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Fredriksson, K; Galle, B; Nystrom, K; Svanberg, Sune

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Lidar system applied in atmospheric pollution monitoring

K. Fredriksson, B. Galle, K. Nyström, and S. Svanberg

A lidar system, incorporating tunable dye lasers and a 25-cm diam Newtonian telescope, has been constructed and applied in atmospheric pollution monitoring. The system, which is fully controlled by a specially designed microcomputer, has been used in several field tests, where stack effluents as well as the ambient air have been monitored. Results from particle, NO₂, and SO₂ measurements are discussed.

Introduction

Remote sensing methods are becoming increasingly important for monitoring the state of the earth’s atmosphere. By using laser techniques, atmospheric pollution can be studied in long-path absorption or atmospheric backscattering experiments. Of the many reviews that have been written on this topic, we would like especially to mention the 1976 monograph edited by Hinkley.¹ The laser radar or lidar approach offers particular advantages because of its single-ended nature and great flexibility. Pulsed laser light is transmitted into the atmosphere, and light backscattered from particles and aerosols is collected by an optical telescope and analyzed. Apart from particle content information, knowledge of the 3-D distribution of certain molecular pollutants can be obtained using the so-called dial (differential absorption lidar) technique. Here the particle backscattering is measured for two wavelengths, where the gas to be investigated absorbs strongly and weakly, respectively. This technique was pioneered by Schotland,² and early laser work was performed by Rothe et al.,³ Grant et al.,⁴ and others. In 1975 we started a program of atmospheric laser probing at this department. In this paper we describe the lidar system that we constructed and discuss results from several field tests, where particles, NO₂, and SO₂ were studied. Our system incorporates tunable dye lasers and uses a Newtonian telescope. A specially designed microcomputer fully controls the measurement procedure.

For a lidar system, information from a distance \( R \) is obtained at a time delay \( t = 2R/c \) after the transmission of the laser pulse, traveling with the velocity of light c. The power \( P_x(R, \Delta R) \) received from the interval \( \Delta R \) at the range \( R \) is given by the general lidar equation:

\[
P_x(R, \Delta R) = \frac{CW_\lambda \sigma_b N(R)}{R^2} \exp \left\{ -2 \int_0^R [\sigma(\lambda)n(r) + \sigma_a N(r)] dr \right\},
\]

where \( W \) is the transmitted pulse power, and \( N(R) \) is the concentration of scatterers with the backscattering cross section \( \sigma_b \). The exponential describes the attenuation of the beam and the backscattered light. The attenuation is caused by absorption due to molecules of concentration \( n(r) \) and absorption cross section \( \sigma(\lambda) \) and scattering and absorption due to particles of concentration \( N(r) \) and cross section \( \sigma_a \). \( C \) is a system constant. In particle studies it is sometimes a reasonable approximation to neglect the absorption term. If, then, \( \sigma_b \) is assumed to be a constant, \( N(R) \) can be directly mapped in three dimensions. As the value of \( \sigma_b \) is not generally known, only relative concentration distributions are usually obtained. In dial measurements of gas concentrations, use is made of the fact that wavelength regions can be found where \( \sigma(\lambda) \) strongly changes for a small wavelength change, while \( \sigma_b \) and \( \sigma_a \) stay approximately constant. In these conditions, the evaluation of the experimental data is quite simple. If more than one gas absorbs in the wavelength region of interest, a summation must be performed in the exponential factor.

Lidar System

The lidar system used in these measurements is schematically shown in Fig. 1. The system includes three different pulsed lasers, a Newtonian telescope, and detection, controlling, and processing electronics.
The choice of laser transmitter is dependent on the measuring technique employed. A 400-kW nitrogen laser with 10-nsec pulses at 337 nm is used in measurements of particles and aerosols. The repetition rate is 10–50 pulses/sec. In NO\textsubscript{2} measurements, a dye laser pumped by the nitrogen laser is used. The dyes Coumarin 45 or Coumarin 47, giving about 0.5 mJ in 5-nsec pulses in the 440–480-nm region, are used to produce light of suitable wavelengths in the NO\textsubscript{2} absorption band. Wavelengths around 300 nm in the UV absorption band of SO\textsubscript{2} are obtained with a frequency-doubled flashlamp pumped dye laser as 1-μsec pulses with the energy 0.4 mJ at a repetition rate of 10–25 Hz. The wavelength setting of the dye lasers is made by stepper motors.

