

## Scattering Matrix Formulation for Substructure Characteristic Mode Analysis

Mats Gustafsson, Johan Lundgren, Kurt Schab, Lukas Jelinek, Miloslav Capek

Electrical and Information Technology, Lund University, Sweden Dept. of Electrical Engineering, Santa Clara University, Santa Clara, USA Dept. of Electromagnetic Field, Czech Technical University, Prague, Czech Republic

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 $\begin{array}{l} \mathsf{M}. \mbox{ Gustafsson et al. "Unified theory of characteristic modes: Part I–Fundamentals". IEEE Trans. Antennas Propag. \\ 70.12 (2022), pp. 11801–11813 \\ \mathsf{M}. \mbox{ Capek et al. "Characteristic Mode Decomposition Using the Scattering Dyadic in Arbitrary Full-Wave Solvers". \\ IEEE Trans. Antennas Propag. 71.1 (2023), pp. 830–839 \\ \end{array}$ 



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 Modal significance |t<sub>n</sub>| = 1/(√1+λ<sub>n</sub><sup>2</sup>)



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► characteristic numbers  $\lambda_n = -\frac{\operatorname{Im}\{t_n\}}{\operatorname{Re}\{t_n\}}$ ,  $\phi_n = 2 \arctan(\lambda_n^{-1})$ ,  $\alpha_n = \pi - \arctan(\lambda_n)$ 

• Modal significance 
$$|t_n| = \frac{1}{\sqrt{1+\lambda_n^2}}$$

Many equivalent representations for CM

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- $\blacktriangleright$  Excitation  $\mathbf{V}_c$  in the antenna region and induced currents  $\mathbf{I}_0$  on the ground plane

$$\begin{bmatrix} \mathbf{Z}_{cc} & \mathbf{Z}_{c0} \\ \mathbf{Z}_{0c} & \mathbf{Z}_{00} \end{bmatrix} \begin{bmatrix} \mathbf{I}_0 \\ \mathbf{I}_c \end{bmatrix} = \begin{bmatrix} \mathbf{0} \\ \mathbf{V}_c \end{bmatrix}$$



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 $\begin{array}{c} \Omega_{c} \\ V_{c} \\ I_{c} \end{array} \qquad \begin{array}{c} \Omega_{0} \\ V_{0} = 0 \\ I_{0} \end{array}$ Note: decompositions based on basis function (DoF) can require a fine mesh to model connected regions

 $\Omega_{\rm G}$ 

 $\Omega_{\rm A}$ 

Reduced MoM system matrix

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J. Ethier and D. McNamara. "Sub-structure characteristic mode concept for antenna shape synthesis". Electronics letters 48.9 (2012), p. 1

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 Substructure CM from the generalized eigenvalue problem [EM12]

$$\widetilde{\mathbf{X}}\widetilde{\mathbf{I}}_n = \lambda_n \widetilde{\mathbf{R}}\widetilde{\mathbf{I}}_n \quad \text{or } \widetilde{\mathbf{Z}}\widetilde{\mathbf{I}}_n = (1 + j\lambda_n)\widetilde{\mathbf{R}}\widetilde{\mathbf{I}}_n$$



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# What is the corresponding scattering formulation for substructure CM?



Scattering matrices



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► The scattering eigenvalues s<sub>n</sub> are related to modal significance |t<sub>n</sub>| and MoM substructure characteristic eigenvalues λ<sub>n</sub> = eig(X̃, R̃) as

$$t_n = rac{s_n - 1}{2} \quad ext{and} \ \lambda_n = -\operatorname{Im}\{t_n^{-1}\} = \mathrm{j}rac{s_n + 1}{s_n - 1}$$



Scaling of the scattered field from the background





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- Unitary scattering matrices (lossless)

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► The scattered fields can be expressed in transition matrices (or scattering dyadics) S = 2T + 1 and S<sub>0</sub> = 2T<sub>0</sub> + 1

$$\frac{1}{2}(\mathbf{S}_{0}^{\mathrm{H}}\mathbf{S}-\mathbf{1})\mathbf{a}_{n} = (2\mathbf{T}_{0}^{\mathrm{H}}\mathbf{T}+\mathbf{T}_{0}^{\mathrm{H}}+\mathbf{T})\mathbf{a}_{n} = t_{n}\mathbf{a}_{n}$$



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Maximal scattering (difference between the scattered fields of the composite object Ta<sub>n</sub> and the background T<sub>0</sub>a<sub>n</sub>)

$$|(\mathbf{T} - \mathbf{T}_0)\mathbf{a}_n|^2 = \mathbf{a}_n^{\mathrm{H}}(\mathbf{T}_0^{\mathrm{H}}\mathbf{T}_0 + \mathbf{T}^{\mathrm{H}}\mathbf{T} - 2\operatorname{Re}\{\mathbf{T}_0^{\mathrm{H}}\mathbf{T}\})\mathbf{a}_n\}$$
  
