustion	Gases.	Туре	F	Enclosed	Space.	ε res	= 0	.10
Opening Factor	$A \cdot \sqrt{H}/A = 0.$	01	m ^{1/2}								

Opening ractor A	"' ¹ t	• • • •										
T	0	.1		0.3		0.	5		1.0		2.0	
q	1.	•5		4.5		7.	5	-	15.0		30.0	
Time h	T	e	m	g	е	r	a	t	u	r	е .	
0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.40 3.60 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.20 5.40 5.60 5.80 6.00	14 2 2 2 1 1 1 1	5187188943415099887759988		252 418 459 459 459 459 459 459 459 459 459 459		25 44 45 51 54 40 36 33 296 20 18 17 19 6 4 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3	86968870624729862507418525473		250 413 498 498 498 443 555 555 555 555 566 489 441 475 444 4375 445 447 447 550 450 460 460 460 460 460 460 460 460 460 46		250308344575555555555555555555555555555555555	

<u>F2</u>

Time Graphs of Temperature	of	Combustion	Gases.	Type I	F Enclosed	Space.	Eres	= (0.10
Opening Factor $A \cdot \sqrt{H}/A_{r} = 0$.	02 :	_m 1/2					100		

Т	0.1		C	.3		0.5	5]	1.0		2.0
Q	3.0		9	0.0		15.0)	30	0.0		60.0
Time h	T 6	9	m	Þ	е	r	8.	t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.50 3.60 3.80 4.00 5.00	382 612 413 354 273 161 143 120 104 97 90 77 63 48 32		66666666666666666666666666666666666666	33 51 58 54 30 33 36 36		382 612 645 645 645 645 7122 732 734 745 745 745 745 745 745 748 748 748 748 748 748 748 748 748 748		5 5 6 6 6 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7	92 52 53 57 50 50 50 50 50 50 50 50 50 50 50 50 50		382 508 508 507 706 706 707 707 708 708 708 708 708 708 708 708

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Time Graphs of Temperature	of Combustion	Gases.	Type F	Enclosed	Space.	ε _{res} =	0.10
Opening Factor $A \cdot \sqrt{H}/A_{\perp} = 0$.	_{04 m} 1/2					•	

opening radiot in this	"t											
T	0.1	L	(0.3		0.5	i]	L.O		2.0	
q	6.0)	18	3.0		30.0	}	60	0.0]	L20.0	
Time	T	е	m	p	е	r	a	t	u	r	е	
h												
0.05	522			22		522 815			22 73		522 674	
0.10 0.15	815 525			15 41		841			13 63		763	
0.20	437			70		870		8	74		874	
0.25	327			92		892			84		884	
0.30	206			41		901			96 05		897	
0.35 0.40	181 163			67 79		909 916			05 13		905 913	
0.45	150			74		921			19		919	
0.50	139			21		855		9	24		924	
0.60	121			14		731			36		936	
0.65	112			59		661 585			39 41		939 941	
0.70 0.80	105 89			02 64		525			դկ		941	
0.90	72			34		469		8	77		948	
1.00	54			05		413			24		951	
1.10	33			78		356			71		953	
1.20 1.30				49 19		296 272			13 51		9 5 5 9 5 7	
1.40				88		250		5	83		959	
1.50				54		229			48		960	
1.60				<u>4</u> 4		208 188		5. 1.	14 79		962 931	
1.70 1.80				39 36		168			19 43		917	
1.90				34		147			07		903	
2.00				32		124			70		888	
2.20						79 44			94 68		851 811	
2.40 2.60						37			00 414		765	
2.80						34		2:	21		720	
3.00						32		19	98		672	
3.20								1.	76 50		627	
3.40 3.60									52 28		581 533	
3.80									02		483	
4.00								•	76		432	
4.20									46		379	
4.40									38 35		325 284	
4.60 4.80								:	33		261	
5.00									32		238	
5.20											214	
5.40											192	
5.60 5.80											169 146	
6.00											120	

<u>F4</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.10 Opening Factor A· $\sqrt{\rm H}/\rm A_t$ = 0.08 m^{1/2}