The Newtonian telescope has a 25-cm mirror with a focal length of 100 cm and is constructed with a mount for vertical and horizontal turnings. The turnings are effected with computer-controlled motors. The laser beam is transmitted to an output mirror (placed inside the telescope) by means of first-surface aluminum mirrors in such a way that the transmitted laser beam and the optical axis of the telescope are almost coaxial at every positioning of the system. Alternatively, the flashlamp pumped dye laser can be placed on a table rotating together with the telescope. The field of view of the telescope normally corresponds to 2 mrad to match the divergence of the dye lasers. Interference filters, 5–10 nm broad with a transmission of 35–55%, are used to suppress the background skylight. Another variation is that the system may be used with a 500-mm Bausch & Lomb grating monochromator.

The spectrally filtered lidar signal is detected by an EMI 9558 QB PMT. A fast transient digitizer, Biomat 8100 with 2048 8-bit channels, is used to convert the signal to time-resolved digital data. The data are transferred to a microcomputer for signal averaging.

This unit has a memory consisting of 4096 24-bit words that may be divided into sixteen or fewer parts. The maximum transient frequency is 400 Hz. The number of transients to be averaged, the laser wavelength setting, the telescope steering, and the readout on a strip-chart recorder are controlled by the microcomputer, which also can do arithmetic operations. Thus the measuring cycle can be fully automatized. The measurement process is monitored on oscilloscopes via continuous readouts of the individual and the averaged lidar signals. In some lidar measurements a boxcar averager is used instead of the transient digitizer. In this case the lidar signals from only one or two distances are studied at each laser pulse. In some measurement modes a chopper is used to block every second or ninth shot from the laser. Then rf interference, caused by the nitrogen laser, and the background signal level are subtracted from the lidar signals. The lidar system is described in more detail in an institute report.\textsuperscript{5} In measurements over the city of Goteborg, the system has been placed in a suitably located laboratory within the Physics Department; during field tests it was housed in a small truck.

**Measurements of Particles**

We have used the lidar system in several studies of particle emissions from industrial smoke stacks. Measurements of relative particle distributions are easy to perform using elastically backscattered light and neglecting weak effects of beam attenuation. To obtain good spatial resolution it is important to use a laser of short pulse length. Normally we have used the 10-nsec UV pulses from the N\textsubscript{2} laser. In Fig. 2 the result of a vertical scan downwind an industrial area is shown.

![Fig. 2. Relative particle distribution downwind an industrial area, obtained in a vertical lidar scan employing a nitrogen laser. An experimental curve illustrating the evaluation procedure is included.](image-url)
measurements was performed using the boxcar averager, the plume (light backscattered from the air before and behind the beam attenuation by the plume, measured in the procedure. For dense plumes it is preferable to measure the extent of the plume echo can be used in the evaluation of the particle concentration. If the particle concentration is not too high, the amplification normally necessary to provide an in-stack calibration. If the particle concentration is not too high, the amplification is normally necessary to provide an in-stack calibration. This approach avoids both influences due to wind and due to condensing water droplets. Because of the complexity of the Mie scattering theory and the lack of detailed information on particle characteristics, it is normally necessary to provide an in-stack calibration. If the particle concentration is not too high, the amplification of the plume echo can be used in the evaluation procedure. For dense plumes it is preferable to measure the beam attenuation by the plume, measured in the light backscattered from the air before and behind the plume ("plume opacity").

