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1970

Link to publication

*Citation for published version (APA):* Magnusson, S. E., & Thelandersson, S. (1970). *Temperature - Time Curves of Complete Process of Fire Development.* (Bulletin of Division of Structural Mechanics and Concrete Construction, Bulletin 16; Vol. Bulletin 16). Lund Institute of Technology.

*Total number of authors:* 2

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S. E. MAGNUSSON and S. THELANDERSSON

TEMPERATURE-TIME CURVES OF COMPLETE PROCESS OF FIRE DEVELOPMENT

## Ci 65

# ACTA POLYTECHNICA SCANDINAVICA

CIVIL ENGINEERING AND BUILDING CONSTRUCTION SERIES No. 65

# Temperature–Time Curves of Complete Process of Fire Development

Theoretical Study of Wood Fuel Fires in Enclosed Spaces

S. E. MAGNUSSON and S. THELANDERSSON Division of Structural Mechanics and Concrete Construction Lund Institute of Technology Lund, Sweden

STOCKHOLM 1970

Printed in Sweden by Elanders Boktryckeri AB Göteborg 1970

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## Principal notations

A	Area of vertical openings in the en-	
$A_{h}$	closed space Area of horizontal openings in the en-	m <sup>2</sup>
<u>n</u>	closed space	m <sup>2</sup>
$A_t$	Total bounding surface area of the en-	111
-	closed space	m <sup>2</sup>
$(A \cdot \sqrt{H})$	Air flow factor (Ventilation factor)	m <sup>5/2</sup>
$(A \cdot \sqrt{H}/A_t)$	Opening factor	m <sup>1/2</sup>
В	Width of an opening in the enclosed	
	space	m
G	Volume of combustion gases produced	
~	per unit weight of fuel	$m^3 \cdot kg^{-1}$
$G_0$	Volume of combustion gases (expressed	
	in Nm <sup>3</sup> ) produced per unit weight of	
IJ	fuel	$\mathrm{Nm}^3 \cdot \mathrm{kg}^{-1}$
H	Height of vertical opening in the en-	
H'	closed space	m
12	Height of vertical opening below the	
H''	neutral zone level in the enclosed space Height of vertical opening above the	m
~*	neutral zone level in the enclosed space	
Ι	Enthalpy	m kcal · m <sup>−3</sup>
$I_{C}$	Heat energy released per unit time	Ktal · III
U	during combustion	kcal · h <sup>−1</sup>
$I_{\mathcal{B}}$	Heat energy stored per unit time in the	Roux 11
	gas volume which is contained in the	
	enclosed space	kcal · h <sup>−1</sup>
$I_L$	Heat energy withdrawn per unit time	
	from the enclosed space owing to the	
<b>.</b>	replacement of hot gases by cold air	kcal $\cdot$ h <sup>-1</sup>
$I_R$	Heat energy withdrawn per unit time	
	from the enclosed space by radiation	
	through openings in the enclosed space	kcal $\cdot$ h <sup>-1</sup>

7	Heat energy withdrawn per unit time	
$I_W$		
	from the enclosed space through wall,	Front h=1
_	roof or ceiling, and floor structures	kcal $\cdot$ h <sup>-1</sup>
L	Quantity of air consumed per unit	NT 2 1 -1
	weight of fuel during combustion	$Nm^3 \cdot kg^{-1}$
M	Quantity of combustible material	kg
Р	Static pressure	$kg \cdot m^{-2}$
$Q_{ m out}$	Rate of flow of the outgoing gases	
	through a vertical opening in the en-	
	closed space	kg · h <sup>−1</sup>
$Q_{ m in}$	Rate of flow of the incoming air through	
2-1N	a vertical opening in the enclosed space	$kg \cdot h^{-1}$
$Q_h$	Rate of flow of the outgoing gases	5
$\mathcal{Q}_h$	through a horizontal opening in the	
	enclosed space	$kg \cdot h^{-1}$
0	—	K5 II
$\mathcal{Q}$	Rate of flow of air supplied to the en-	$m^3 \cdot s^{-1}$
-	closed space by means of fans	
R	Rate of combustion	kg of wood per unit time
_		•PACE • •
$R_{\max}$	Maximum rate of combustion deter-	kg of wood per
	mined by the rate of air supply	unit time
Т	Duration of the fire defined as the du-	
	ration of the flame phase	h
W	Heat value of the fuel	kcal · kg <sup>−1</sup>
с	Specific heat	kcal $\cdot$ m <sup>-3</sup> $\cdot$ °C <sup>-1</sup>
Cp	Specific heat of the combustion gases	kcal $\cdot$ m <sup>-3</sup> $\cdot$ °C <sup>-1</sup>
g	Acceleration of gravity	$m \cdot s^{-2}$
ĥ	Difference in level between the centre of	
	a vertical opening and a horizontal	
	opening	m
h'	Difference in level between the neutral	
	zone and a horizontal opening	m
а	Fire load	Mcal $\cdot$ m <sup>-2</sup> of bound-
q	ine load	ing surface area
v	Hydraulic radius	cm
r		$m \cdot h^{-1}$
$v_y, v_z$	Velocity of flow	111 11
V <sub>h</sub>	Velocity of flow through a horizontal	m . h-1
	opening in the enclosed space	$m \cdot h^{-1}$
t	Time co-ordinate	h
x	Position co-ordinate	m
α <sub>i</sub>	Coefficient of heat transfer at a surface	
	exposed to fire (internal surface)	kcal $\cdot$ m <sup>-2</sup> $\cdot$ h <sup>-1</sup> $\cdot$ °C <sup>-1</sup>

$\alpha_u$	Coefficient of heat transfer at a surface	
	not exposed to fire (external surface)	kcal $\cdot$ m <sup>-2</sup> $\cdot$ h <sup>-1</sup> $\cdot$ °C <sup>-1</sup>
γ	Weight per unit volume	$kg \cdot m^{-3}$
$\varepsilon_{\rm res}$	Resultant emissivity for radiation be-	kg m
	tween flames, combustion gases, and a	
	surface exposed to fire (internal surface)	
$\varepsilon_{fl}$	Emissivity of flames	
$\varepsilon_i$	Emissivity of a surface exposed to fire	
Q	Temperature	°C
Ъ°	Temperature of the outside air	°C
$\vartheta_g$	Temperature of the combustion gases	°C
$\vartheta_i$	Temperature of a surface exposed to	C
	fire (internal surface)	°C
$\vartheta_u$	Temperature of a surface not exposed	0
	to fire (external surface)	°C
$\Delta \vartheta$	Temperature difference between the	C
	combustion gases and the outside air	°C
λ	Thermal conductivity	kcal $\cdot$ m <sup>-1</sup> $\cdot$ h <sup>-1</sup> $\cdot$ °C <sup>-1</sup>
μ	Coefficient of contraction	
ho	Density	kg ⋅ m <sup>-3</sup>
$ ho_0$	Density of the outside air	$kg \cdot m^{-3}$
$ ho_g$	Density of the combustion gases	$kg \cdot m^{-3}$
	Super Super	ng m

#### 1. Introduction

The efforts made during the past decade in the field of structural fire engineering research have paved the way for differentiated, functionally correct structural fire engineering design carried out on the basis of theoretical calculations. This was rendered possible by investigations which can on the whole be classified in one or several of the main groups enumerated below. At the same time, these groups may be regarded as the essential stages or steps in an appropriate procedure for fire engineering design of load-bearing and separating structures [1].

(a) Determination of the characteristics of the fire load in an enclosed space under exposure to fire.

(b) Study of the variations in the development of energy, in the requisite air supply, and in the evolution of gases, with the time in the course of a fire. Determination of the temperature of the combustion gases in the enclosed space as a function of the time.

(c) Determination of the thermal properties of the materials used for structures in the temperature range which is of interest in connection with fires.

(d) Determination of the non-stationary temperature fields which are produced in a fire-exposed structure on the assumption that the temperature-time curve for the combustion gases is given, cf. (b).

(e) Determination of the structural behaviour and the load-bearing capacity of a fire-exposed structure on the basis of the temperature fields defined under (d), and with the help of the available information on those changes in the strength and deformation characteristics of the materials which take place under such conditions.

The object of the present investigation is to make a close study of the stage (b) in order to determine the complete temperature-time curve for the gaseous products of combustion under different conditions, and in particular the temperature-time curve in the cooling phase,<sup>1</sup>) for fires of the wood fuel type in enclosed spaces.

<sup>&</sup>lt;sup>1</sup>) The characteristics of the different phases of the process of fire development are represented in Fig. 1. The term "cooling phase" will be used in this publication to designate the smoulder phase and the cooling phase taken together.



Fig. 1. Phases of the process of fire development as defined in the present publication. The definition in [1] distinguishes between a smoulder phase and a cooling phase, which are here regarded as a single phase designated by the term "cooling phase".

The present state of research in this field is clearly reflected in the sections dealing with fire protection in the Swedish Building Regulations 1967 (abbreviated SBR 67) and in the Draft Specification "Aluminium Structures". In comparison with the relevant regulations which are in force in most other countries, the Swedish rules represent substantial progress on the road to judicious structural fire engineering design. This is primarily due to the fact, that, when the designer has to choose that temperature-time curve which characterises the process of fire development, and which must serve as a basis for all theoretical structural fire engineering design, these rules enables the designer to be guided by all the results, which have been obtained from research in this field during recent years.

On an international plane, it is found that standard temperature-time curves for the process of fire development have been adopted in several countries [1]. If the fuel supply is unlimited, then the agreement between these curves is relatively close, see Fig. 2 a. Under practical conditions, when the fuel supply is limited, the standard specifications used in various countries stipulate that the variation in the temperature with the time shall be in conformity with the standard curve during a certain definite period of time, which is designated by the term "duration of the fire", T, and is defined as the duration of the flame phase. A comparison of the relations between the duration of the fire and the fire load which are employed in various countries is represented in Fig. 2 b [1]. This comparison shows very great differences in an assumption which is fundamental for structural fire engineering design. The wide dispersion between the curves reproduced in Fig. 2 b indicates that there exists no univalued relation between the fire load and the duration of the fire. Concerning



Fig. 2a. Standard curves used in some countries to represent the variation in the temperature,  $\vartheta$ , in an enclosed space exposed to fire with the time, t. The symbol  $\vartheta_0$  denotes the temperature in the enclosed space at the time t=0.

1. ISO/TC 92; INSTA 28/2; DIN 4102-62.

- 2. EMPA, Switzerland.
- 3. ASTM 119 (1953), USA.

4. V 1076 (1955), Netherlands; BS 476 (1953), United Kingdom.

5. A 1304, Japan.



Fig. 2b. Relations between the fire load, in kg of wood per m<sup>2</sup> of floor area, and the duration of the fire, in h, which are stipulated in the standard specifications used in Sweden (Curve 1), United States of America (Curve 2), and United Kingdom (Curve 3), as well as in a Swiss draft standard specification (Curve 4).

the cooling phase of the process of fire development, it should be noted that it is as a rule completely disregarded.

The relevant Swedish standard specifications provide the designer with three alternative methods of design. Just as in other countries, it is permissible to carry out the design in a roughly simplified and stereotyped manner by using Curve 1 in Fig. 2 a as a point of departure. For the flame phase, this curve gives the temperature of the combustion gases,  $\vartheta_g$ , in the enclosed space in accordance with the equation

$$\vartheta_{e} - \vartheta_{0} = 1325 - 430 \ e^{-0.2t} - 270 \ e^{-1.7t} - 635 \ e^{-19t}$$
 (1.1)

where t is the time, in hours, and  $\vartheta_0$  is the temperature in the enclosed space at the time t=0. The differences in the combustion characteristics of various fuels, or the fact that the rate of combustion varies within wide limits with the dimensions of the openings in the enclosed space, are not taken into account in this equation. The above-mentioned curve is closely in agreement with that temperature-time curve which is recommended by the ISO for fire tests on building components.

Alternatively, for certain definite types of fire loads and enclosed spaces, the designer may use a method which is simplified, but is nevertheless more differentiated, in comparison with the design procedure outlined in the above. The applicability of this alternative method presupposes that it is possible to comply with the two necessary conditions which are stated in what follows. In the first place, it is required to demonstrate that the characteristics of the fire load in respect of rate of combustion and radiation are approximately in accordance with those which apply in the case of wood fuel. In the second place, it is stipulated that the opening factor of the enclosed space, which is given by the expression  $A\sqrt{H}/A_t$ , where A is the total opening area of windows and doors, in  $m^2$ , H is a weighted average of the vertical dimensions of these openings, in m, and  $A_t$  is the total area of the surfaces bounding the enclosed space, in m<sup>2</sup>, shall be known during all phases of the process of fire development. If these two conditions are satisfied, then it is allowed to carry out the design on the basis of a specific curve representing the variation in the temperature of the combustion gases in the enclosed space,  $\vartheta_g$ , with the time. For the flame phase, this curve is determined by the opening factor, see Fig. 2 c, in the course of the duration of the fire, T, which is defined by the equation

$$T = qA_t / (25A) / H) \quad \text{min} \tag{1.2}$$

where q is the fire load, in Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area. The dash-line curve represents the INSTA curve expressed by Eq. (1.1).

Finally, and generally, in the cases where the quantities of combustible materials which constitute the fire load, as well as the rate of combustion, are accurately known, the above-mentioned two Swedish specifications allow



Fig. 2c. Variation in the temperature of the combustion gases,  $\vartheta_g$ , with the time, t, at different values of the opening factor,  $A \cdot \sqrt[3]{H}/A_t$ . Curves published in the Swedish Building Regulations 1967 (SBR 67) and in the Swedish draft specification "Aluminium Structures". The dash-line curve is the standard curve calculated by means of Eq. (1.1).

the fire resistance of a building component to be determined on the basis of the variation in the temperature of the combustion gases with the time, calculated with the help of the formula which is known as the equation of heat balance. This equation, which constitutes a fundamental description of the energy balance of the process of fire development, and hence also serves as a basis for Eq. (1.2), will be discussed at some length in Chapter 2.

What has been said up to this point relates only to the ignition and flame phases of the process of fire development. As regards the cooling phase, it is stipulated in a summary manner merely that the time graph of the temperature of the combustion gases shall be chosen so as to be linear, and that the rate of decrease in the gas temperature shall be taken to be  $10^{\circ}C \cdot \min^{-1}$ , unless other assumptions can be demonstrated to be more correct. If the design is carried out by means of the second or third alternative method, each of which is functionally realistic, then this implies that two phases of the same continuous process are represented in such a way that their descriptions are entirely different in the degree of accuracy as well as in the extent to which the actual conditions are taken into account. It is obvious that this gives rise to a considerable unbalance in the basis for design. To show how important it is that the cooling phase should also be described in a differentiated manner, it may be useful to give two examples of structural members which are characterised by low and high thermal inertia, respectively.

For an enclosed space, where the fire load is  $q = 12 \text{ Mcal} \cdot \text{m}^{-2}$  of bounding surface area, and the opening factor is  $A\sqrt{H}/A_t = 0.08 \text{ m}^{0.5}$ , Fig. 3 represents a calculated temperature-time curve for a steel column exposed to fire and characterized by the ratio  $F_s/V_s = 100 \text{ m}^{-1}$  and by  $\varepsilon_r = 0.5$ , where  $F_s$  is the total bounding surface area of the column, which is equal to its fire-exposed area, in m<sup>2</sup>, V<sub>s</sub> is the steel volume of the column, in m<sup>3</sup>, and  $\varepsilon_r$  is the resultant emissivity for heat transfer from the flames and the combustion gases to the



Fig. 3. Variation in the temperature,  $\vartheta_s$ , of a steel column with the time, t, calculated on the basis of temperature-time curves, which differ in shape during the cooling phase (dash-line curves).