	U										
Ţ	0.3	L.	(3.0		0.5	5		1.0		2.0
q	12.0)	36	0.0		60.0)	120	0.0	2	240.0
Time	T	е	m	Þ	е	r	a	t	u	r	е
h											
0.05 0.10	628 971			28 71		628			28		628
0.15	620		10			971 1012		. 0.	23 18		822 917
0.20	493		103			1032		102		-	L025
0.25	355		101			1043		10			L037
0.30	208		97			1052		10)	+7]	047
0.35	188		87			1058		10			.055
0.40 0.45	172		77			1062		106			.060
0.50	160 150		61 58			1066 973		106			.065
0.60	131		45			822		106			.068 .072
0.65	122		36			740		10			.075
0.70	113		31			648		107			.077
0.80	96		28			581		107	9		.079
0.90 1.00	77		25			516		99			.081
1.10	55 30		21 18			451 384		93 87			083 084
1.20	٥٦		15			315		80			086
1.30			12			290		73			087
1.40			8			267		65	0	1	088
1.50			4			244		60			089
1.60 1.70			3			221		56	9		090
1.80			3			198 176		52 48			048 032
1.90				•		153		44			014
2.00						128		40	1		996
2.20						75		31			952
2.40 2.60						36		28			905
2.80						31		26.			857
3.00								23 21			806 750
3.20								18			597
3.40								16	1		5¥2
3.60								13			584
3.80 4.00								10			527
4.20								76 31			₊ 68 ₊07
4.40								32			346
4.60								30			301
4.80											76
5.00										2	251
5.20 5.40											226
5.60											201
5.80											.76 .50
6.00											.22
										_	

<u>F5</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.10 Opening Factor A·VH/A_t = 0.12 m^{1/2}

1 0	T t										
T	0.1	L	0	3		0.5	5	:	1.0		2.0
q	18.0)	51	0.4		90.0)	180	0.0	:	360.0
Time h	T	е	m	p	e	r	a	t	u,	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.00 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	677 1060 657 514 364 205 189 175 164 135 125 98 78 55 27		3	60 66 66 64 68 68 68 68 68 68 68 68 68 68 68 68 68		677 1060 1086 1114 1112 1128 1130 1021 859 771 603 321 297 225 202 179 128 73		86 94 100 111 112 113 114 113 114 115 116 116 117 117 117 117 117 117 117 117	10 117 227 29 31 38 40 46 51 51 51 51 51 51 51 51 51 51 51 51 51		677 880 1097 1112 1112 1112 1112 1113 1114 11

<u>F6</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.35 Opening Factor A· $\sqrt{\rm H}/\rm A_{\rm t}$ = 0.01 m^{1/2}

T	0.1		C	.3		0.5	;	J	L.O		2.0
q	1.5		4	•5		7.5	i	15	5.0		30.0
Time h	T	е	m	р	е	r	а	t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.60 2.80 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	241 387 283 243 193 115 116 102 93 85 73 68 63 54 44 34		41 44 44 44 44 44 45 45 45 45 45 45 45 45	37 19 17 55 30 51 51 63 33 81 13 86 12 12 12		241 387 419 447 451 461 461 461 471 451 461 461 461 471 451 461 461 461 461 461 461 461 461 461 46		33 34 44 44 44 44 44 44 44 44 44 44 44 4	335 355 355 355 355 357 366 366 366 366 366 366 366 366		24729 4471380 4471380 4489935780 4471380 448935780 4471380 449990 4470244 43333222311674 11211086

<u>F7</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.35 Opening Factor A· $\sqrt{\rm H}/\rm A_t$ = 0.02 m^{1/2}

=	ь				
Ţ	0.1	0.3	0.5	1.0	2.0
q	3.0	9.0	15.0	30.0	60.0
Time h	T e	m p	e r a	t u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.40 4.60 4.80 5.00 5.60 6.00	364 560 394 332 257 172 147 131 110 97 984 72 584 29	364 560 588 614 581 475 581 415 380 312 275 236 207 185 143 39 49 40 35 30	364 560 588 615 620 625 630 590 513 470 385 315 214 199 189 153 121 69 41 35 32	364 472 537 6114 6618 6632 6632 6633 6633 6633 6633 6633 663	364 4737114861889913470 6636666666666666666666666666666666666

Time Graphs of	Temperature of	f Combustion	Gases.	Type F	Enclosed	Space.	€	= 0.35
Opening Factor	$A \cdot \sqrt{H}/A_{\perp} = 0.6$	յկ ոլ/2					res	