= - Re{ $\mathbf{a}_n^{\mathrm{H}}(2\mathbf{T}_0^{\mathrm{H}}\mathbf{T} + \mathbf{T}_0^{\mathrm{H}} + \mathbf{T})\mathbf{a}_n\}$  = - Re{ $\{t_n\}|\mathbf{a}_n|^2 = |t_n|^2|\mathbf{a}_n|^2$ 


• factorizing the radiation matrix  $\mathbf{R} = \operatorname{Re}\{\mathbf{Z}\} = \mathbf{U}^{\mathrm{T}}\mathbf{U}$  into spherical waves [Gus+22a] with  $\mathbf{U} = \begin{bmatrix} \mathbf{U}_0 & \mathbf{U}_c \end{bmatrix}$  and setting

$$\widetilde{\mathbf{U}} = \mathbf{U}_{\mathrm{c}} - \mathbf{U}_{0}\mathbf{Z}_{00}^{-1}\mathbf{Z}_{00}$$

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(\*)

► reformulates the CM eigenvalue problem to  $\widetilde{\mathbf{Z}}\widetilde{\mathbf{I}}_n = (1 + j\lambda_n)\widetilde{\mathbf{U}}^{\mathrm{H}}\widetilde{\mathbf{U}}\widetilde{\mathbf{I}}_n \Rightarrow -\widetilde{\mathbf{U}}\widetilde{\mathbf{Z}}^{-1}\widetilde{\mathbf{U}}^{\mathrm{H}}\mathbf{f}_n = t_n\mathbf{f}_n$ where  $\mathbf{f}_n = -\widetilde{\mathbf{U}}\widetilde{\mathbf{I}}_n$  and  $t_n = -1/(1 + j\lambda_n)$ 

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where  $\mathbf{f}_n = -\widetilde{\mathbf{U}}\widetilde{\mathbf{I}}_n$  and  $t_n = -1/(1 + \mathrm{j}\lambda_n)$ 

$$\mathbf{T} = -\mathbf{U}\mathbf{Z}^{-1}\mathbf{U}^{\mathrm{T}} = -\begin{bmatrix} \mathbf{U}_0 & \mathbf{U}_\mathrm{c} \end{bmatrix} \begin{bmatrix} \mathbf{Z}_{00} & \mathbf{Z}_{0\mathrm{c}} \\ \mathbf{Z}_{\mathrm{c}0} & \mathbf{Z}_{\mathrm{cc}} \end{bmatrix}^{-1} \begin{bmatrix} \mathbf{U}_0^{\mathrm{T}} \\ \mathbf{U}_\mathrm{c}^{\mathrm{T}} \end{bmatrix} \quad \mathbf{T}_0 = -\mathbf{U}_0\mathbf{Z}_{00}^{-1}\mathbf{U}_0^{\mathrm{T}}$$

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Transition matrices of the composite object Ω and background object Ω<sub>1</sub> are expressed in MoM system matrices

$$\mathbf{T} = -\mathbf{U}\mathbf{Z}^{-1}\mathbf{U}^{\mathrm{T}} = -\begin{bmatrix} \mathbf{U}_0 & \mathbf{U}_c \end{bmatrix} \begin{bmatrix} \mathbf{Z}_{00} & \mathbf{Z}_{0c} \\ \mathbf{Z}_{c0} & \mathbf{Z}_{cc} \end{bmatrix}^{-1} \begin{bmatrix} \mathbf{U}_0^{\mathrm{T}} \\ \mathbf{U}_c^{\mathrm{T}} \end{bmatrix} \quad \mathbf{T}_0 = -\mathbf{U}_0\mathbf{Z}_{00}^{-1}\mathbf{U}_0^{\mathrm{T}}$$

▶ Block inversion shows that MoM matrix (\*) and scattering based substructure modes  $2\mathbf{T}_0^H\mathbf{T} + \mathbf{T}_0^H + \mathbf{T}$  are equivalent

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### Numerical example



- (left) full structure (right) substructure
- Adapted from [EM12],  $\ell = 120 \,\mathrm{mm}$ ,  $w = 60 \,\mathrm{mm}$ ,  $h = 15 \,\mathrm{mm}$ ,  $d = 30 \,\mathrm{mm}$
- Negligible differences for such as a sheet resistance  $0.01 \Omega/\Box$

# Generalizations

Many possible generalizations:

- Embedded structures
- Cavities
- Combination with ports



# Generalizations



Mats Gustafsson, Lund University, Sweden, 10

#### Scattering based formulation of substructure CM



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 $\Omega_{C}$  $\Omega_{\rm A}$  $\Omega_2$  $\overline{\mathbf{V}_{2}}$  $\mathbf{I}_{2}$ 

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