 $F_s/V_s$  = Ratio, in m<sup>-1</sup>, of the fire-exposed surface area, i.e. the total bounding surface area, of the column to the steel volume of the column.

$$e_r = \text{Resultant emissivity}$$

 $-\vartheta_s$  $-\vartheta_g$  at  $A\sqrt{H}/A_t = 0.08 \text{ m}^{1/2}.$ 

= 12 Mcal  $\cdot$  m<sup>2</sup>.

- 1. Instantaneous cooling. Rate of decrease in temperature  $\infty^{\circ}C\cdot h^{-1}.$
- 2. Linear rate of decrease in temperature,  $600^{\circ}$ C · h<sup>-1</sup>.
- 3. Linear rate of decrease in temperature, 1200°C · h<sup>-1</sup>.
- 4. Linear rate of decrease in temperature, 6 min  $\leq t \leq 12$  min, 4600°C  $\cdot$  h<sup>-1</sup>;  $t \geq 12$  min, 600°C · h<sup>-1</sup>.

 $F_s/V_s = 100 \text{ m}^{-1} \varepsilon_r = 0.5.$ 

steel column. The characteristics of the flame phase for the temperature-time curve of the enclosed space have been chosen in conformity with the Swedish Building Regulations 1967. In Fig. 3, the full-line curves represent the temperature of the steel,  $\vartheta_s$ , and the dash-line curves show the temperature of the combustion gases,  $\vartheta_g$ , on the basis of the four alternative assumptions concerning the cooling phase of the process of fire development which are stated in what follows.

(1) After the duration of the fire  $T = qA_t/(25A\sqrt{H}) = 6$  min, the temperature of the combustion gases drops instantaneously to ordinary room temperature.

(2) The temperature of the combustion gases decreases in accordance with the Swedish Building Regulations 1967 at a linear rate of  $10^{\circ}$ C · min<sup>-1</sup>.

(3) The temperature of the combustion gases decreases at a linear rate of  $20^{\circ}$ C · min<sup>-1</sup>.

(4) The temperature of the combustion gases is assumed to vary in a more realistic manner, that is to say, it drops to half their maximum temperature during the first 6 min of the cooling phase, and then decreases at a linear rate of  $10^{\circ}$ C · min<sup>-1</sup>.

At the end of the flame phase, the temperature of the steel is  $303^{\circ}$ C. After that, the temperature of the steel continues to increase during the cooling phase of the process of fire development to its respective maximum values corresponding to the four alternative assumptions, viz., 303, 799, 719, and 518°C.

Fig. 4 shows the effects on the load-bearing capacity of a reinforced concrete slab which are produced by different slopes of the linear cooling phase of the time-temperature curves. Each *T*-value on the horisontal axis corresponds to a specific time-temperature curve. *T* is the duration of the flame phase, and depends on the fire load. The distance from the centre lines of the reinforcing bars to the fire-exposed surface of the slab, is assumed to be 2 cm. The emissivity of the flames is taken to be 0.7. The temperature in the enclosed space is supposed to vary with the time during the flame phase in accordance with Eq. (1.1). The temperature of the combustion gases is assumed to decrease at linear rates of 5, 10, and  $20^{\circ}$ C · min<sup>-1</sup>, or to drop instantaneously to ordinary room temperature ( $\infty^{\circ}$ C · h<sup>-1</sup>). For a static load which causes failure at a temperature of the reinforcing bars  $\vartheta_{scr} = 450^{\circ}$ C, we obtain a fire resistance period,  $t_{fr}$ , which varies from 0.52 to 0.82 h, and for  $\vartheta_{scr} = 500^{\circ}$ C, the corresponding variation in the fire resistance period ranges from 0.72 to 1.01 h.

These examples show that it is necessary to calculate the fire resistance period of a building component so as to take account of that reduction in its load-bearing capacity, or in its separating capacity, which occurs during the cooling phase. Moreover, they indicate that those temperature-time curves for enclosed spaces which are to serve as a basis for such calculations should also be differentiated and as realistic as possible in the cooling phase.





the temperature-time curve.

----- Cooling phase taken into account.

---- Cooling phase not taken into account.

Concrete cover 2 cm.

 $\varepsilon_{fl}$  = Emissivity of the flames=0.7.

 $t_{fr}$  = Fire resistance period.

 $\vartheta_{scr}$  = Critical temperature of the steel, i.e. the temperature at which the reinforcement fails.

 $t_{fr} = 0.52$  to 0.82 h at  $\vartheta_{scr} = 450^{\circ}$ C.

 $t_{fr} = 0.72$  to 1.01 h at  $\vartheta_{scr} = 500^{\circ}$ C.

The object of the present investigation is therefore to evolve a method which shall be applicable to different combinations of the values of the air flow factor (ventilation factor),  $A\sqrt{H}$ , and the fire load, as well as the type of material of the structures bounding the enclosed space under exposure to fire of the wood fuel type. This shall enable that the temperature-time curves for the enclosed space to be calculated by means of a theoretical procedure so as to cover the whole process of fire development, and shall thus make it possible to carry out judicious structural fire engineering design on the basis of the variation in the temperature with the time during all phases of the fire.

In connection with a general treatment of the equation of heat balance, Chapter 2 deals with the problems which are met with when a theoretical calculation of the temperature-time curve for the combustion gases is extended so as to comprise the cooling phase. Chapter 3 describes the methods which have been used to tackle these problems, and the modifications of the equation of heat balance which have been necessary for this purpose. In Chapter 4, an account is given of the computer programme which has been prepared for the calculations, and which is represented in the form of a flow chart. In Chapter 5, the full-scale tests which have served as a basis for the present investigation are subjected to comparative theoretical analysis. The time graphs of the rate of combustion which have been determined with the help of these theoretical analyses are presented in Chapter 6. Finally, in Chapter 7, these graphs are used as a basis for the calculation of complete temperaturetime curves for combustion gases in enclosed spaces which vary in the values of the opening factor and the fire load, as well as in type of material employed in the structures bounding the enclosed space.

#### 2. Equation of heat balance of process of fire development

The papers published by *Kawagoe* and *Sekine* [2], as well as by  $\ddot{O}deen$  [3], in the early 1960ies have made it possible to carry out theoretical calculations of temperature-time curves for combustion gases in the flame phase of the process of fire development to a degree of accuracy that is sufficient for practical purposes. These three authors have studied the energy balance during the process of fire development. The quantity of energy released per unit time, just as the volume of combustion gases evolved, during combustion were assumed to be known. With the help of the calculation of the quantity of energy which was lost per unit time by conduction and radiation from the enclosed space through its bounding structures, it was possible to deduce an equation of heat balance, and to solve it so as to obtain the temperature of the combustion gases. The treatment of this problem was based on the simplified assumptions which are reproduced in what follows.

(a) The temperature in the interior of the whole enclosed space is uniform at any given instant.

(b) The coefficient of heat transfer to the interior bounding surfaces of the enclosed space is uniform at every point.

(c) The heat flow through the bounding structures of the enclosed space is one-dimensional and, except for the window and door openings, if any, uniformly distributed.

Kawagoe and Sekine, as well as Ödeen, confined themselves throughout their papers to a study of the flame phase of the process of fire development, and the equation which they have deduced cannot be applied directly to the cooling phase. Primarily, the calculation of the temperature-time curve for the combustion gases during the cooling period requires an analytical investigation of two fundamental sub-problems, which have been but little studied up to the present time. In the first place, it is necessary to determine the quantity of energy liberated per unit time when this quantity, as is the case in the cooling phase, is no longer determined by the rate of air supply. In the second place, it is required to investigate the thermodynamic conditions which are encountered when the rate of combustion is no longer limited by the dimensions of the openings in the enclosed space. The equation of energy balance which has been deduced by Kawagoe and Sekine and by Ödeen, as

well as that extension of the theory which is required for the calculation of the temperature-time curve in the cooling phase of the process of fire development, will be dealt with in what follows.

The above-mentioned equation expresses for any given instant t, the balance between the respective quantities of heat energy generated and lost per unit time in the enclosed space under consideration. In its complete form, this equation is

$$I_C = I_L + I_W + I_R + I_B (2.1)$$

where

 $I_c$  = the heat energy released per unit time during combustion,

- $I_L$  = the heat energy withdrawn per unit time from the enclosed space owing to the replacement of hot gases by cold air,
- $I_W$  = the heat energy withdrawn per unit time from the enclosed space through wall, roof or ceiling, and floor structures,
- $I_R$  = the heat energy withdrawn per unit time from the enclosed space by radiation through the openings in the enclosed space,
- $I_B$  = the heat energy stored per unit time in the gas volume which is contained in the enclosed space.

The terms entering into the above equation are schematically illustrated in Fig. 5.



Fig. 5. Schematic illustration of the terms entering into the equation of heat balance.

#### Term $I_B$

In comparison with the quantities of energy which are involved in fires, the quantity of energy which can be stored in the gas volume contained in the enclosed space is of minor importance. Therefore, the term  $I_B$  can to a close approximation be put equal to zero.

#### Term $I_C$

In the case of fires of the wood fuel type, the term  $I_c$ , which expresses the quantity of heat released per unit time during combustion, is a factor which is difficult to determine in the equation of heat balance. In order to obtain this term, Kawagoe and Sekine, as well as Ödeen, chose as a point of departure the rate of combustion, expressed in terms of kilogrammes of wood per unit time, which they multiplied by a heat value.

This gives the equation

$$I_c = R \cdot W \quad \text{kcal} \cdot h^{-1} \tag{2.2}$$

where

R = the rate of combustion, in kg of wood per h, W = the heat value, in kcal per kg of wood.

With the exception of a small number of calculations which were based on a triangular variation in the rate of combustion, R, with the time, Ödeen assumed that R is constant, and is arbitrarily chosen, and that the quantity of heat, W, released during the flame phase is 4120 kcal per kg of wood. Thus, the heat energy liberated per unit time during combustion,  $I_c$ , is supposed to remain constant until the fuel has burnt up, and the cooling phase is characterized by the fact that no additional energy is supplied to the enclosed space. This description of the quantity of energy developed per unit time is applicable to burning of liquid fuels, but does not take account of the real combustion characteristics of wood fuels. As a rule, these fuels give off 30 to 50 per cent of the total quantity of energy after the end of the flame phase.

The investigation made by Kawagoe and Sekine is more differentiated in respect of the characterization of the process of fire development. As *Kawagoe* had shown in an earlier paper [5], during that period of the fire when the rate of combustion reaches a maximum, i.e. during the flame phase, the rate of air supply to the enclosed space, and hence also the maximum rate of combustion,  $R_{max}$ , are proportional to the air flow factor,  $A\sqrt{H}$ .

If the areas are expressed in  $m^2$ , and the maximum rate of combustion,  $R_{max}$ , is expressed in kg of wood per min, then we have the approximate relation

$$R_{\text{max}} = 5.5 \cdot A \sqrt{H}$$
 kg of wood per min (2.3)

Furthermore, in the papers published by Kawagoe and Sekine, the quantity of heat liberated during the flame phase, W, is stated to be 2575 kcal per kg of wood. This value was obtained by reducing the nominal heat value of wood so as to take account of the degree of incomplete combustion. The degree of incomplete combustion was estimated with the help of those analyses of the composition of the combustion gases which were carried out during fire tests [5]. As regards the cooling phase, Kawagoe and Sekine had made a few isolated comparative calculations based on a polygon-shaped time graph of the rate of combustion, and then found that the temperatures obtained when the cooling phase was characterized by a linear decrease in temperature at a rate of 7 or  $10^{\circ}$ C · min<sup>-1</sup> were much too high.

Accordingly, if the results of the investigations made by Ödeen, as well as by Kawagoe and Sekine, are to be applied to fires where the fuel is of the wood type, then the calculations have to be confined to the flame phase of the process of fire development.

When the treatment of this problem is extended so as to comprise the cooling phase of the process of fire development, Eq. (2.3) is not generally applicable. For the quantity of energy released per unit time during combustion, Eq. (2.3), in combination with Eq. (2.2), gives only the theoretical upper limit, which is determined by the available rate of air supply. During the cooling phase, the energy liberated per unit time will be governed by other factors. For this reason, and since no systematic investigation has so far been made in order to determine the relations between the three quantities which are of interest in this connection, viz., the reduction in the weight of fuel, the quantity of energy developed per unit weight of fuel, and the requisite rate of air supply, the quantity of energy released per unit time during the cooling phase of the process of fire development had to be determined by means of the method described in what follows. The procedure in calculation for the solution of the equation of heat balance was programmed for a CD 3600 computer. A study of the literature was carried out in order to examine the available publications on full-scale tests. A number of these tests were selected in the cases where the reported data were so complete as to enable numerical treatment. After that, the computer was used to calculate the temperature-time curve on the basis of an assumed form of the time graph of the rate of combustion for the complete process of fire development. The time graph of the rate of combustion was then varied until the agreement between the experimental and theoretical temperature-time curves was as close as possible. The only absolute requirement to be fulfilled in this connection was that the total quantity of energy liberated during the whole process of fire development should be equal to the total energy of combustion of the fuel. When an adequate range of variation in the opening factor and in the fire load was considered to have been covered, the time graphs of the rate of combustion obtained in this way were systematized. For a given fire load and a given opening factor, it was then possible to assume that the curve showing the variation in the rate of combustion with the time was known on the basis of this systematization. The investigation referred to in the above is described in Chapter 5, and its results, expressed in terms of time graphs of the rate of combustion in a dimensionless form, are presented in Chapter 6.

In connection with the treatment of the term  $I_c$ , it should be pointed out that the fire load must be described in combustion engineering terms in such a way that it may be associated with the equation of heat balance, Eq. (2.1).

In most countries, the fire load is expressed in terms of the quantity of wood that is equivalent to it in heat value per unit *floor area*. This characterization must be replaced by a parameter which has a physical significance when it is treated in calculations. This has been done in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures", which stipulate that the fire load shall be stated as that total quantity of heat, q, in Mcal  $\cdot$  m<sup>-2</sup> of total bounding surface area of the enclosed space exposed to fire, which is liberated on the assumption of complete combustion of all the combustible material contained in the enclosed space. A still more refined description of the fire load, which should take into account the variation with the time in the quantity of energy released by combustion, as well as the emissivities of the flames and the combustion gases, is an urgently recommended subject for research in this field.

#### Term $I_L$

The term  $I_L$  in the equation of heat balance expresses that quantity of heat which is withdrawn per unit time from the enclosed space owing to the replacement of hot gases by cold air through the openings in the enclosed space. For the determination of  $I_L$ , Ödeen, as well as Kawagoe and Sekine, used the equation

$$I_L = R \cdot G_0 \cdot (\vartheta_g - \vartheta_0) \cdot c_p \tag{2.4}$$

where

R = the rate of combustion, in kg of wood per h,

 $G_0$  = the volume of combustion gases produced by the fire, in Nm<sup>3</sup> · kg<sup>-1</sup> of fuel,<sup>1</sup>)

 $c_p$  = the specific heat of the combustion gases, in kcal  $\cdot$  Nm<sup>-3</sup>  $\cdot$  °C<sup>-1</sup>,

 $\vartheta_g$  = the temperature of the combustion gases, in °C,

 $\vartheta_0$  = the temperature of the air outside the enclosed space, in °C.