		•										
T		0.1		(3,0		0.5	;	J	.0		2.0
q		6.0	1	19	0.8		30.0	1	60	0.0	:	120.0
Time		T	e	m	р	е	r	8.	t	u	r	e
h												
0.05 0.10 0.15 0.20 0.35 0.40 0.45 0.50 0.65 0.70 0.65 0.70 1.20 1.30 1.40 1.50 1.70 1.80 1.90 2.20 2.80 3.36 3.36 4.60 4.60 5.50 6.50		473 751 501 407 305 191 168 152 115 107 100 85 50 30		7 7 7 7 7 7 7 7 7 7 7 6 6 5 5 5 4 4 4 4 4 4 4 4 4 4 4 4 4 4 4	751 63 779 807 778 807 778 807 778 807 778 807 778 807 778 807 807		473 751 763 770 783 796 796 793 737 634 426 254 426 254 255 197 139 34 31		626 657 77 77 77 77 77 77 77 77 88 80 80 80 80 80 80 80 80 80 80 80 80	98 98 98 98 99 99 99 99 99 99 99 99 99 9		478 698 48 778 7799 935 68 813 815 66 813 778 66 62 78 84 85 778 80 86 81 81 81 81 81 81 81 81 81 81 81 81 81
0.00												

<u>F9</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.35 Opening Factor A· $\sqrt{\rm H}/\rm A_t$ = 0.08 m^{1/2}

T	0.1		0	•3		0.5	;	1	.0		2.0
q.	12.0		36	.0		60.0		120	.0	;	240.0
Time h	T	е	m	р	e	r	a	t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.80 4.00 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	593 892 593 464 338 198 1128 1128 119 128 128 128 128 128 128 128 128 128 128		4 3	92 21 33 37 98 39 88 86 22 23 23 22		593 892 921 923 934 934 943 860 738 868 590 258 236 214 171 148 124 72 33		77 84 92 92 93 91 91 91 91 91 91 91 91 91 91 91 91 91	236336924478944465775740480227395170		578 999999999999999999999988817766555444833296496269 978 99936924 78 94955544 57 944 06 14 39304664 92184 96 26 9

<u>F10</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.35 Opening Factor A· $\sqrt{\rm H}/\rm A_{\rm t}$ = 0.12 m^{1/2}

	-							_	_		
T	0.1			.3		0.5			0		2.0
q	18.0		54	.0		90.0	•	180	0.0		360.0
Time	Ţ	e	m	p	е	r	a	t	u	r	е
h											
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.40 3.60 3.80 4.00 4.20 4.40 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	640 985 631 489 350 197 182 170 160 150 132 123 114 96 77 53 26		3	35 30 30 34 46 86 37 32 51 36 36 37 55)))	640 985 1000 1013 1016 1018 800 721 633 572 509 311 288 266 243 220 198 1151 125 71 30		899 922 1000 1001 1001 1002 1002 1002 100	21 26 26 27 27 28 28 29 29 29 29 29 29 29 29 29 29 29 29 29		640 852 921 1006 1005 1012 1015 1017 1021 1022 1023 1024 1025 1027 1030 1030 1030 1030 1030 1030 1030 1030 1030 104 979 964 948 909 868 824 777 726 678 678 678 678 678 678 678 67

<u>F11</u>

Time Graphs of Opening Factor	Temperature of	f Combust	ion Gases.	Type F E	nclosed Spac	e. _{Eres} = 0.60
T		0.3	0.5	1.0	2.0	
	0.1 1.5	4.5	7.5	15.0	30.0	
Q Time h	т е	m p	e r a	t u	r e	
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.00 3.40 3.60 3.60 3.60 3.60 3.60 4.20 4.40 4.60 4.80 5.00 5.60 5.80 6.00	231 364 269 229 180 125 105 93 85 78 68 63 59 49 40 31	231 364 405 413 419 408 362 317 288 263 217 193 167 146 130 115 30	231 364 405 413 425 425 425 427 402 351 323 292 265 242 217 192 165 129 118 108 75 32	231 359 416 419 428 428 433 431 431 390 278 433 431 431 390 278 218 201 152 140 78 50 34 50 115 50 78 50 50 50 50 50 50 50 50 50 50 50 50 50	231 311 311 311 311 311 310 410 412 412 412 413 413 414 414 414 414 414 414 414 414	

<u>F12</u>

Time Graphs of	Temperature	of	Combustion	Gases.	Type :	F Enclosed	Space.	٤,,,,	=	0.60
Opening Factor	$A \cdot \sqrt{H}/A_{\perp} = 0$.	02	_m 1/2					160		