In Fig. 3 a particle measurement is illustrated. The remote sensing curve displays the amplitude of the plume echo vs time. The measurement was performed using the boxcar averager, set to sample the intensity at a delay time corresponding to the light round-trip time to the plume, in the actual case situated at 220-m distance from the lidar system. Comparison is made with the in-stack curve obtained with a calibrated recording instrument. The parts lacking on each curve correspond to zero calibration intervals. A good correlation between the curves was found except during the situations when the in-stack instrument was out of range. From the curves the proper calibration for the particular type of smoke (from an iron-alloy melting furnace) could be established.

**Measurements of NO₂**

In measurements of gaseous pollutants, the dial technique is normally employed. As mentioned above, this technique is based on the difference in absorption cross section of a certain molecular gas at two or more close lying wavelengths. At normal atmospheric conditions the backscattering and the absorption of particles and aerosols are slowly varying functions of the wavelength. Thus, a fast wavelength dependence in the lidar Eq. (1) will be solely due to absorbing gases. The system constant C may also have a fast wavelength dependence, e.g., from Fabry-Perot effects in the transmitting and receiving optics, but this is automatically taken care of by the evaluation procedure. The dial equation is obtained from Eq. (1) as

\[
\frac{P_{\lambda 1}(R)}{P_{\lambda 2}(R)} = C' \cdot \exp\left\{-2 \int_0^R \left[\sigma(\lambda_1) - \sigma(\lambda_2)\right] n(r) dr\right\},
\]

where C' is a new constant, and it is assumed that only one atmospheric gas has a differential absorption at the two wavelengths \(\lambda_1\) and \(\lambda_2\). If \(\sigma(\lambda_1)\) is larger than \(\sigma(\lambda_2)\), it follows from Eq. (2) that the dial signal will decrease with the distance R in the presence of the studied gas.

We have measured NO₂ concentrations with the dial technique in the 440–460-nm wavelength region. The cross-section values used in the calculations are those given by Grant et al., with wavelength extensions made in absorption measurements in our laboratory fitted to these values. Deviations from these, as obtained in laboratory measurements by Woods and Jolliffe, are within the uncertainty limits. Possible corrections at high temperatures in smokestack plumes have not been considered in these measurements.

Our first measurements on gaseous pollutants concerned NO₂ concentrations in smokestack plumes of an oil refinery. The lidar system was located at a distance of 2 km from the stack studied. In Fig. 4 a measurement is shown in which the lidar was directed just above the top of the smokestack. The backscattered intensity from the atmosphere surrounding the plume, measured for two different wavelengths, is displayed. The lidar signal of the plume is far beyond scale and is cutoff in the figure. This measurement implies an integrated concentration of 190 ± 50 mg/m³ · m, which equals 30 ± 10 mg/m³ as an average in the plume. The uncer-
30 mg/m$^3$. Measurements of exhausts from an aeroengine manufacturer during tests of jet engines implied very high NO$_2$ contents; values up to 900 mg/m$^3$. were obtained. These NO$_2$ measurements were performed at a range of 1 km. NO$_2$ concentrations in the ambient air have been measured in different types of areas. Tens of parts per billion have been recorded in the city of Goteborg. When average concentrations over large distances are to be measured, one may use some distant topographic target instead of the atmosphere for the backscattering of the laser beam. Such a measurement is shown in Fig. 5, where a mountain at a distance of 5.5 km is used as the backscattering target. This measurement yields an average concentration of 2 ± 2-ppb NO$_2$.

**Measurements of SO$_2$**

In the dial measurements of SO$_2$, wavelengths in the absorption band around 300 nm are used. The absorption cross-section values used in our evaluations are those given by Thompson et al. $^9$ Corrections to these values at high temperatures, as discussed by Jolliffe and Wood$^10$ and Exton, $^11$ are made in measurements of smokestack plumes. Wavelength calibration was made using absorption cells containing SO$_2$.