Eq. (2.4) states that the term  $I_L$  is put equal to the heat content of the combustion gases produced by the fire with reference to that of the outside air. On account of the difference in density between the cold outside air and the hot gases in the interior of the enclosed space, an exchange of heat by convection takes place in the openings of the enclosed space. The rate of this heat exchange determines the maximum value of the rate of combustion so far as the supply of oxygen is concerned. Therefore, Eq. (2.4) can be used as an expression for  $I_L$  when the rate of combustion is determined by the rate

<sup>&</sup>lt;sup>1</sup>)  $Nm^3$ =normal cubic metre=the quantity of a gas which occupies a volume of 1 cubic metre at 0°C and 760 mm barometric pressure.

of air supply, in spite of the fact that this equation in itself does not describe  $I_L$ , but expresses the heat content of the combustion gases produced by the fire. However, if the rate of combustion is limited by factors other than the rate of air supply, e.g. by the available quantity of fuel, or if the combustion is completed, then it is obvious that Eq. (2.4) does not hold good. Accordingly, a theoretical treatment of the cooling phase of the process of fire development requires an expression for  $I_L$  which is more generally applicable, and which is based on the rate of air exchange. This problem has been studied in a thesis for degree of Master of Engineering prepared by Ahlquist and Thelandersson [6], in the Division of Structural Mechanics and Concrete Construction, Lund Institute of Technology, Lund, Sweden. These authors based their study on the assumption that the static pressure distribution in the enclosed space varies linearly from the floor to the ceiling or roof, and that there exists a level (the neutral layer) at which the pressure in the enclosed space is equal to the pressure outside the enclosed space. This simplified model has been used by several authors, e.g. by Kawagoe [5] in the theoretical deduction of Eq. (2.3), and by Thomas [7] in studies which dealt with venting in the course of fires. In this connection, Kawagoe showed by means of a number of tests that the assumed pressure was actually applicable to a close approximation. On the assumption that the openings in the enclosed space are vertical only, Kawagoe used Bernoulli's equation to determine the quantities of outgoing combustion gases and incoming cold air as functions of the difference in temperature and the position of the neutral layer. On the basis of the condition that the difference between the quantities of gases flowing into and out of the enclosed space shall be equal to the difference between the quantities of gases produced and consumed by combustion, Ahlquist and Thelandersson calculated the position of the neutral layer as a function of the temperature and the rate of combustion. In this calculation, it was assumed that the rate of combustion may vary from zero to a maximum value, which is dependent on the dimensions of the openings. Furthermore, it was assumed that the liberation of a certain definite quantity of energy is associated with the consumption of the same quantity of air and the production of the same quantity of combustion gases, irrespective of the rate of combustion.

After the position of the neutral layer had been determined in this way, it was possible to obtain an expression for  $I_L$  at different values of the temperature and the rate of combustion. In the above-mentioned thesis [6], the treatment was also extended so as to comprise the modifications which are necessary when the enclosed space is provided with a vent. The deduction of the equations which are required for this purpose is reproduced in its main features in Chapter 3 of the present publication. Moreover, this chapter also contains a summary treatment of the case where the roof of the enclosed space is provided with horizontal openings.

#### Term $I_W$

The term  $I_W$  denotes the quantity of heat which is withdrawn per unit time from the enclosed space through the structures bounding this space. The term  $I_W$  is determined by solving the general equation of heat conduction in the one-dimensional case under non-steady flow conditions so as to take into account those thermal properties of the materials which are dependent on the temperature, the evaporation of occluded water, and the possible structural transformations in the materials entering into the bounding structures. This equation is

$$c \cdot \gamma \cdot \frac{\partial \vartheta}{\partial t} = \frac{\partial}{\partial x} \left( \lambda_x \cdot \frac{\partial \vartheta}{\partial x} \right)$$

where

c = the specific heat of the wall material,

 $\gamma$  = the weight per unit volume of the wall material,

 $\lambda_x$  = the thermal conductivity of the wall material,

 $\vartheta$  = the temperature in the interior of the wall material,

t =the time,

x = the position co-ordinate.

The above equation is solved by means of a numerical procedure which had been described by *Odemark* [8], among others, and which has subsequently been further developed by *Ödeen*, and others, [9]. The walls, ceiling, roof, and floor structures which bound the enclosed space are divided into *n* layers having the thickness  $\Delta x_k$  each, and the equation of heat balance is written for each one of these layers. If the temperature at the centre of the layer *k* at the time *t* is denoted by  $\vartheta_k$ , and that at the time  $t + \Delta t$  is designated by  $\vartheta_k + \Delta \vartheta_k$ , then the application of this procedure gives the relations

$$\varphi_{1} \cdot \frac{\Delta \vartheta_{1}}{\Delta t} = \psi_{1}(\vartheta_{g} - \vartheta_{1}) - \psi_{2}(\vartheta_{1} - \vartheta_{2})$$

$$\vdots$$

$$\varphi_{k} \cdot \frac{\Delta \vartheta_{k}}{\Delta t} = \psi_{k}(\vartheta_{k-1} - \vartheta_{k}) - \psi_{k+1}(\vartheta_{k} - \vartheta_{k+1})$$

$$\vdots$$

$$\varphi_{n} \cdot \frac{\Delta \vartheta_{n}}{\Delta t} = \psi_{n}(\vartheta_{n-1} - \vartheta_{n}) - \psi_{n+1}(\vartheta_{n} - \vartheta_{0})$$
(2.6)

(2.5)

where

$$\begin{split} \vartheta_{g} &= \text{the temperature of the gases in the enclosed space,} \\ \vartheta_{0} &= \text{the temperature of the outside air,} \\ \varphi_{k} &= \Delta x_{k} \cdot c(x, \vartheta) \cdot \gamma \\ \psi_{1} &= \frac{1}{\frac{1}{\frac{1}{\alpha_{i}(\vartheta)} + \frac{\Delta x_{1}}{2 \cdot \lambda(x, \vartheta)}}} \\ \vdots \\ \psi_{k} &= \frac{1}{\frac{\Delta x_{k-1}}{2 \cdot \lambda(x, \vartheta)} + \frac{\Delta x_{k}}{2 \cdot \lambda(x, \vartheta)}} \\ \vdots \\ \psi_{n+1} &= \frac{1}{\frac{\Delta x_{n}}{2 \cdot \lambda(x, \vartheta)} + \frac{1}{\alpha_{n}(\vartheta)}} \end{split}$$

where

 $\alpha_t(\vartheta) =$  the coefficient of heat transfer at the internal surface,  $\alpha_u(\vartheta) =$  the coefficient of heat transfer at the external surface,  $\lambda(x, \vartheta) =$  the thermal conductivity at the section x,  $c(x, \vartheta) =$  the specific heat at the section x,  $\gamma =$  the weight per unit volume at the section x.

The coefficient of heat transfer,  $\alpha_i$ , at the internal surface exposed to fire may be supposed to consist of two components, viz., first, a radiation component, which is markedly predominant at the high temperatures in question, and second, a convection component, which can be chosen with adequate accuracy so as to be constant, and to be equal to 20 kcal  $\cdot$  m<sup>-2</sup>  $\cdot$  h<sup>-1</sup>  $\cdot$  °C<sup>-1</sup> [9]. By applying the Stefan-Boltzmann law, this gives, for  $\alpha_i$ , the relation

$$\alpha_{i} = \frac{4.96 \cdot \varepsilon_{res}}{\vartheta_{g} - \vartheta_{i}} \left[ \left( \frac{\vartheta_{g} + 273}{100} \right)^{4} - \left( \frac{\vartheta_{i} + 273}{100} \right)^{4} \right] + 20 \text{ kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$$

$$(2.7a)$$

where

 $\vartheta_i$  = the temperature of the internal surface,

 $\varepsilon_{res}$  = the resultant emissivity for radiation between flames, combustion gases, and the internal surface.

The resultant emissivity,  $\varepsilon_{res}$ , is determined from the formula

$$\frac{1}{\varepsilon_{res}} = \frac{1}{\varepsilon_{fl}} + \frac{1}{\varepsilon_i} - 1 \tag{2.7b}$$

where

 $\varepsilon_{fl}$  = the emissivity of the flames,

 $\varepsilon_i$  = the emissivity of the surface exposed to fire.

Properly speaking, Eq. (2.7b) represents the emissivity for radiation between two parallel surfaces, but it was considered to be the best available approximation in the case of radiation between flames and a surface exposed to fire.

According to [9], the coefficient of heat transfer,  $\alpha_u$ , at the external surface, which is not exposed to fire, can be represented by the approximate expression

$$\alpha_u = 7.5 + 0.028 \cdot \vartheta_u \quad \text{kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$$
(2.8)

where

 $\vartheta_u$  = the temperature of the external surface.

The system of differential equations of the first order, Eq. (2.6), is solved numerically (see Chapter 4), and then the term  $I_W$  is given by the relation

$$I_W = A_t \cdot \psi_1 \cdot (\vartheta_g - \vartheta_1) \tag{2.9}$$

where

 $A_t$  = the total area of the surfaces bounding the enclosed space.

If the structures bounding the enclosed space consist of different materials, or if they differ in thickness, as is usually the case in practice, then the abovementioned operations are carried out separately for each type of structure, and after that the term  $I_W$  is obtained from the expression

$$I_{W} = \sum_{j} I_{W,j} = \sum_{j} A_{j} \cdot \psi_{1,j} \cdot (\vartheta_{g} - \vartheta_{1,j})$$

$$(2.10)$$

The present section, which deals with the term  $I_W$ , is based in its entirety on the publications of Kawagoe and Sekine as well as Ödeen referred to in the above.

#### Term $I_R$

Kawagoe and Sekine calculated the term  $I_R$  from the following formula, which is a generalization of the Stefan-Boltzmann law:

$$I_R = A \cdot (E_g - E_0) \tag{2.11}$$

where

A = the area of the opening,

$$E_g = 4.96 \cdot \left(\frac{\vartheta_g + 273}{100}\right)^4$$
$$E_0 = 4.96 \cdot \left(\frac{\vartheta_0 + 273}{100}\right)^4$$

This formula is applicable to the whole duration of the process of fire development, and is used in its unchanged form for the calculations in the present publication.

# 3. Study of combustion gas flow and heat flow through openings in enclosed spaces

As has been shown in the section dealing with the term  $I_L$  in Chapter 2, the expression given by Eq. (2.4) is applicable only during the flame phase of the process of fire development. In that section, the term  $I_L$  was determined on the basis of a maximum rate of combustion. If this expression is to be extended so as to be valid for the cooling phase of the fire also, then this requires that the rates of gas and air flow through the openings in the enclosed space shall be determined directly, and that these rates shall then be used as a point of departure for determining the quantity of heat which is withdrawn per unit time from the enclosed space. Similar problems, which were defined in thermodynamic terms, and which related to fires in enclosed spaces, have been treated by Kawagoe [5], among others. The primary prerequisites to such a treatment are the assumptions that the pressure distribution in a vertical direction is linear, and that there exists a neutral layer or zone, i.e. a level at which the static pressure in the interior of the enclosed space is equal to the atmospheric pressure outside the enclosed space. From these assumptions, Kawagoe deduced the expression for the maximum rate of combustion,  $R_{\text{max}}$ , which is given by Eq. (2.3). On the assumption that the position of the neutral zone is the unknown variable, he used the Bernoulli equation to calculate the respective quantities of gases and air which flow out of and into the enclosed space per unit time. After that, he determined the position of the neutral zone from the condition that the rate of flow of the incoming gases shall be equal to the rate of flow of the gases which are consumed by combustion, and that the rate of flow of the outgoing gases shall be equal to the rate of flow of the gases which are produced by combustion. Finally, on the assumption that the quantity of air consumed per unit weight of fuel is known, he calculated the maximum rate of combustion,  $R_{\text{max}}$ .

In this chapter, a similar analysis will be carried out in what follows. The purpose of this analysis is to determine the term  $I_L$  by an expression which is more general than that given by Eq. (2.4). A detailed deduction will be presented for the case where the enclosed space is provided with one or several vertical openings which are equal in height. After that, we shall deal with the modifications which are required in the applications which involve openings of other types.

#### Flow conditions in vertical rectangular opening in enclosed space

The interchange between the gaseous products of combustion and the combustion air takes place because the density of the hot gases is lower than that of the cold air outside the enclosed space. On the assumption that the temperature in the whole enclosed space is uniform, and that there exists a neutral zone, the velocities of gas and air flow can be determined theoretically. After that, if the dimensions of the opening are known, it is possible to calculate the respective rates of flow, i.e. the masses of the outgoing gases and the incoming air per unit time.

The velocity distribution in a vertical rectangular opening is schematically represented in Fig. 6.

The difference in static pressure between the outside and the inside is equal to zero at the level of the neutral zone. Accordingly, if use is made of the notations given in Fig. 6, the pressure difference,  $P_{y}$ , above the neutral zone i

$$P_y = (\rho_0 - \rho_g) \cdot y \tag{3.1a}$$

and the pressure difference,  $P_z$ , below the neutral zone is

$$P_z = -(\rho_0 - \rho_g) \cdot z \tag{3.1b}$$

where

 $\rho_0 =$  density of the outside air,

 $\rho_g$  = density of the combustion gases.

The density of the gaseous products of combustion is assumed to be equal to the density of the air at the same temperature [10].



Fig. 6. Gas flow in an enclosed space provided with a vertical opening.

A = Area of the vertical opening in the enclosed space.

H = Height of the vertical opening in the enclosed space.

H'' = Height of the vertical opening above the neutral zone level.

H' = Height of the vertical opening below the neutral zone level.

 $v_y$  = Velocity of gas flow above the neutral zone level.

 $v_z$  = Velocity of gas flow below the neutral zone level.

From Bernoulli's theorem we obtain the following expressions for the variation in the velocity of flow with the distance from the neutral zone

$$v_y = \sqrt{2g \cdot y \cdot \frac{\rho_0 - \rho_g}{\rho_g}}$$
(3.2a)

and

$$v_z = \sqrt{2g \cdot z \cdot \frac{\rho_0 - \rho_g}{\rho_0}}$$
(3.2b)

Then the rate of flow of the outgoing gases is

$$Q_{\text{out}} = \mu \cdot B \cdot \rho_g \cdot \int_0^{H''} v_y \cdot dy$$
(3.3a)

and the rate of flow of the incoming air is

$$Q_{\rm in} = \mu \cdot B \cdot \rho_0 \cdot \int_0^{H'} v_z \cdot dz \tag{3.3b}$$

where

 $\mu$  = the coefficient of contraction,

B = the width of the opening.

By substituting Eqs. (3.2a) and (3.2b) in Eqs. (3.3a) and (3.3b), respectively, we can directly calculate  $Q_{out}$  and  $Q_{in}$ . We find

$$Q_{\text{out}} = \frac{2}{3} \cdot \mu \cdot B \cdot (H'')^{3/2} \cdot \sqrt{2g \cdot \rho_g \cdot (\rho_0 - \rho_g)}$$
(3.4a)

and

$$Q_{\rm in} = \frac{2}{3} \cdot \mu \cdot \boldsymbol{B} \cdot (H')^{3/2} \cdot \sqrt{2g \cdot \rho_0 \cdot (\rho_0 - \rho_g)}$$
(3.4b)

The position of the neutral zone is determined by the equation of gas interchange in the enclosed space. This equation states that the difference between the rates of flow of the outgoing gases and the incoming air shall be equal to the difference between the rates of flow of the gases which are produced and consumed by combustion.

The mass of air contained in the enclosed space is assumed to be constant during the whole period of time under consideration. That total error, referred to the whole duration of the process of fire development, which is caused by this assumption in the calculation of heat flow is not greater than the heat content of the volume of air in the enclosed space. In comparison with the quantities of heat which are associated with fires, and in view of the other approximations which have been made in connection with the application of Eq. (2.1), this error may be regarded as negligible.

#### Deduction of maximum rate of combustion determined by rate of air supply in accordance with [5]

On the assumption that the rate of combustion, R, is determined by the rate of air supply  $(R = R_{max})$ , we have

$$Q_{\rm out} = R_{\rm max} \cdot G_0 \cdot \rho_0 \tag{3.5a}$$

and

$$Q_{\rm in} = R_{\rm max} \cdot L \cdot \rho_0 \tag{3.5b}$$

where

- $G_0$  = the volume of combustion gases, in Nm<sup>3</sup>, produced by the combustion of 1 kg of fuel,
- L = the volume of air, in Nm<sup>3</sup>, consumed by the combustion of 1 kg of fuel.