	_	τ										
T		0.1		C	.3		0.5	i	1	0		2.0
q		3.0		9	.0		15.0	I	30	0.0		60.0
Time		T	е	m	p	е	r	a	t	u	r	ę
h 0.05 0.10 0.15 0.20 0.35 0.45 0.60 0.65 0.60 0.65 0.60 0.65 0.60 0.70 0.80 0.10 0.10 0.10 0.10 0.10 0.10 0.20 0.2		348 525 374 310 240 159 136 121 111 104 85 80 68 55 41 27		505 56 55 55 55 55 55 55 55 55 55 55 55 55	33 78 48 37 54 18 93 73 73		348 525 547 560 560 571 570 534 426 385 322 289 254 426 172 158 1128 97 31		144 565 566 577 577 575 575 575 575 544 444 336 336 242 211 111 111 111 111 111 111 111 111	71 73 74 77 79 83 83 83 83 83 83 83 83 83 83 83 83 83		34555555555555555555555555555555555555

<u>F13</u>

Time Graphs of Temporal Opening Factor $A \cdot \sqrt{H}$	erature o: /A _t = 0.01	f Combustion	Gases.	Type F E	nclosed Space	• ε _{res} = 0.60
Т	0.1	0.3	0.5	1.0	2.0	
q	6.0	18.0	30.0	60.0	120.0	
Time h	T e	m p e	r a	t u	r e	
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.40 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.20 5.40 5.80 6.00 5.20 5.20 5.40 6.50 6.60	450 708 470 381 288 180 158 144 134 125 110 97 82 66 48 28	450 708 709 710 669 613 537 467 431 352 308 261 233 208 185 161 135 107 78 44 35 31	450 708 703 709 710 714 717 719 722 668 578 474 437 355 309 261 243 206 189 171 152 133 112 69 36 31	450 597 650 720 711 716 718 728 729 731 732 686 652 458 433 407 381 353 326 220 201 181 160 117 94 99 39 32 30	450 597 650 720 711 716 718 723 728 729 731 732 733 741 732 739 740 739 740 698 690 664 637 608 578 514 482 410 372 234 410 372 234 417 55 133	

<u>F14</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.60 Opening Factor A·VH/A_t = 0.08 m^{1/2}

	u										
T	0.1		0	•3		0.5			0		2.0
q.	12.0		36	.0		60.0		120	0.0		240.0
Time h	T	e	m	р	е	r	а	t	u	r	е
0.05 0.10 0.15 0.25 0.35 0.45 0.65 0.65 0.65 0.780 0.65 0.780 0.90 1.10 1.30 1.40 1.50 1.40 2.20 2.40 2.40 2.30 3.30 3.40 3.50 3.40 3.50 3.60	568 840 560 440 190 172 160 150 142 125 1168 92 73 51 27		3	0 24 0 22 0 6 9 8 3 4 3 5 8 8 8 8 8		558 845 862 8667 862 8667 862 8667 862 862 863 863 863 863 863 863 863 863 863 863		71 79 85 85	56 5524 668 0234 56376 1021 0727 196 3073701		561 77855 868 868 8778 878 878 878 878 87

<u>F15</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space. $\epsilon_{\rm res}$ = 0.60 Opening Factor A· $\sqrt{\rm H}/\rm A_{\rm t}$ = 0.12 m^{1/2}

T	0.1		0	•3		0.5		1	.0		2.0
q .	18.0		54	.0		90.0		180	.0		360.0
Time h	T	е	m	р	е	r	a	t	u	r	e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	913 931 599 468 339 192 177 148 130 121 112 95 76 52		3	31 33 33 33 33 33 33 33 33 33 33 34 34 34		613 931 933 936 945 949 873 752 680 601 491 432 370 287 173 149 124 70 28			18 18 18 18 18 18 18 18 18 18 18 18 18 1		68187071167999999999999999999999999999999999

<u>G1</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A+ $\sqrt{H}/\Lambda_{\rm t}$ = 0.01 m^{1/2}

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	T	e m	ũ	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.60 5.60 5.60 5.60 5.60 5.60 5.6	287 473 199 1395 138 122 117 108 102 108 109 109 109 109 109 109 109 109 109 109	287 471 478 456 3924 305 247 197 187 180 141 100 77 166 141 100 141 100 141 141 141 141 141 141	287 478 490 423 350 278 423 251 151 108 88 777 605 528 43 408 33 33 33 33 33 33 33 33 33 33 33 33 33	287184005004702155515649551963068038396555444140987 478465050147021555120222116306807685554443109887 444333332222211163068076855554444140987	287 362 440 498 5007 5117 518 573 574 441 376 413 376 376 376 377 377 377 377 377 377 37	287 361 498 5517 5517 575 575 575 575 575 575 575 5	281 498 550 511 578 661 671 661 671 671 671 671 671 671 671	287 361 498 498 507 517 521 586 631 674 696 671 716 638 633 771 638 633 774 706 833 833 838 839 838 839 838 839 838 839 838 839 838 838