In addition to dial measurements at two wavelengths, a modified type of dial technique was applied when measuring high SO$_2$ concentrations in smokestack plumes. Then the attenuated lidar signal as measured through the plume is compared with the signal obtained in a direction beside the plume. Only one laser wavelength, with rather low SO$_2$ cross section, is used, and thus the method is useful only when the attenuation caused by particles is negligible. This technique is preferable at high concentrations, since the laser light at the strongly absorbed wavelength is then almost completely blocked by the gas. Alternatively, one may use the ordinary dial technique in the region of weaker absorption around 315 nm.

Measurements of SO$_2$ contents in plumes from an oil-burning power station were made. Concentrations in the 300–550-ppm range were obtained using the modified dial technique. In Fig. 6 such a measurement of SO$_2$ is shown. The registration curve displays the result obtained from a division of the lidar recordings in a direction through the plume and beside it, respectively. Compared with Eq. (2), this equals a situation where $\sigma(\lambda_2)$ is zero. The 1-μsec pulse length of the flashlamp pumped dye laser reduces the spatial resolution; therefore, the lidar plume echo in Fig. 6 is rather broad. Measurements of outlets from an iron-alloy plant and from other pollution sources were also made as well as studies of the spreading of SO$_2$ plumes.

We have also performed measurements of ambient SO$_2$ concentrations in urban as well as in industrial areas. Figure 7 shows a measurement of the lidar signal from the atmosphere at a distance of 1 km at six different wavelengths. Each value has been normalized to the short-range atmospheric lidar signal. The values
Fig. 6. Lidar measurement of SO₂ emission from an oil-burning power station. Two lidar curves, measured through and beside the smoke plume, have been divided, displaying the smoke-particle scattering as well as the attenuation due to SO₂. A SO₂ concentration of 760 ± 150 mg/m³ was inferred across the 4.5-m diam plume.

show the time of the individual measurements. Comparison is made with the day average value obtained by the local Public Health Board with a conventional chemical measuring method. As can be seen, the remote sensing results are fully compatible with the reference value.

Conclusions

The lidar system described in this paper has proved efficient and useful for atmospheric probing. Within its basic constraints it has been gradually refined as a result of the experience gained in an extended field-test program. A more detailed account of the different types of measurements discussed is given in Refs. 5 and 12-15. Based on this work, we are now developing a fully mobile dedicated lidar facility for the Swedish Space Corporation. The system, housed in a 6-m covered truck, is based on a Nd:YAG pumped dye laser and agree with the absorption feature of SO₂, and an average concentration of 25 ± 8 ppb is implied from this measurement.

A typical two-wavelength dial measurement of SO₂ in the ambient air is shown in Fig. 8. The upper curve displays the dial signal according to Eq. (2), and the vertical scale expresses the SO₂ concentration as integrated from 500 m. The lower curve shows the concentration obtained when the dial signal is evaluated over 500-m intervals. This measurement was made in an area where the background level is rather low and where the atmosphere up to a distance of 2 km is polluted from an industrial area.

Several studies of the SO₂ pollution situation in Goteborg were made. A measurement of the time-varying concentration over a period of 7 h is shown in Fig. 9. Evaluations of the backscattered return from the atmosphere at 2 km as well as from a topographic target at a distance of 3 km are shown in the figure. The vertical bars in the figure indicate the estimated maximal errors in the measurements; the horizontal bars show the time of the individual measurements. Comparison is made with the day average value obtained by the local Public Health Board with a conventional chemical measuring method. As can be seen, the remote sensing results are fully compatible with the reference value.

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a fixed vertical 30-cm telescope. A beam-steering plane mirror plus the transmitter and receiver electronics are controlled by a minicomputer. The mobile system will be operated by the Swedish Environmental Protection Board for monitoring particles, NO₂, SO₂, and O₃. Future extensions to CO, CH₄, and HCl are foreseen. A system of this kind is also useful for hydrospheric probing as has been demonstrated, e.g., in this laboratory. Thus the lidar concept has important implications to environmental diagnostics.

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