By substituting Eqs. (3.4a) and (3.4b) in Eqs. (3.5a) and (3.5b), respectively, we can calculate the position of the neutral zone, i.e. H' and H''. After that, by substituting H' in Eq. (3.5b), we obtain, for  $R_{max}$ , the expression

$$R_{\max} = \kappa(\Delta \vartheta) \cdot A \cdot \sqrt{H} \tag{3.6}$$

where

 $\kappa(\Delta \vartheta)$  is a coefficient, which depends on  $\Delta \vartheta$ 

and

 $\Delta \vartheta = \vartheta_q - \vartheta_0$ 

The values of  $\kappa(\Delta \theta)$  is calculated for two fuels. First, for fires of the wood fuel type, which are most characteristic of actual fires, because wood usually constitutes the predominant fire load. Second, for fires, where the fuel consists of alcohol. The numerical values used in these calculations are given in what follows [10].

$$\mu = 0.7$$

$$G_0 = 4.86 \text{ Nm}^3 \cdot \text{kg}^{-1} \text{ for wood}$$

$$6.22 \text{ Nm}^3 \cdot \text{kg}^{-1} \text{ for alcohol}$$

$$L = 3.98 \text{ Nm}^3 \cdot \text{kg}^{-1} \text{ for wood}$$

$$5.23 \text{ Nm}^3 \cdot \text{kg}^{-1} \text{ for alcohol}$$

The relation between  $\kappa$  and  $\Delta \vartheta$  is represented in Fig. 7. This graph shows that the variation in the value of  $\kappa$  with the temperature is very slight in the temperature range which is met with in fires. For practical applications, Kawagoe put the value of  $\kappa$  for wood fires equal to 330 kg  $\cdot$  h<sup>-1</sup>  $\cdot$  m<sup>-5/2</sup>, irrespective of the temperature. The value in question was used in deducing Eq. (2.3). This equation has been verified experimentally by several authors in model tests as well as in full-scale tests, see [5]. Eqs. (3.6) and (2.3) are applicable only to one or several openings of equal height, *H*. In these equations, *A* denotes the sum of the areas of the individual openings.



Fig. 7. Relation between the coefficient  $\varkappa$  in the equation  $R_{\max} = \varkappa A \sqrt{H}$  and the temperature difference,  $\Delta \vartheta$ , between the combustion gases and the outside air. Wood fires and alcohol fires.

# Determination of quantity of heat, $I_L$ , withdrawn per unit time through openings in enclosed space during the whole process of fire development

In order that an expression may be valid throughout the duration of the process of fire development, it is necessary to presuppose that the rate of combustion, R, may assume all values in the interval extending from zero to the maximum value which is given by Eq. (2.3). Therefore, R can be written

$$R = a \cdot 330 \cdot A \cdot \sqrt{H} \text{ kg} \cdot \text{h}^{-1}$$

where

 $0 \le a \le 1$ 

The balance between the respective rates of flow of the outgoing gases and the incoming air is given by the relation

$$Q_{\text{out}} - Q_{\text{in}} = G_0 \cdot R \cdot \rho_0 - L \cdot R \cdot \rho_0 \tag{3.8}$$

After substitution of  $Q_{out}$  and  $Q_{in}$  from Eqs. (3.4a) and (3.4b), respectively, and after simplification, we obtain

$$\left(\frac{H''}{H}\right)^{3/2} \cdot \sqrt{\rho_g(\rho_0 - \rho_g)} - \left(1 - \frac{H''}{H}\right)^{3/2} \cdot \sqrt{\rho_0(\rho_0 - \rho_g)} =$$
$$= \frac{(G_0 - L) \cdot \rho_0 \cdot 330a}{\frac{2}{3}\mu \cdot \sqrt{2g}}$$
(3.9)

With the help of Gay-Lussac's law of volumes for gases

$$\rho_0 = \rho_g \left( 1 + \frac{\Delta \vartheta}{273} \right) \tag{3.10}$$

the ratio H''/H can be determined for different values of a and  $\Delta \vartheta$  from Eq. (3.9).

The quantity of heat which is withdrawn per unit time from the enclosed space can be written

$$I_L = Q_{\text{out}} \cdot c_p \cdot \frac{\Delta \vartheta}{\rho_0} \tag{3.11a}$$

where

 $c_p$ =the specific heat of the outgoing gases, in kcal  $\cdot$  m<sup>-3</sup>  $\cdot$  °C<sup>-1</sup>. By substituting Eq. (3.4a) in Eq. (3.11a), we get

$$I_L = \varphi(\Delta \vartheta) \cdot c_p \cdot \Delta \vartheta \cdot A \cdot \sqrt{H}$$
(3.11b)

where

$$\varphi(\Delta\vartheta) = \frac{2}{3}\mu \sqrt[3]{2g} \frac{\sqrt{\frac{\Delta\vartheta}{273}}}{1 + \frac{\Delta\vartheta}{273}} \cdot \left(\frac{H''}{H}\right)^{3/2}$$
(3.11c)

If the ratio H''/H is determined from Eq. (3.9), then Eq. (3.11c) yields  $\varphi(\Delta \vartheta)$  for different values of a and  $\Delta \vartheta$ . For combustion of wood fuel, this relation is represented in Fig. 8, which is based on the values of  $\mu$ ,  $G_0$ , and L given on p. 30. It is seen from this graph that the family of curves in question



Fig. 8. Relation between the coefficient  $\varphi$  in the equation  $I_L = \varphi \cdot c_p \cdot \Delta \vartheta \cdot A \cdot \sqrt{H}$  and the ratio  $a = R/330 \cdot A \cdot \sqrt{H}$  for various values of the temperature difference  $\Delta \vartheta$ .

consists of approximately parallel straight lines. Therefore, we can write

 $\varphi = \varphi_0 + 120 \cdot a$ 

where

 $\varphi_0$  is the value of  $\varphi$  for a=0.

Table 1 gives  $\varphi_0$  for various values of  $\Delta \vartheta$ .

For  $\Delta \vartheta > 300$  °C, the value of  $\varphi_0$  is nearly independent of the temperature, and may be assumed to range from 1500 to 1600.

If the rate of combustion, and hence the factor a, are known, then  $I_L$  can be calculated for any instant, t, during the process of fire development. The assumptions which have been chosen for the deduction of the expression for the term  $I_L$  will be briefly discussed in what follows.

In spite of the fact that steady-state conditions have been assumed in the calculation of the flow through the opening, the results obtained from this calculation can also be applied under non-steady-state conditions, because a change in a position of equilibrium almost immediately gives rise to the establishment of a new position of equilibrium [7].

	$\varphi_0,$		$\varphi_0,$
Δϑ, °C	$m^{1/2} \cdot h^{-1}$	∆ϑ, °C	$m^{1/2} \cdot h^{-1}$
10	515	500	1597
50	991	600	1587
100	1254	700	1567
200	1476	800	1551
300	1567	900	1552
400	1595	1000	1510

The calculation of  $I_L$  requires that the specific heat of the outgoing gases shall be known. The air content of these gases is dependent on the rate of combustion. When a=0, the outgoing gases consist of air alone, and when a=1, they consist of gaseous products of combustion only. However, the difference in the specific heat between the air and the combustion gases is very slight, and the specific heat may therefore to a close approximation be regarded as independent of the rate of combustion.

Furthermore, in the calculation of  $\varphi$ , it is assumed that the values of  $G_0$  and L remain constant during the whole process of fire development, irrespective of the rate of combustion. This assumption has not been verified by any physical considerations, but if we examine the right-hand member of Eq. (3.9), then we find that the effect produced by an error in the difference between  $G_0$  and L on the value of  $\varphi$  is comparable to that of an error in a, that is to say, this effect is very slight.

33

(3.12)

In his treatment which relates to the flame phase only, Kawagoe has assumed that the temperature is uniform in the whole enclosed space. In the present publication, this assumption has been extended so as to be applicable during the whole process of fire development. The assumption that the variation in temperature in a vertical direction is relatively slight when the intensity of the fire decreases has been confirmed by the full-scale tests which are described in Chapter 5 of this publication. In most of these tests, the dispersion in the temperature measured at different points in the enclosed space during the cooling phase was found to be smaller than during the flame phase.

#### Modification of treatment in cases where enclosed spaces are provided with several openings which differ in height

The deduction carried out in the preceding section is applicable only in the cases where the air is supplied to the enclosed space through one or several openings which are equal in height, and which have a common neutral zone. If the enclosed space is provided with several openings which differ in height, then a corresponding deduction can be made in each individual case. For the determination of the maximum rate of combustion,  $R_{max}$ , Yokoi [13] has described an approximate method. It consists in the determination of a fictitious air flow factor, which is used in the original formula, Eq. (2.3). The fictitious air flow factor is determined from the expression

$$(A \cdot \sqrt{H})_{\text{fict}} = \sum_{i} A_{i} \cdot \sqrt{H}_{i}$$
(3.13)

For some cases, Kawagoe [14] has compared the values of  $R_{max}$  which were obtained from Eq. (3.13) with those which were determined by means of accurate calculations. He found that Eq. (3.13) gives values which are sufficiently accurate for practical uses when the differences in the height and in the vertical position of the openings are not too great. The Swedish Building Regulations 1967 recommended another acceptable approximation, namely, that a weighted average of the heights of the individual openings should be used as a value of H in the calculation of the air flow factor.

After a fictitious value of the air flow factor has been determined from Eq. (3.13), this value can be used instead of  $A \cdot \sqrt{H}$  in Eq. (3.11b) where the factor  $\varphi$  is calculated from Eq. (3.12) as before.

#### Modification of treatment in cases where enclosed spaces are provided with horizontal openings in roofs

In the preceding two sections, it was assumed that all openings in the enclosed space are vertical. In the present section, we shall expound a more general theory which makes it possible to take account of the presence of
horizontal openings in the roof of the enclosed space. We suppose the enclosed space to be in conformity with Fig. 9.

We assume that there exists a neutral zone in the enclosed space, and that the level of this zone is not higher than the upper edge of the vertical opening, and not lower than its lower edge. A condition which is prerequisite to this assumption will be stated further on in the present section. For the vertical opening, the rate of flow of the outgoing gases,  $Q_{out}$ , is obtained from Eq. (3.4a), and the rate of flow of the incoming air,  $Q_{in}$ , is computed from Eq. (3.4b). The velocity of the gases which flow out through the horizontal opening is (Bernoulli's theorem)

$$v_h = \sqrt{2gh' \cdot \frac{\rho_0 - \rho_g}{\rho_a}} \tag{3.14}$$

This formula has been verified experimentally in connection with studies of venting fires [7]. The rate flow of the outgoing gases through the horizontal opening is

$$Q_h = \mu \cdot A_h \cdot \nu_h \cdot \rho_g \tag{3.15}$$



Fig. 9. Gas flow in an enclosed space provided with a vertical opening and a horizontal opening.

- A = Area of the vertical opening.
- $A_h$  = Area of the horizontal opening.
- H = Height of the vertical opening.
- H'' = Height of the vertical opening above the neutral zone level.
- H' = Height of the vertical opening below the neutral zone level.
- h = Vertical distance from the centre of the vertical opening to the level of the horizontal opening.
- h' = Vertical distance from the neutral zone to the level of the horizontal opening.

At a maximum rate of combustion,  $R = R_{max}$ , the equation of mass balance of gases requires that

$$\begin{array}{c}
Q_{\text{out}} + Q_{k} = G_{0} \cdot R_{\text{max}} \cdot \rho_{0} \\
Q_{\text{in}} = L \cdot R_{\text{max}} \cdot \rho_{0}
\end{array}$$
(3.16)

These two relations can be used to determine the position of the neutral zone which is modified in view of the presence of the horizontal opening, and then the maximum rate of combustion,  $R_{\text{max}}$ , can be calculated.  $R_{\text{max}}$  is a function of the term  $\frac{A_h \cdot \sqrt{h^2}}{A \cdot \sqrt{H}}$  at a given temperature. The value of  $R_{\text{max}}$  varies slightly with the temperature. If we write

$$R_{\rm max} = 330 \cdot (A \cdot \sqrt{H})_{\rm fict} \tag{3.17}$$

then  $(A \cdot \sqrt{H})_{\text{fict}}$  can be determined from the alignment chart in Fig. 10, which is entered at the value of  $A_h \cdot \sqrt{h}/A \cdot \sqrt{H}$ .

It is to be expected that the value of  $(A \cdot \sqrt{H})_{\text{fict}}$  determined in this manner may be used to an adequate degree of accuracy in the same way as the air flow factor,  $A \cdot \sqrt{H}$ , to characterize a fire. Accordingly, for an enclosed space with horizontal openings in the roof, the opening factor is given by the expression  $(A \cdot \sqrt{H})_{\text{fict}}/A_t$ , where  $(A \cdot \sqrt{H})_{\text{fict}}$  is determined from the alignment chart in Fig. 10. The term  $I_L$  in the equation of heat balance, Eq. (2.1), of the process of fire development is obtained by analogy with Eq. (3.11b) from

$$I_L = \varphi(\Delta \vartheta) \cdot c_p \cdot \Delta \vartheta \cdot (A \cdot \sqrt{H})_{\text{fict}}$$
(3.18)

where  $\varphi(\Delta \vartheta)$  is given by Eq. (3.12).



Fig. 10. Alignment chart for the calculation of the value of the modified air flow factor,  $(A/\overline{H})_{\text{fict}}$ , on the basis of known geometrical data. For notations, see Fig. 9.

This procedure presupposes that the flow through the horizontal opening in the roof is not predominant. Consequently, the factor  $A_h \cdot \sqrt{h'}/A \cdot \sqrt{H}$  has an upper limit at which the above model of the flow conditions ceases to be relevant. This upper limit is

$$\frac{A_h \cdot \sqrt{h'}}{A \cdot \sqrt{H}} = \begin{cases} 1.76 \text{ at } 1000^{\circ}\text{C} \\ 1.37 \text{ at } 500^{\circ}\text{C} \end{cases}$$

At this limit, the neutral zone is on a level with the upper edge of the vertical opening, and then h' is identical with the vertical distance from the level of the horizontal opening in the roof to the upper edge of the vertical opening. Tests [7] have indicated that the model used is relevant up to this upper limit. For values of  $A_h \cdot \sqrt{h'}/A \cdot \sqrt{H}$  which are higher than this limit, all gaseous products of combustion will be vented through the horizontal opening in the roof. This is sometimes intentional if it is desired that the spread of fire to adjoining rooms should be prevented by venting the fire [15]. If the air and the combustion gases flow in the main through horizontal openings, then the flow becomes unstable, and it is difficult to represent it by a simple theoretical model [16].

## Example showing how to use alignment chart in Fig. 10

Calculate the maximum rate of combustion,  $R_{\text{max}}$ , during the flame phase at 1000°C in the enclosed space characterized by the following data:

$$A = 2 \text{ m}^2$$
,  $H = 1 \text{ m}$ ,  $A_h = 1 \text{ m}^2$ ,  $h = 1.5 \text{ m}$ .  
 $\frac{A_h \cdot \sqrt{h}}{A \cdot \sqrt{H}} = \frac{1 \cdot \sqrt{1.5}}{2 \cdot \sqrt{1}} = 0.61; \quad \frac{A_h}{A} = 0.5$ 

The dash line in the alignment chart gives

$$(A \cdot \sqrt[]{H})_{\text{fict}} = 2.18 \cdot A \cdot \sqrt[]{H} = 4.36 \text{ m}^{5/2}$$
  
 $R_{\text{max}} = 330 \cdot 4.36 \text{ kg} \cdot \text{h}^{-1} = 1440 \text{ kg} \cdot \text{h}^{-1}$ 

This alignment chart can also be used for enclosed spaces where the horizontal opening in the roof is replaced by a ventilation duct. In such cases, the height h is replaced by the height of the gas column (the static head), with the reduction of the losses due to friction, which can be expressed in terms of the equivalent loss in static head. In an ordinary flat or office equipped with common ventilators made of non-combustible materials, the effect of the ventilation ducts is usually negligible.