<u>G2</u>

Time Graphs of	Temperature of	Combustion	Gases.	Туре	G Enclosed	Space.
Opening Factor	$A \cdot \sqrt{H}/A_{+} = 0.02$	m ^{1/2}				

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.04 m^{1/2}

	T								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
Q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0	
Time h	Т	e n	ı p	е	r	a t	u	r e	
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.80 6.00	545 842 481 403 297 187 190 161 152 135 71 103 40 38 34 39	546 849 775 654 654 654 651 839 651 855 845 833 833 833 833 833 833 833 833 833 83	546 842 852 853 876 873 876 873 876 873 876 872 873 874 875 875 875 875 875 875 875 875 875 875	545 848 8515 902 733 8515 902 733 8515 902 733 8605 8605 8605 8605 8605 8605 8605 8605	54553453999997764576410693859384913547318754333333333333333333333333333333333333	546 695 763 859 939 938 939 988 988 9716 80 988 988 988 988 988 988 988 988 988	546 6953 8551 9389 9389 9389 9389 9389 9389 9389 938	546 695 763 8551 915 915 915 915 915 915 915 915 915	

<u>G4</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.06 m $^{1/2}$

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	p	е	r	a t	u	r
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.20 3.40 3.60 3.80 4.00 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	593 947 532 439 316 129 124 165 138 129 121 104 86 46 43 37 33 33 30 28	593 947 943 863 7594 965 967 967 967 967 967 967 967 967 967 967	59473593057588217994286636848733040734222226 59435305755821794286635553044334322226	593 947 947 953 9007 1031 998 1031 998 1031 998 1031 998 1031 998 1031 1051 1051 1051 1051 1051 1051 1051	593 780 960 9738 9738 1044 1059 1008 1057 1001 938 816 7637 109 8816 109 8816 109 8816 109 8816 109 8816 8816 8816 8816 8816 8816 8816 881	593 780 960 973 1023 1044 1059 1008 1057 1001 1005 1045 1059 1005 1045 1045 1059 1000 1045 1059 1000 1045 1045 1059 1044 1059 1044 1059 1045 1059 1044 1059 1044 1059 1044 1059 1044 1059 1044 1059 1059 1059 1059 1059 1059 1059 1059	593 780 860 960 973 998 1023 1044 1059 1002 1045 1080 1047 1074 1076	593 780 860 960 973 998 1023 1044 1057 1001 1002 1045 1047 1074 1076

<u>G5</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_t=0.08~m^{1/2}$

_	_							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	Т	e m	р р	е	r s	t t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.60 5.60 6.00	643 1015 574 460 327 188 186 176 157 139 130 121 103 866 41 38 34 32	640 1015 1009 939 802 634 566 491 330 271 255 180 142 101 555 48 34 33 33 33	640 1015 1009 1024 1049 980 797 673 610 484 418 351 207 173 100 60 55 43 39 36 33 31 29	640 1015 1009 1024 1049 1071 1074 1074 1074 1074 1074 1074 1074	641 640 641 640 641 641 641 642 643 644 644 644 644 644 644 644	642 842 918 918 1038 1046 1067	640 842 918 1027 1038 1062 1090 1046 1067 1087 1069 1097 1117	640 842 918 1027 1038 1062 1090 1046 1067 1098 1067 1069 1097 1117

<u>G6</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$ = 0.12 m^{1/2}

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time h	T	e m	ă	е	r a	ı t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.70 1.80 1.90 2.00 2.40 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.40 4.60 4.80 5.00 5.60 5.60 6.00	726 1099 615 486 340 188 187 177 168 159 140 131 103 84 63 35 33 31 29 28 27 26	723 1099 1091 1009 853 669 593 508 422 233 274 255 216 140 37 41 40 37 35 33 29 27 25	723 1099 1091 1104 1125 1048 900 796 670 610 477 407 336 301 267 233 200 167 130 92 46 44 40 37 35 33 30 28 26	723 1099 1091 1104 1118 1134 1153 1052 856 772 674 614 547 405 333 308 284 260 236 211 187 163 136 79 39 34 31 28	726 907 988 1105 1117 11136 11147 1120 11147 11085 960 869 753 662 605 559 442 386 329 315 128 235 208 315 126 336 337 337 337 337 337 337 337 337 33	723 907 988 1104 1117 1136 1157 1120 1141 1158 1131 1147 1148 1071	723 907 988 1104 1117 1136 1157 1120 1141 1158 1131 1147 1163 1148 1169	723 907 988 11.04 11.17 11.36 11.57 11.20 11.41 11.58 11.31 11.47 11.63 11.48 11.69