If the enclosed space is ventilated through air inlets and outlets by means of a fan installation, then the corresponding fictitious air flow factor can be calculated in an analogous manner. If the quantity of gases exhausted per unit time is  $Q_{out}^{\nu}$  in kg  $\cdot$  h<sup>-1</sup>, and the quantity of air supplied per unit time is  $Q_{in}^{\nu}$  in kg  $\cdot$  h<sup>-1</sup>, then, for  $R = R_{max}$ , we have the conditions

$$\begin{array}{c}
Q_{\text{out}} + Q_{\text{out}}^{\nu} = G_0 \cdot R_{\text{max}} \cdot \rho_0 \\
Q_{\text{in}} + Q_{\text{in}}^{\nu} = L \cdot R_{\text{max}} \cdot \rho_0
\end{array}$$
(3.19)

from which  $R_{\max}$  and  $(A \cdot \sqrt{H})_{\text{fict}}$  can be determined by analogy with the above.

## 4. Description of programme for digital computer

The integration of the system of equations given by Eq. (2.6) was carried out by using the Runge-Kutta method in a modified form which has been suggested by Merson [22]. This modified method enables the computer to choose that interval of integration,  $\Delta t$ , which is required in order to ensure a certain definite degree of accuracy. To integrate the above-mentioned system of equations in the time interval from t to  $t + \Delta t$ , this system was evaluated in five individual operations. A determination of the temperature of the combustion gases is required for each one of these operations. According to Eqs. (2.1), (2.10), (2.11), and (3.11), the equation of heat balance of the process of fire development can be written

$$I_C = I_L + I_W + I_R \tag{4.1}$$

where

$$I_{L} = \varphi(\Delta \vartheta) \cdot c_{p} \cdot A \sqrt{H}(\vartheta_{g} - \vartheta_{0})$$
$$I_{W} = \sum_{j} I_{W, j} = \sum A_{j} \psi_{1, j}(\vartheta_{g} - \vartheta_{1, j})$$
$$I_{R} = A \cdot (E_{g} - E_{0})$$

If  $I_R$ ,  $\varphi$  and  $c_p$  are calculated on the basis of that value of the temperature of the combustion gases which has been obtained from the next preceding determination, then Eq. (4.1) can be solved for the temperature of the combustion gases,  $\vartheta_a$ , in an explicit form

$$\vartheta_{g} = \frac{I_{C} + \varphi(\Delta \vartheta) \cdot A \sqrt{H} \cdot \vartheta_{0} + \sum_{j} A_{j} \cdot \psi_{1, j} \cdot \vartheta_{1, j} - I_{R}}{\varphi(\Delta \vartheta) \cdot c_{p} \cdot A \sqrt{H} + \sum_{j} A_{j} \psi_{1, j}}$$
(4.2)

which can be substituted in the system of equations represented by Eq. (2.6).

The programme for the computer has been prepared in such a way as to be applicable to enclosed spaces bounded by structures which were assumed to be of up to three different types. Two of these structures were supposed to be homogeneous, whereas the third might be divided into two or three layers consisting of different materials, e.g. plasterboard panels, mineral wool, and brick.



enclosed space are supposed to comprise, first, a roof or/and ceiling and a floor made of homogeneous material, and second, partitions which consist of a Fig. 11. Flow chart for the calculation of the temperature-time curve of the process of fire development in an enclosed space. The structures which bound the load-bearing frame made of steel studs and covered on the inside as well as on the outside with two plasterboard panels, 13 mm in thickness each. The specific heat of most structural materials, e.g. concrete having a nonnegligible moisture content or plasterboard panels, varies discontinuously with the temperature when these materials are subjected to physical or chemical transformations. Therefore the programme used the enthalpy I, in kcal  $\cdot$  m<sup>-3</sup>, of the different materials as a dependent variable in the calculation of the temperature fields.

Fig. 11 represents a schematic flow chart which shows the programming procedure in the calculation of the combustion gases in an enclosed space where the floor and the roof or/and ceiling are made of concrete, while all the walls consist of a lightweight frame built of steel which are covered on the inside as well as on the outside with two plasterboard panels, 13 mm in thickness each. This type of wall exhibits two characteristic features, viz., first, experiments have shown that a plasterboard panel exposed to fire disintegrates when the temperature of the panel has reached a certain definite value, and second, the relation between the temperature and the enthalpy is dependent on the rate of temperature rise. Furthermore, it was necessary to choose different enthalpy-temperature curves depending on whether the temperature of the plasterboard panel in question was assumed to be increasing or decreasing. For further particulars, reference is made to the description of the calculations for the Type G enclosed space in Chapter 7.

# 5. Calculation of time graphs of rate of combustion for some full-scale tests described in literature

As has been mentioned in Chapter 2, the present chapter will deal with the comparative calculations which have led to a determination of the variation in the rate of combustion, expressed in kcal  $\cdot$  h<sup>-1</sup>, with the time. The method of successive approximations employed for this purpose consisted in making calculations which were based on different forms of the time graph of the rate of combustion. These calculations were repeated until they resulted in that curve which corresponded to the closest agreement between experimental and calculated curves representing the variation in the temperature of the combustion gases with the time. This method required a certain systematization of the description which represents the variation in the quantity of energy released by combustion with the time. What can be assumed to be known to a sufficient degree of accuracy in this connection is solely the total quantity of energy that can be liberated during the whole process of fire development, i.e. the fire load. Furthermore, it can be assumed that Eq. (2.3)

 $R_{\rm max} = 330 A \sqrt{H}$  kg of wood per h

expresses the maximum rate of combustion, in kg of wood per h. In order that a theoretical determination of the temperature of the combustion gases may be possible, it is moreover necessary to determine the relation between released energy and weight loss of fuel, W, in

 $I_{\rm C} = R \cdot W \quad \text{kcal} \cdot h^{-1}$ 

This determination is rendered difficult by the fact that the combustion of the gases formed by pyrolysis and that of solid wood fuel constitute a complicated process which involves a series of chemical reactions in different phases, see [4]. In respect of energy conditions, some of these reactions are endothermic, others exothermic. So far as the Authors know, no systematic investigation has been made up to now in order to carry out a quantitative analysis of the liberation of energy during the individual phases of the process of fire development.

These considerations have necessitated certain assumptions which concern the form of the curve showing the variation in the rate of combustion with the time. These assumptions are stated in what follows. The quantity of energy

liberated per unit time during the ignition phase was supposed to increase according to a polygonal function of the time to a level which corresponds to the rate of combustion during the flame phase. The determination of this level was based on Eq. (2.3),  $R_{\text{max}} = 330 A \sqrt{H}$  kg of wood per h. When the quantity of energy released per kg of wood fuel during the flame phase was assumed to range from 2500 to 2800 kcal, it was found that the calculated temperature-time curves were closely in agreement with the results of the full-scale tests in respect of the maximum temperature and the duration of the flame phase. In order to adapt these assumptions to the temperature-time curves which have been published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures", and which are used in Sweden as a basis for calculations of the fire endurance of structural components in conformity with standard specifications, the quantity of heat evolved during the flame phase was supposed to be 2575 kcal · kg<sup>-1</sup> of wood, since this value had been employed for calculating the above-mentioned temperature-time curves. Cf. Chapter 2. The value in question has originally been stated by Kawagoe, who used it as a measure of the quantity of energy that is liberated by incomplete combustion of 1 kg of wood. By analysing the composition of the combustion gases in fire tests, see e.g. [5], it has been found that they contain considerable quantities of carbon monoxide, and this indicates that the combustion is not complete. It seems that the analysis of the combustion gases was performed during the flame phase, and the value 2575 therefore applies to this phase only. During the cooling phase, the weight loss of fuel per liberated energy unit is considerably less than during the flame phase. This means, that, when the whole process of fire is considered, the energy released by the combustion of 1 kg of wood must be higher than 2575 kcal even if the combustion is incomplete. As a rule, there was scant basis for an accurate prediction of how the rate of completeness of combustion varied during the different phases. Consequently, in our comparative analysis the nominal heat value of wood, ranging from 3500 to 4500 kcal  $\cdot$  kg<sup>-1</sup>, was used as the energy liberated during the whole course of fire. This means that combustion is assumed to be complete throughout the present publication. The resulting time-temperature curves for the combustion gases therefore as a rule are found to be more in agreement with the maximum temperature curves than with the mean temperature curves obtained in the full scale tests.

Consequently, during the flame phase, the value of the rate of combustion will be constant,  $330 \cdot A\sqrt{H} \cdot 2575 \text{ kcal} \cdot \text{h}^{-1}$ . The cooling phase is characterized by a rate of combustion which decreases in conformity with a polygonal time graph in such a way that the slope of each individual side of the polygon is dependent on the duration of the flame phase. In every case, the area between the rate of combustion curve and the time axis must equal the fire load, expressed in energy units. The use of a polygonal time graph of the rate of

combustion makes it easier to check the agreement between the total quantity of energy liberated during the process of fire development and the fire load which is given from the outset.

The factors which must be taken into account in a comparative theoretical calculation of the temperature of the combustion gases in fire tests are enumerated in what follows.

(1) Characteristics of fuel: Quantity, moisture content, porosity factor (hydraulic radius), distribution in the enclosed space.

(2) Geometric characteristics of the enclosed space: Opening factor and its variation with the time (e.g. in the case where the fire burns through a door), shape of openings, cross-sectional area of ventilation ducts (if any).

(3) Characteristics of the structures bounding the enclosed space: Structural design, thermal properties and temperatures of disintegration (if any) of the materials entering into the structures, emissivity characteristics of the surfaces.

By studying the available literature on full-scale wood fire tests, it was found, first, that the number of published tests is relatively small, and second, that those data which are so detailed as to render possible a comparative theoretical calculation have been stated only for a few tests out of this number. For this reason, it was necessary to confine the comparative calculations to only four test series reported in the literature or in other sources.<sup>1</sup> These test series are enumerated below.

(1) Test series A. Tests carried out by *Sjölin*, and dealt with in a thesis for a L. Techn. degree, [17].

(2) Test series B. Tests made by Kawagoe, and published in [5].

(3) Test series C. Tests performed by *Ödeen* at the Royal Institute of Technology, Stockholm, and described in his doctoral thesis [18].

(4) Test series D. Tests directed by *Pettersson* and *Ödeen*, which were carried out in a test house of the Atomic Energy Co., Ltd. (AB Atomenergi) at Studsvik, Sweden.

## Test series A. Tests carried out by Sjölin. Calculation of time graphs of rate of combustion

The tests made by Sjölin were undertaken in order to study the spread of fire, and the process of fire development, in rooms and combinations of rooms exposed to ignition at a single point by heat radiation emitted by an explosion of nuclear weapons. The fire tests were carried out in a test house which was provided with concrete floor structures and concrete or lightweight concrete

<sup>&</sup>lt;sup>1</sup> Not until too late in the publication of this paper did the authors learn about the full scale fire tests carried out at Fire Research Station, Boreham Hood, London [11, 12].

This is the more regrettable as these experiments make an excellent basis for comparative theoretical calculations.

walls. The test house was designed in such a way that various enclosed spaces might be formed by individual rooms, or by several rooms connected together. The available model scales were 1 to 1, 1 to 2, and 1 to 4. The variables recorded in these tests were expressed as functions of the time, and comprised the temperature and the velocity of gas flow at characteristic points, the intensity of radiation, the composition of the gases, and the rate of combustion, which was determined by continuous weighing of the quantity of fuel in the enclosed space.

The fire load in all these tests consisted of authentic furniture. This was an extraordinarily valuable feature of the tests, seeing that all the other full-scale tests which are dealt with in the present publication were made by using fire loads of the wood crib type.

Seven of the fire tests included in this test series were found to be suited for the present theoretical study. In the other tests, the ignition did not cause the fuel to take fire. Table 2 shows the scope of the seven tests under consideration.

			Table 2	. Test series A	· · · · · · · · · · · · · · · · · · ·
Test No.	Type of room	Window area, m <sup>2</sup>	Opening factor, $\frac{A \cdot \sqrt[4]{H}}{A_t}$ $m^{1/2}$	Fire load, kg · m <sup>-2</sup> of bounding surface area	Remarks
1	В	1.16	0.0237-0.06	3.5	(The fire burnt through a door, 1.6
2	В	1.16	0.0237-0.06	4.4	$m^2$ in area, during the time interval from $t=0$ min. to $t=6$ min.
3	L	1,16	0.0160-0.0356	4.9	(The fire burnt through a door, 1.6
4	L	1.88	0.0278-0.0486	5.6	$\begin{cases} m^2 \text{ in area, during the time interval} \\ \text{from } t=8 \text{ min. to } t=12 \text{ min.} \end{cases}$
5	L	1.88	0.0548	5.0	
6	L	2.95	0.068	5.7	
7	B+L	3.20	0.040	8.1	

B = Furnished two-person bedroom, 10.4 m<sup>2</sup> in floor area.

L = Furnished living room, 18.8 m<sup>2</sup> in floor area.

B+L =Combination of B and L.

The curves representing the variation in the temperature of the combustion gases with the time in the above-mentioned seven tests, as well as the corresponding curves obtained by calculations with the help of automatic computer, are reproduced in Appendix 1. In the test No. 7, the enclosed space consisted of two contiguous rooms, which communicated through an open door. The partition between these two rooms was considered in the theoretical calculations to be an enclosed structure which possessed a heat-absorbing capacity. For this test, the curves representing the temperature of the combustion gases as a function of the time are shown separately for the bedroom and the livingroom.

All the fires in these tests were characterized by a protracted process of ignition, which was followed by a rapid transition to the flame phase. For the calculated curves, the time was put equal to zero at the instant when the flame phase began, i.e. when the fuel took fire, in the actual fire tests. The time graphs of the rate of combustion which were finally obtained from the calculations are also reproduced in the respective diagrams. In the tests Nos. 1 to 4, the opening factor was changed during the process of fire development because the fire had burnt through a door. This change was taken into account in the calculations, and constituted the cause of the somewhat unusual shape of the time graphs of the rate of combustion which refer to these four tests. In these cases, the curves were plotted in such a way as to relate the rate of combustion to that value of the opening factor which was obtained after the fire had burnt through the door. In the test No. 3, the fact that the fire has burnt through a door is reflected very clearly in the curve, which shows that the temperature of the combustion gases remained constant at 500 to 600°C, and then rapidly rose to about 800°C when the fire burnt through the door.

Since the calculations were based on the opening factor and on the total energy content of the fuel, it was possible to choose the time graphs of the rate of combustion in such a manner that the agreement between the observed and calculated time graphs of the temperature of the combustion gases was very close in all the tests except the test No. 7. In the test No. 7, the calculated time graph of the combustion gas temperature was compared with the corresponding observed curves for the living-room as well as for the bedroom. The agreement between these curves was relatively close in the second case, but not in the first, where the curve is slightly displaced in time with reference to the curve for the bedroom. Moreover, when use was made of the opening factor which was determined geometrically, the temperatures obtained for the flame phase were found to be somewhat too low. However, it is not correct to regard the above-mentioned two rooms as a single enclosed space, since there existed quite a considerable difference in temperature between these rooms. Nor are the two rooms in question to be regarded as two separate enclosed spaces, since a certain heat exchange took place between these rooms.

## Test series B. Tests made by Kawagoe [5]. Calculation of time graphs of rate of combustion

In [5], *Kawagoe* has described a large number of fire tests which had been carried out in Japan. In this investigation, he primarily studied the relation between the reduction in the weight of fuel per unit time and the dimensions

of the openings in the enclosed space. Among other things, he also deduced the equation which is reproduced in Eq. (2.3) in the present publication. The variables measured in these tests were the reduction in the weight of fuel, the temperature of the combustion gases, the composition of the combustion gases, the gas velocities, the intensity of radiation, and the pressure distribution in the window openings. On account of the above-mentioned main purpose of Kawagoe's investigation, it is only in three of these tests that the results of measurements, the geometric data, and the data on the materials entering into the structures which bounded the enclosed space are presented in such a way as to make it possible to carry out theoretical calculations of the type under consideration. These three tests were performed in a test house which was provided with walls made of hollow concrete blocks and with concrete floor and roof structures. The other test data are given in Table 3.

Table 3. Test series B.						
	Fire load, kg of	Bounding surface	Window dimensions,	Opening fac <sub>t</sub> or, $A \cdot \sqrt{H}$		
Test		area,	width $ imes$ height,	$\overline{A_t}$		
No.	wood	m² .	m×m	m <sup>1</sup> /2		
1	400	48	0.93×1.8	0.0467		
2	900	48	0.93  imes 1.8	0.0467		
3	1000	48	$0.93 \times 1.8$	0.0467		

Each hollow concrete block used for the walls comprised a single large cavity, without any subdivisions. The volume of the cavity was estimated at 30 to 40 per cent of the total volume of the block. In the calculations, the walls were considered to be composed of two different structures. One of them consisted of concrete alone, while the other comprised three layers, viz., concrete, air-filled cavity, and concrete, respectively. The second structure represented that part of the wall surface which corresponded to the cavities of the hollow concrete blocks, while the first structure was equivalent to the remaining part of the wall surface.

The observed and calculated temperature-time curves, as well as the time graphs of the rate of combustion used in the calculations, are reproduced in Appendix 2. The fire loads in the tests Nos. 2 and 3 were relatively high and the duration of the fires in these tests was therefore long. In order that the calculated values should agree with the values observed in these tests, it was necessary to choose a comparatively flat slope for the ascending branch of the time graph of the rate of combustion.

The data on the cooling phases in the tests Nos. 2 and 3 reported in [5] are not complete, and this is the reason why the curves relating to these tests break off at such an early stage.

## Test series C. Tests published by Ödeen [18]. Calculation of time graphs of rate of combustion

The tests described by Ödeen in [18] were carried out in a tunnel building of an approximately semi-circular shape, which had been specially constructed for this purpose. It was provided with a concrete wall, 20 cm in thickness, its total bounding surface area was 75 m<sup>2</sup>, and its total enclosed volume was 46 m<sup>3</sup>. A fan system made it possible to regulate and to measure the quantity of air which was supplied per unit time of the fire. In addition, a vent, 0.5 m<sup>2</sup> in cross-sectional area, for conveying the combustion gases to the outside air was provided in the upper part of each end wall of the tunnel.

A series of fire tests using fir wood as fuel has been carried out in this test building. A study was made of the effects produced on the process of fire development by the factors which are enumerated in what follows.

(1) The volume of air supplied per unit time to the tunnel, Q, in m<sup>3</sup> · s<sup>-1</sup>.

(2) The quantity of combustible material (fire load), M, in kg of wood.

(3) The hydraulic radius of the fuel, r, in cm. This factor expresses the ratio of the total volume of the fuel to its total bounding surface area.

The scope of the test series using wood fuel is shown in Table 4, which was extracted from [18].

In order that the results of these tests may be compared by means of calculations with those of fire tests in ordinary enclosed spaces, where the rate

			le 4. Test series C.		<b>.</b>
Test No.	Fire load, <i>M</i> , kg of wood	Rate of air supply, Q, m <sup>3</sup> · s <sup>-1</sup>	Hydraulic radius, <i>r</i> , cm	Moisture content of fuel, per cent	Energy content of fuel, Mcal
1	270	1.0	_	9	1129
2	675	2.0	1.0	17	2565
3	675	1.0	1.0	17	2565
4	675	0.7	1.0	22	2468
5	675	1,5	1.0	22	2468
6	675	MIN	1.0	21	2501
7	675	1.0	1.7	21	2501
8	675	0.7	1.7	21	2501
9	675	1.0	0.6	21	2501
10	270	0.7	1.0	21	1000
11	405	0.7	1.0	22	1481
12	405	0.7	2.4	28	1440
13	135	0.7	1.0	28	481
14	945	1.0	1.6	16	3659
15	1350	2.0	1,4	17	5130
16	405	0.7	0.4	17	1539

of air supply is determined by the openings in the enclosed space, the quantity of air supplied per unit time, Q, must be converted into an air flow factor,  $A\sqrt{H}$ , or into an opening factor,  $A\sqrt{H}/A_t$ . In this connection, it is necessary to take account of the air flow which may possibly enter into the tunnel through the outlets for combustion gases at low values of Q, and may therefore increase the value of the air flow factor.

In an ordinary enclosed space provided with a vertical opening, the rate of flow of the incoming air,  $Q_{in}$ , is given according to Eq. (3.4b), by the relation

$$Q_{\rm in} = 2/3 \ \mu B(H')^{3/2} \left[ \sqrt{2g \cdot \frac{(\rho_0 - \rho_\theta)}{\rho_0}} \ \mathrm{m}^3 \cdot \mathrm{s}^{-1} = \right]$$
$$= 2/3 \cdot 0.7 \cdot A \sqrt[3]{H} \sqrt{2g} (H'/H)^{3/2} \left( \frac{\Delta \theta}{273}}{1 + \frac{\Delta \theta}{273}} \right)^{1/2} \ \mathrm{m}^3 \cdot \mathrm{s}^{-1} \quad (5.1)$$

The notations used in Eq. (5.1) have been explained in Chapter 3.

With the help of Eqs. (3.9) and (3.10), we obtain, for a=1, the relation between H''/H and  $\Delta \vartheta$  shown in Table 5.

	Table 5. Relation b	etween $n/n$ and 2	1 <i>v</i> .	
 $\Delta \vartheta$	H''/H	Δϑ	H''/H	
400	0.61	800	0.64	
500	0.62	900	0.65	
600	0.63	1000	0.66	

Table 5. Relation between H''/H and  $\Delta \vartheta$ .

Since H' = H - H'', we can substitute H'/H in Eq. (5.1). For different values of the temperature  $\vartheta_g = \vartheta_0 + \Delta \vartheta$ , this yields the simplified expression

$$Q_{in} = K \cdot A \sqrt{H}$$

(5.2a)

where K varies with the temperature of the combustion gases,  $\vartheta_g$ , in conformity with Table 6, where  $\vartheta_0$  was put equal to zero.

	Table 6. Relation between K and $\vartheta_g$ .						
$\vartheta_g$	K	θ <sub>g</sub>	K				
400	0.40	900	0.38				
500	0.40	1000	0.38				
600	0.40	1100	0.37				
700	0.39	1200	0.37				
800	0.39						

4

In the case under consideration, the rate of flow of the incoming air,  $Q_{in}$ , was equal to the quantity of air supplied per unit time by the fan system. Hence it follows that this quantity of air can be replaced by a fictitious air flow factor,  $(A\sqrt{H})_{fict}$ , by means of Eq. (5.2a).

If the constant K is put equal to its value at  $600^{\circ}$ C, then we obtain

$$Q = 0.40(A/\overline{H})_{\text{fict}} \tag{5.2b}$$

By substituting  $A_t = 75 \text{ m}^2$ , we get

$$(A/H/A_t)_{\rm fict} = 0.033Q$$
 (5.3)

This relation between the fictitious opening factor,  $(A\sqrt{H}/A_t)_{\text{fict}}$ , and the rate of air flow, Q, shall include the effect of the air which may possibly enter into the tunnel through the paraboliform outlets for combustion gases, 0.5 m<sup>2</sup> in cross-sectional area each. In the calculations, these outlets were assumed to be replaced by two rectangular openings having a base B=1.13 m and a height H=0.44 m each.

If the rate of flow of the incoming air through these openings is put equal to  $Q_{in}$ , then  $Q_{in}$  at a maximum rate of combustion is given by the equation

$$Q_{\rm in} + Q\rho_0 = LR_{\rm max}\rho_0$$

For notations in this equation, see Chapter 3. By substituting  $Q_{in}$  from Eq. (3.4b), we find

$$2/3 \ \sqrt{2gH'} \ \frac{\rho_0 - \rho_g}{\rho_0} \ \mu B \cdot H' \rho_0 + Q \rho_0 = L R_{\max} \rho_0$$
(5.4)

where  $R_{\text{max}}$  is determined from Eq. (3.5a).

$$Q_{\rm out} = G_0 \cdot \rho_0 \cdot R_{\rm ma},$$

or

$$2/3 \cdot \left| \right/ 2gH'' \frac{\rho_0 - \rho_g}{\rho_g} \cdot \mu \cdot B \cdot H'' \cdot \rho_g = G_0 \cdot \rho_0 \cdot R_{\max}$$
(5.5)

By substituting B = 1.13 m, H = 0.44 m,  $\mu = 0.7$ ,  $G_0 = 4.86$  Nm<sup>3</sup> · kg<sup>-1</sup>, L = 3.98 Nm<sup>3</sup> · kg<sup>-1</sup> and  $R_{max}$  from Eq. (5.5) into Eq. (5.4) we obtain

$$0.83(H'/H)^{3/2} + 0.722Q = 0.38(H''/H)^{3/2}$$
(5.6)

In the calculation of  $\rho_g$ , the temperature of the combustion gases was assumed to be 600°C.

For those values of Q which are of interest in this connection, i.e.  $Q \ge 0.7$  m<sup>3</sup> · s<sup>-1</sup> in conformity with Table 4, Eq. (5.6) has no solution in the interval  $0 \le H'/H \le 1$ . This means that the flow is possible in an outward direction only.

In the test No. 6, the fans were switched off, and the exchange of air took place through the outlets for combustion gases alone. If these outlets are supposed, as before, to be approximately represented by rectangular openings having a base of 1.13 m and a height of 0.44 m each, then we obtain an opening factor  $A\sqrt{H}/A_t = 0.0088$  m<sup>1/2</sup>, which corresponds to a rate of air flow  $Q_{\min} = 0.26$  m<sup>3</sup> · s<sup>-1</sup> according to Eq. (5.3). Cf. the average value of the rate of air flow, 0.19 m<sup>3</sup> · s<sup>-1</sup>, which has been computed by Ödeen for the whole process of fire development.

Thus, in the tests under consideration, the fictitious opening factor,  $(A\sqrt{H}/A_t)_{\text{fiet}}$ , and hence the maximum rate of combustion, are determined approximately in conformity with the above from Eq. (5.3).

For the 16 tests comprised in this series, Table 7 gives, first, the opening factor calculated by means of the above relations, and second, the opening factor,  $(A \cdot \sqrt{H}/A_t)_{exp}$ , that has proved to give theoretical results which are in agreement with the experimental values. The test results, the theoretical results, and the time graphs of the rate of combustion are shown in Appendix 3. As regards the time graphs of the combustion gas temperature, it is to be noted that the six full-line curves represent the temperatures at different

Test No.	Fire load, <i>M</i> , kg of wood	Hydraulic radius, r, cm	Rate of air supply, Q, $m^3 \cdot s^{-1}$	Opening factor, theoretical value, $(A\sqrt{H}/A_t)_{theor}$ calculated from the formula $A\sqrt{H}/A_t = 0.0334 \cdot Q$	Opening factor, experimental value, $(A\sqrt{H}/A_t)_{exp}$	$\frac{(A\sqrt[3]{H}/A_t)_{exp}}{(A\sqrt[3]{H}/A_t)_{theorem}}$
4	675	1.0	0.7	0.023	0.035	1,52
8	675	1.7	0.7	0.023	0.015	0.65
10	270	1.0	0.7	0.023	0.015	0.65
11	405	1.0	0.7	0.023	0.020	0.87
12	405	2.4	0.7	0.023	0.012	0.52
13	135	1.0	0.7	0.023	0.005	0.22
16	405	0.4	0.7	0.023	0.060	2.60
1	270	-	1.0	0.033	0.023	0.70
3	675	1.0	1.0	0.033	0.043	1.30
7	675	1.7	1.0	0.033	0.037	1.12
9	675	0.6	1.0	0.033	0.051	1.54
14	945	1.6	1.0	0.033	0.037	1.12
5	675	1.0	1.5	0.050	0.055	1.10
2	675	1.0	2.0	0.067	0.060	0.90
15	1350	1.4	2.0 about	0.067	0.060	0.90
6	675	1.0	0.25	0.009	0.010	1,11

Table 7. Theoretical and experimental values of the opening factor.

Test series C [18].

points in the enclosed space. The fine dash-line curve summarises the values recorded in the radiation measurements, and the heavy dash-line curve is the calculated curve. The variation in the rate of combustion with the time,  $I_c$ , is represented in terms of  $330 \cdot (A\sqrt{H})_{exp} \cdot 2575 \text{ kcal} \cdot h^{-1}$  put equal to unity. The rate of flow of the incoming air, Q, conveyed by the fan system was constant during the whole process of fire development. The radiation measurements have provided certain indications for choosing the instant at which the rate of combustion had decreased to zero. The difference between the temperature at the level of the floor surface and the average temperature in the other parts of the enclosed space has been taken into account. In the test series under review, this difference in temperature has probably been increased owing to the fact that air was supplied to the enclosed space by means of fans at the floor surface level. The temperature difference was taken into consideration by assuming that the coefficient of heat transfer at the floor surface was equal to 80 per cent of the corresponding coefficient for the other surfaces. In all cases when the rate of burning was controlled by the fuel bed and not by the ventilation it was taken into account that heat energy was withdrawn from the enclosed space by that part of the incoming air which did not take part in the combustion.

As may be seen from Table 7, the positive as well as negative differences between the value of the opening factor,  $(A\sqrt{H}/A_t)_{exp}$ , determined from the test results, i.e. the actual maximum values of the rate of combustion, and the corresponding values obtained on the assumption that the rate of combustion is limited by the rate of air supply, were found to be great in some tests. However, the calculated curves were as a rule closely in agreement with the observed values. The agreement between the maximum temperature and the duration of the flame phase indicates that those values of the quantity of energy released per unit time which were used for the theoretical calculations were on the whole correct.

Furthermore, Table 7 shows two other factors among those which, in addition to the air flow factor, determine the rate of combustion. These factors are the hydraulic radius and the amount of fuel (the fire load). The effect of the first-mentioned factor can be demonstrated, for instance, by a comparison between the tests. Nos 16 and 12. The values of the air flow factor, as well as those of the fire load, in these two tests were equal, whereas the respective values of the hydraulic radius were 0.4 and 2.4 cm. In consequence of the difference in the hydraulic radius, the actual maximum rate of combustion in the test No. 16 was about 5 times as high as in the test No. 12. This may roughly be explained by the simplified study of the mechanism of combustion in what follows.

In a wood fuel consisting of comparatively large pieces of wood, pyrolysis takes place in several forms at the same time. In a certain definite inner zone of the wood, where the temperature is relatively low, say, below about  $250^{\circ}$ C, the reactions are endothermic, whereas in the outer zones, where the temperature is higher, the reactions are exothermic, and in certain cases, e.g. in the secondary pyrolysis of tar products, markedly exothermic. The combustion of the products of pyrolysis generates heat, which increases the temperature of the fuel by conduction, and hence renders possible an exothermic decomposition in the inner zone. In the test No. 16, the wood fuel consisted of concrete form timber, for which the ratio of the volume to the exposed surface was 0.4 cm. Accordingly, if the width, the length, and the thickness of a piece of wood are denoted by *b*, *l*, and *t*, respectively, then we have

$$r = 0.4 \text{ cm} = \frac{b \cdot l \cdot t}{2l(b+t)}$$

If we put b=t (square cross section), then we obtain a thickness of 1.6 cm, and if we set  $b \ge t$ , then we get a thickness of 0.8 cm. In view of the small thickness, in combination with the mechanism of heat return to the fuel described in the above, it is probable that a few minutes after the fuel has taken fire the whole quantity of fuel is in a state of active exothermic pyrolysis. Tests [23] have shown that the progression of the charred layer on a wooden beam exposed to fire is about 0.6 mm  $\cdot$  min<sup>-1</sup>. This value is applicable to a firwood beam exposed to fire in conformity with a standard temperature-time curve. The variation in the rate of carbonization with the intensity of the process of fire development is a problem which appears to be wholly unexplored at the present time. But if the above-mentioned value, 0.6 mm  $\cdot$  min<sup>-1</sup>, is assumed to be correct, then this implies that the whole amount of fuel would be charred in the course of 5 to 10 min. Since from one half to two thirds of the total quantity of energy is liberated during the flame phase, this can explain the intense release of energy immediately after the fuel has taken fire.

The test No. 12 shows that when the fire-exposed surface area diminishes below a certain definite limit, the quantity of energy which can be developed per unit time is determined by the rate of progression of the charred layer, and not by the air flow factor. In this test, the hydraulic radius was 2.4 cm. If the above-mentioned value, 0.6 mm  $\cdot$  min<sup>-1</sup>, which is probably too high in view of the low temperature during the process of fire development in this case, is used as a measure of the progression of the charred layer, then this corresponds to a maximum rate of combustion of  $405 \cdot 0.6 \cdot 10^{-2}/0.24 =$  $10.1 \text{ kg} \cdot \text{min}^{-1}$ . The value of the maximum rate of combustion computed from the formula  $R = 5.5A \sqrt{H}$  is 9.7 kg  $\cdot \text{min}^{-1}$ , and this implies that the quantity of energy released per unit time during a fire is determined by the surface area exposed to fire, and not by the rate of air supply, at least during certain phases of the fire.

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Furthermore, a closer study of Table 7 also demonstrates the marked effect produced by the amount of fuel on the rate of combustion. This can likewise be explained by means of the mechanism of return of the heat evolved by combustion which has been outlined in the above. In the tests Nos. 13, 10, 11, and 4, the values of the hydraulic radius, r, were equal, just as those of the rate of air supply Q, whereas the fire load, M, was varied. For the values of the fire load M=135, 270, 405, and 675 kg, which approximately corresponded to 7, 15, 20, and 35 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area of the enclosed space, the respective values of the maximum rate of combustion calculated from the test results were found to be equal to 0.22, 0.65, 0.87, and 1.5 times the rate of combustion which was computed on the basis of the air flow factor. In [18], it is stated that the flame phase was slightly developed in the tests Nos. 10 and 11, whereas the fuel did not take fire at all in the test No. 13.

Moreover, it is seen from Table 7 that the variation in the rate of combustion with the hydraulic radius seems to decrease as the opening factor or the fire load increases.

Finally, it may be useful to touch on the question to what extent these analytic fire tests, which were carried out under conditions that were idealized so far as possible, and which exhibited characteristics of combustion that in several tests markedly differed from those predicted by the theory, can be utilized as a basis for predicting the behaviour of more conventional fires. An ordinary fire in an enclosed space is governed to a varying degree by two feed-back mechanisms. In the first place, the lower density of the combustion gases forces unconsumed air by natural convection towards the flames, and hence increases the rate of combustion, as well as the evolution of combustion gases. In the second place, part of the heat generated by combustion returns to the fuel, and increases the rate of energy release. So long as these mechanisms are negative, the combustion remains stable. Cf. e.g. the test No. 13. An essential difference between Ödeen's tests described in [18] and an ordinary fire is that the return mechanism which governs the rate of air supply was eliminated in the tests. The effect of this circumstance is difficult to determine, but it may be mentioned for comparison that all the theoretical and experimental results in the test series (A, B, and D) where the exchange of air was self-regulated were found to be closely in agreement if the rate of combustion was assumed to be determined by the rate of air supply. However, a comparison of the time graphs of the rate of combustion for the test series A to D shows that, in the cases where differences were present, the results obtained from the test series C deviated from those of the other test series only in respect of the maximum quantity of energy released per unit time. If the time graphs of the rate of combustion are represented in terms of  $330 \cdot (A \sqrt{H})_{exp} \cdot 2575$ kcal · h<sup>-1</sup> put equal to unity, then the curves for this test series are found

to be closely in agreement with those for the test series A, B, and D when the fire load and the opening factor are given. All the same, since a dispersion in the values of the quantity of energy liberated per unit time has been observed in the test series C even when the tests were identical in respect of the rate of air supply and the wall material, it should be noted that this dispersion indicates the need for determining the factors which, in addition to those mentioned in the above, govern the rate of release of energy.

## Test series D. Tests made at the National Swedish Institute for Materials Testing. Calculation of time graphs of rate of combustion

Under the direction of the Fire Engineering Laboratory of the National Swedish Institute for Materials Testing, a test house for model-scale and fullscale fire studies has been erected at the Studsvik Test Station of the Atomic Energy Co., Ltd. The primary object of the investigations carried out in this house was to study the spread of fire and smoke along the exterior walls and along the ventilation ducts in the case of fire in an individual enclosed space in a multi-storeyed building. Extensive measurements of the reduction in the weight of fuel, the temperature, the intensity of radiation, the gas flow, and the composition of combustion gases were made in these tests, and the test results are therefore well suited for theoretical comparisons.

Fig. 12 shows the test house, which was three storeys high, and which consisted of a load-bearing steel frame clad with lightweight concrete elements. The results of the tests carried out up to now have not yet been published, but the test programme and the test equipment are described in [20] and [21]. The Authors of the present publication were afforded an opportunity to acquaint themselves with the results of the first four full-scale tests. The data for these tests are reproduced in Table 8.

The fires were initiated in the lowermost storey, which was connected with the outside air by means of a vertical ventilation duct by-passing the storeys

			Table 8. Te	st series D.			
Test No.	Fire load, kg of wood	Window dimensions, width×height m×m	Bounding surface area, m <sup>2</sup>	Air flow factor, $A \cdot \sqrt{H}$ , $m^{5/2}$	Air flow factor, fictitious value, $(A \overline{H})_{\text{fict}}$ m <sup>5/2</sup>	Moisture content of fuel, per cent	Moisture content of lightweight concrete, per cent
1	350	1.3×1.3	75	1.93	2.10	7.7	4.0
2	200	$1 \times 1$	75	1.00	1.10	8.8	3.1
3	115	$0.8 \times 0.8$	75	0.57	0.69	7.4	2.1
4	1150	$2 \times 2$	75	5.63	5.65	9.7	2.3



Fig. 12. Photograph of the test house used in the test series D to study the spread of fire and smoke along the external walls and through a vertical ventilation duct in the case of fire in a certain definite storey in a multi-storey building.

situated above. Calculations showed that the effect of the ventilation duct on the air flow factor was not to be disregarded, and a fictitious air flow factor,  $(A\sqrt[3]{H})_{\text{fict}}$ , was therefore computed in accordance with the principles stated in Chapter 3. The magnitude of the corresponding correction can be seen from Table 8. The time graphs of the temperature of the combustion gases which were obtained from these tests, and which are expressed in terms of the mean value of the temperatures observed during the tests at 21 points in the enclosed space, are reproduced in Appendix 4. In order to give an idea of the dispersion about this mean value, the diagram relating to the test No. 2 shows the temperature-time curves for a point situated 45 cm above the floor and a point located 45 cm below the ceiling. Furthermore, the four diagrams in Appendix 4 also comprise the calculated curves which represent the variation in the temperature of the combustion gases with the time, as well as the time graphs of the rate of combustion in a dimensionless form which were used in the calculations.

It proved possible to bring the calculated and observed curves into close agreement by choosing the time graphs of the rate of combustion which were similar in shape in all the tests, and were based on a maximum rate of combustion  $I_c = 330 \cdot A \sqrt{H} \cdot 2575$  kcal  $\cdot h^{-1}$ . The value of the air flow factor used in the calculations was the value  $(A \sqrt{H})_{\text{fict}}$ , which was corrected so as to take account of the effect of the ventilation duct.

#### Summary

All the comparative calculations dealt with in the present chapter were based on the assumption that the energy conditions during the process of combustion can be characterized by an ignition phase in which the quantity of energy released per unit time increases from a zero value in accordance with a polygonal function of the time to a value that is given by the air flow factor. This phase is followed by a flame phase, during which the rate of combustion was supposed to be constant. After that the rate of combustion decreases to zero as a polygonal function of the time in the course of the cooling phase, during which the slopes of the individual sides of the polygon vary in a marked manner with the fire load. The higher the fire load, the slower the decrease in the rate of combustion. Of course, these assumptions give a simplified picture of the variation in the liberation of energy per unit time during the process of fire development. Thus, for most types of fire loads, it is to be expected that the plane part of the curve which represents the rate of combustion during the flame phase is rather to be regarded as the mean value of the quantity of energy released per unit time. This is illustrated in Fig. 13, which represents the variations in the observed rate of combustion (expressed in terms of the rate of reduction in the weight of fuel) and in the oxygen content of the combustion gases during the test No. 1 in the test series A. On account of technical difficulties in measurements, the values of the rate of reduction in the weight of fuel were somewhat uncertain immediately after the fuel had taken fire, as the rate of combustion was then liable to very wide variations. However, these values were confirmed by the fact that the oxygen content of the combustion gases in the enclosed space exhibited corresponding variations. Even when the curve which represents the variation in the rate of reduction in the weight of fuel with the time is known, our present knowledge of the relation between the rate of reduction in the weight of fuel and the rate of release of energy during the different phases of the process of fire development does not make it possible to determine the quantity of energy liberated per unit time during combustion. For the test No. 3 in the test series D, Fig. 14 shows three theoretical temperature-time curves calculated on the basis



Fig. 13. Variation in the rate of combustion, in kg  $\cdot$  min<sup>-1</sup>, with the time, *t*, determined by measuring the reduction in the weight of fuel in one of the tests comprised in the test series A, see Chapter 5. Furthermore, this figure also shows the time graph of the oxygen content of the combustion gases in the same test (dash-line curve). During the interval from the 55th to the 60th minute, the weighing of the fuel was disturbed by the fact that parts of the ceiling of the enclosed space fell down.



Fig. 14. Time graphs of the temperature of the combustion gases calculated for three different heat values of the fuel, viz., W=2575, 3000, and 4230 kcal  $\cdot$  kg<sup>-1</sup> of wood, on the basis of an observed time graph of weight loss of fuel (inset). An experimental temperature-time curve is represented by the dash-line curve.

of a measured rate of weight loss. It was assumed that the heat value W corresponding to 1 kg of weight loss was constant during the whole process of fire development and equal to 2575,3000 and 4230 kcal, respectively. The dash line curve represents the values observed in the test. The observed rate of reduction in the weight of fuel is represented in a separate diagram in fig. 14. As seen from this diagram, the maximum rate of reduction in the weight of fuel is closely in agreement with the theoretical value  $R = 5.5 \times A \sqrt{H} = 3.80 \text{ kg} \cdot \text{min}^{-1}$  (horisontal dash line). It is also seen from fig. 14 that the assumption of a constant value of W obviously is incorrect. A comparison between the curves indicates that the quantity of energy released per unit weight of fuel, at least during the first part of the cooling phase is greater that that during the flame phase.

Accordingly, the relative values of the rate of combustion obtained from the above comparative calculations cannot be directly expressed in terms of the rate of reduction in the weight of fuel, in kg  $\cdot$  min<sup>-1</sup>. If the quantity of energy liberated per unit time is expressed directly in kcal  $\cdot$  min<sup>-1</sup>, then this obviates the difficulty of determining the heat of combustion of the wood fuel during the various phases of the process of fire development.

A prerequisite in most of the calculations has been that the fire process is controlled by ventilation. This condition is far from being generally realized when it comes to actual fires. As a rule, a combination of small ventilation area and high fire load gives a process of fire development where the rate of burning is proportional to the air flow factor. If the fire load is low and the ventilation area large, the combustion will proceed as if in the open. This means that the rate of burning depends on the fire load density (fire exposed surface) and that an increase in ventilation will not result in a corresponding increase in rate of burning.

It is, however, impossible to say in advance if the process of fire development will be ventilation controlled or not even if the air flow factor and the fire load are known. The orientation and the distribution of the fuel in the enclosed space and the thickness or the porosity of the fuel will be a decisive factor in each particular case. An assumption that the rate of burning is determined by the air flow factor ought to give time-temperature curves which are on the safe side in practically every case. If such an assumption is made and the combustion in spite of this happens as in the open, i.e. is fuel bed controlled, the result will be a fire process of lower maximum temperature and, at least in some cases, of longer duration. The longer duration will not increase the severity of the fire to a corresponding degree. This is due to the fact that part of energy released by the combustion will be withdrawn from the enclosed space by the surplus air. In this way the temperature of the combustion gases will be lower compared to the case when the duration is the same but the process controlled by the ventilation. To sum up, a comparative theoretical analysis of the results obtained from some thirty full-scale fire tests of the wood fuel type has been carried out in Chapter 5. The calculations made for this purpose covered relatively wide variations in fire load, opening factor, and hydraulic radius, as well as in the thermal properties of the structures bounding the enclosed space. As a result of these calculations, the time graph of the quantity of energy released per unit time may be assumed to be known within this range of variation.

## 6. Determination of general time graphs of quantity of energy released per unit time during different phases of process of fire development

In order to afford a basis for the calculation of the curve which represents the variation in the temperature of the combustion gases with the time during the process of fire development under varying conditions, it is necessary to systematize the time graphs of the rate of combustion which have been obtained in Chapter 5. A detailed investigation has been made of these graphs in order to find out how they vary with the fire load and with the opening factor. This investigation indicated the possibility of the simplification outlined in what follows.

If the ratio of the fire load, which is given from the outset, to the air flow factor,  $A\sqrt{H}$ , is constant, that is to say, if the duration of the fire is constant, then the time graph of the rate of combustion, expressed in a relative form in terms of the maximum rate of combustion,  $330 \cdot A\sqrt{H} \cdot 2575 \text{ kcal} \cdot h^{-1}$ , put equal to unity, is independent of the opening factor. This implies, for instance, that an enclosed space where the opening factor is  $A\sqrt{H}/A_t = 0.01 \text{ m}^{1/2}$  and the fire load is  $q=5 \text{ Mcal} \cdot \text{m}^{-2}$  can be characterized by the same graph of the rate of combustion as an enclosed space where the opening factor is  $0.04 \text{ m}^{1/2}$  and the fire load is  $20 \text{ Mcal} \cdot \text{m}^{-2}$ . Accordingly, the results of the calculations in Chapter 5, which show how the quantity of energy released per unit time varies with the time, can be represented by a graph which comprises a separate curve for each value of the ratio  $qA_t/A\sqrt{H}$ . In this connection, it is convenient to introduce the duration of the fire, T, defined as the duration of the flame phase, cf. Eq. (1.2),

$$T = qA_t/(1500A \cdot \sqrt{H})$$
 h

as the variable at which the graph shall be entered. In this formula for calculating the duration of the flame phase, the product of the constant 330 in the expression  $R=330 \cdot A\sqrt{H}$  kg  $\cdot$  h<sup>-1</sup> and the heat value of the wood fuel, i.e. 4.5 Mcal  $\cdot$  kg<sup>-1</sup>, has been put equal to 1500. For the values of the duration of the fire defined in this way, T=0.1, 0.2, 0.3, 0.5, 0.75, 1.0, 1.5 and 2.0 h, which correspond to the respective fire loads (150, 300, 450, 750, 1125, 1500, 2250, and 3000)  $\cdot A\sqrt{H}/A_t$  Mcal  $\cdot$  m<sup>-2</sup> of the total surface area bounding

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the enclosed space, Fig. 15 shows the variation in the rate of combustion with the time.

In order to make it easier to check the agreement between the total quantity of energy liberated during the process of fire development and the fire load, which is given from the beginning, and since a more accurate representation would be illusory, considering the character of the available data, the curve form has been assumed to be polygonal, just as in the comparative calculations in Chapter 5. In Fig. 15,  $330A\sqrt{H} \cdot 2575$  kcal  $\cdot$  h<sup>-1</sup> has been put equal to unity. The respective areas between the above-mentioned curves and the axis of time shall therefore be  $\frac{1}{330 \times 2575}$  (150, 300, 450, 750, 1125, 1500, 2250, and 3000) area units. In the relevant Swedish regulations, the quantity *T*, determined by the relation



Fig. 15. Time graphs of the energy released per unit time in the process of combustion,  $I_c$ , expressed in a relative form by putting  $330 \cdot 2575 \cdot A \sqrt{H}$  equal to unity. The eight curves shown in this figure correspond to different values of the duration of the fire defined as the duration of the flame phase by the expression  $T = q \cdot A_t / (1500 \cdot A \cdot \sqrt{H})$ . The dash-line portion of the curve for the ignition phase belongs to the curves relating to the lowest four values of the duration of the fire.

## $T = qA_t/(25A\sqrt{H})$ min

designates the instant which marks the end of the flame phase and the beginning of the linear cooling phase. On the basis of the experiences derived from the comparative theoretical analyses, the instant at which the rate of combustion begins to decrease has been chosen so as to be slightly anterior to the instant defined by T. For the values of the duration of the fire T=6, 12, 18, and 30 min, the time graphs of the rate of combustion during the ignition phase have been given a slightly different shape, which implies that the fuel takes fire within a shorter period of time.

## 7. Calculation of time graphs of temperature of combustion gases for characteristic types of enclosed spaces varying in opening factor and in fire load

The time graphs of the rate of combustion represented in a relative form in Fig. 15 for fires of the wood fuel type in enclosed spaces are utilized in the present chapter as a basis for the calculation of complete time graphs of the temperature of the combustion gases. This is done for varying values of the opening factor and the fire load in enclosed spaces of the seven types dealt with in what follows, which differ in respect of the bounding structures.

### Type A enclosed space

Bounding structures.

All the surfaces which bound the enclosed space are supposed to consist of a material, 20 cm in thickness, whose thermal properties are characterized by the average values given below, which apply to structural materials of such types as concrete, brick, and lightweight concrete.

Thermal conductivity,  $\lambda = 0.7 \text{ kcal} \cdot \text{m}^{-1} \cdot \text{h}^{-1} \cdot \text{°C}^{-1}$ .

Product of the specific heat and the weight per unit volume,

 $c \cdot \gamma = 400 \text{ kcal} \cdot \text{m}^{-3} \cdot \text{°C}^{-1}$ .

The same data on the properties of materials had also been used for the calculation of those temperature-time curves for the flame phase of the process of fire development which have been published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures".

## Type B enclosed space

Bounding structures.

Concrete, 20 cm in thickness. Thermal conductivity,  $\lambda = 1.4 \cdot e^{-0.001 \cdot \vartheta} \text{ kcal} \cdot \text{m}^{-1} \cdot \text{h}^{-1} \cdot \text{°C}^{-1}$  [1]. Enthalpy, *I*, see Fig. 17.

## Type C enclosed space

Bounding structures.

Lightweight concrete, 20 cm in thickness. Weight per unit volume,  $\gamma = 500 \text{ kg} \cdot \text{m}^{-3}$ . Thermal conductivity,  $\lambda$ , see Fig. 16.

Enthalpy, I, see Fig. 17.

The specific heat and the weight per unit volume of the lightweight concrete are assumed to be independent of the temperature. Consequently, the enthalpytemperature curve is rectilinear. The variation in the thermal conductivity,  $\lambda$ , with the temperature is based on a determination which has been made in connection with the test series D described in Chapter 5.



Fig. 16. Relations between the thermal conductivity,  $\lambda$ , and the temperature,  $\vartheta$ , used in the calculations for brick, diabase wool, plasterboard, and lightweight concrete.



Fig. 17. Relations between the enthalpy, I, and the temperature,  $\vartheta$ , used in the calculations for steel, concrete, brick, plasterboard, and lightweight concrete.

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## Type D enclosed space

Bounding structures.

Concrete, 50 per cent of the total bounding surface area.

Lightweight concrete, 50 per cent of the total bounding surface area. Thicknesses, weight per unit volume, and thermal properties as in the Type B and Type C enclosed spaces, respectively.

### Type E enclosed space

Bounding structures.

Lightweight concrete, 50 per cent of the total bounding surface area. Thickness, weight per unit volume, and thermal properties as in the Type C enclosed space.

Concrete, 33 per cent of the total bounding surface area.

Thickness and thermal properties as in the Type B enclosed space. Other structural components, 17 per cent of the total bounding surface area, enumerated in the order from the interior to the exterior:

Plasterboard panel, 13 mm in thickness.

Weight per unit volume,  $\gamma = 790 \text{ kg} \cdot \text{m}^{-3}$ .

Diabase wool, 10 cm in thickness.

Weight per unit volume,  $\gamma = 50 \text{ kg} \cdot \text{m}^{-3}$ .

Brickwork, 20 cm in thickness.

Weight per unit volume,  $\gamma = 1800 \text{ kg} \cdot \text{m}^{-3}$ .

Thermal properties of plasterboard, diabase wool, and brick, see Figs. 16, 17, and 18.

The enthalpy-temperature curve chosen for brick is based on a value of the specific heat which is supposed to be independent of the temperature. In reality, the specific heat of brick slightly varies with the temperature [1]. However, in the present case, the values of the temperature rise in the brickwork, which is most remote from the surface exposed to fire, are so low that



Fig. 18. Relation between the enthalpy, I, and the temperature,  $\vartheta$ , used in the calculations for diabase wool.

the effect of this variation may be disregarded. The variation in the enthalpy of diabase wool with the temperature is based on that relation between the specific heat and the temperature which was published in [1]. The thermal properties of plasterboard are based on the curves which were published in [19]. It is assumed that the plasterboard panels will not fall down or disintegrate.

The values of the thermal conductivity,  $\lambda$ , of brick and diabase wool were taken from [1].

#### Type F enclosed space

Bounding structures.

Sheet steel, 2 mm in thickness, 80 per cent of the total bounding surface area.

Concrete, 20 cm in thickness, 20 per cent of the total bounding surface area.

Thermal properties as in the Type B enclosed space.

Curves representing the variations in the enthalpy and in the thermal conductivity of sheet steel with the temperature, see Figs. 17 and 19, respectively.

This type of enclosed space corresponds to a storage space, or the like with a sheet steel roof, sheet steel walls, and a concrete floor.

#### Type G enclosed space

Bounding structures.

Concrete, 20 per cent of the total bounding surface area.

Thickness and thermal properties as in the Type B enclosed space.

Other structural components, 80 per cent of the total bounding surface area, enumerated in the order from the interior to the exterior:

Two plasterboard panels,  $2 \times 13$  mm in thickness.



Fig. 19. Relation between the thermal conductivity,  $\lambda$ , and the temperature,  $\vartheta$ , used in the calculations for unalloyed steel.

Weight per unit volume,  $\gamma = 790 \text{ kg} \cdot \text{m}^{-3}$ . Cavity, 10 cm in width.

Two plasterboard panels,  $2 \times 13$  mm in thickness.

Thermal properties of plasterboard, see Figs. 16 and 20.

This structure represents a type of partition which is becoming more and more common, and which consists of two plasterboard panels on each side, supported on steel stud framing. It is assumed that the steel studs have no thermal conductivity and no thermal absorptivity.

The test results published in [19] have shown that plasterboard panels which are not fibre-filled disintegrate when their temperature on the side that is not exposed to fire reaches about  $550^{\circ}$ C. However, this does not apply to the outermost, i.e. the fourth, plasterboard panel. This panel is in contact with the air, which has a temperature of  $20^{\circ}$ C. Therefore, this panel never reaches a surface temperature of  $550^{\circ}$ C. In fact, tests have demonstrated that a plasterboard panel in this position disintegrates when the temperature at its centre rise to about  $750^{\circ}$ C. These criteria have been used in calculating the time graph of the combustion gas temperature for enclosed spaces of the type in question. The calculations were discontinued when they had been carried out to the instant at which the fire was expected to burn through the wall, that is to say, after all four plasterboard panels had disintegrated.

As may be seen from Fig. 20, which was taken from [19], the variation in the enthalpy of plasterboard with the temperature is dependent on whether



Fig. 20. Relation between the enthalpy, *I*, and the temperature, θ, used in the calculations for plasterboard. Curve 1: High rate of temperature rise. Curve 2: Low rate of temperature rise.

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the rate of temperature rise is high or low. This circumstance was taken into account in the calculations by assuming a high rate of temperature rise for that plasterboard panel which is nearest to the fire and a low rate of temperature rise for all the other plasterboard panels. Since the structural transformations of plasterboard require an additional quantity of heat that cannot be recovered, the enthalpy-temperature curve during the cooling period is not identical with that during the heating period.

In the calculations, it was assumed that the variation in the enthalpy with the temperature during the cooling period is represented by a straight line which corresponds to a constant value of the product  $c \cdot \gamma = 200$  kcal × × m<sup>-3</sup> · °C<sup>-1</sup>.

For the surfaces exposed to fire in the enclosed spaces of all the types dealt with in the present chapter, the coefficient of heat transfer,  $\alpha_i$ , was calculated by means of Eq. (2.7a), where it was assumed that  $\varepsilon_{fi} = 0.7$  and  $\varepsilon_i = 0.8$  for the Type A to the Type E enclosed spaces, and for the Type G enclosed space. These values give  $\varepsilon_{res} \sim 0.60$ . For the Type F enclosed space, the calculations were based on three values of  $\varepsilon_i$ , viz., 0.1, 0.4, and 0.8. Hence, for  $\varepsilon_{fi} = 0.7$ , the respective values of  $\varepsilon_{res}$  were found to be 0.1, 0.35, and 0.6. For the exterior surfaces of the structures bounding the enclosed space, the coefficient of heat transfer,  $\alpha_u$ , for the Type A to Type E enclosed spaces, and for the Type G enclosed space, was supposed to vary with the surface temperature,  $\vartheta_u$ , in accordance with Eq. (2.8), while its value for the Type F enclosed space was supposed to vary in conformity with the relation

$$\alpha_{u} = 7.5 + \frac{4.96 \cdot \varepsilon_{\text{res}}}{\vartheta_{u} - \vartheta_{0}} \left[ \left( \frac{\vartheta_{u} + 273}{100} \right)^{4} - \left( \frac{\vartheta_{0} + 273}{100} \right)^{4} \right] \text{kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1}. \quad ^{\circ}\text{C}^{-1}$$
(7.1)

where  $\varepsilon_{\rm res}$  was chosen in the same way as in the above-mentioned calculation of  $\alpha_i$ .

The variation in the specific heat,  $c_p$ , of the combustion gases with the temperature is shown in Fig. 21 [3]. For the rest, the calculations were based on the principles which have been stated in Chapters 2 and 4.

The input values used in the calculations were only the opening factor  $A \cdot \sqrt{H}/A_t$ , in m<sup>1/2</sup>, and the fire load, q, in Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area. If the radiation term,  $I_R$ , is disregarded, then Eq. (2.1) becomes independent of dimensions. In other words, if the values of  $A \cdot \sqrt{H}/A_t$  and q are given, then the result will be independent of the terms A, H, and  $A_t$  entering into the opening factor. However,  $I_R$  is proportional to the total opening area, A, and, in order that  $I_R$  may be taken into account, it is necessary to specify the above-mentioned terms. This has been done in the calculations dealing with test results in Chapter 5. For the determination of the total

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Fig. 21. Relation between the specific heat of the combustion gases,  $c_p$ , and their temperature,  $\vartheta_q$ , used in the calculations.

opening area, A, in the calculations described in the present chapter, it was supposed that the dimensions of the enclosed space were the same as those which had been assumed as a basis for the curves published in the Swedish Building Regulations 1967, viz., a total bounding surface area  $A_t = 10,000 \text{ m}^2$ and a square opening. Consequently, the ratio of the total opening area to the total bounding surface area,  $A/A_t$ , was in all cases lower than those values which can be expected to be met with in ordinary buildings, and all the results will therefore be on the safe side. The value  $A_t = 10,000 \text{ m}^2$  was used only to determine the value of the ratio  $A/A_t$ , which was then substituted in the term  $I_R$  in the equation of heat balance. For a value of the opening factor  $A \cdot \sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$ , Table 9 gives the respective values of the ratio  $A/A_t$ which correspond to  $A_t = 10,000 \text{ m}^2$  and  $A_t = 1 \text{ m}^2$  on the assumption that the opening is square, or that it has a height H=1 m.

	Table 9. Values of the ratio $A/A_t$ .				
	Square opening	<i>H</i> =1 m			
$A_t = 1 \text{ m}^2$	0.075	0.04			
$A_t = 10,000 \text{ m}^2$	0.012	0.04			

In order to illustrate the consequence of this variation in the ratio  $A/A_t$ , Fig. 22 shows the temperature-time curve for an enclosed space characterized by an opening factor  $A \cdot \sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$ , a fire load  $q=30 \text{ Mcal} \cdot \text{m}^{-2}$ , as well as by the two extreme values of the ratio  $A/A_t$ , i.e. 0.012 and 0.075. As is seen from this graph, the effect of the difference in the value of the ratio  $A/A_t$  is practically negligible.

Temperature-time curves have been calculated for each one of the seven enclosed spaces, Types A to G, for varying values of the duration of the fire and the opening factor. As regards the relation between the duration of the fire and the fire load, reference is made to Chapter 6. For the Type A to the


Fig. 22. Calculated temperature-time curves for the Type A enclosed space. Opening factor  $A\sqrt[4]{H}/A_t = 0.04 \text{ m}^{1/2}$ . Fire load  $q = 30 \text{ Mcal} \cdot \text{m}^{-2}$  of bounding surface area. Ratio  $A/A_t$  set equal to 0.012 and 0.075, respectively.

Type E enclosed space, as well as for the Type G enclosed space, the temperature-time curves were computed on the basis of 6 different values of the opening factor, viz.,  $A \cdot \sqrt{H}/A_t = 0.01$ , 0.02, 0.04, 0.06, 0.08, and 0.12 m<sup>1/2</sup>. For each of these values the curves were computed for 8 different values of the duration of the flame phase of the fire for which the time graphs of the rate of combustion have been constructed in Chapter 6. This means that 48 temperature-time curves have been obtained for each type of enclosed space. In the case of the Type F enclosed space, the calculations were carried out for 3 different values of the resultant emissivity. Furthermore, for each one of these values, use was made of 5 different values of the opening factor, viz., 0.01, 0.02, 0.04, 0.08, and 0.12  $m^{1/2}$ . For each one of these combinations of values, the curves were calculated on the basis of 5 different values of the fire load corresponding to 5 different values of the duration of the flame phase, viz., 0.1, 0.3, 0.5, 1.0, and 2.0 h, if computed by means of Eq. (2.1). All these curves are shown in Appendix 5. The curves are denoted by the symbols A1 to G6, where the letter A refers to the Type A enclosed space, etc. All these curves are represented in an approximate form after smoothing out the irregularities which were caused by the polygonal shape of the time graphs of the rate of combustion. Furthermore, in order to render the graphs in Appendix 5 more readily legible, that part of each one of the curves which represents the ignition phase was based on that ascending branch of the time graph of the rate of combustion which corresponds to the lower four values of the fire load. The exact results of the calculations are reproduced in tabular form in Appendix 6. In the case of the Type G enclosed space, where the calculation of the temperature-time curves was carried out with reference to the instants when the plasterboard panels fell down, each one of these instants is marked with a circle on the corresponding curve. When a plasterboard panel falls down, this corresponds in the calculations to an instantaneous temperature drop in the enclosed space, as may be seen from the relevant curves.

In practical design, it should be possible to proceed in three steps, viz., first, to choose that type of enclosed space which is most closely similar in respect of the thermal properties of the bounding structures to the case under consideration; second, to determine the opening factor and the fire load; and third, to interpolate linearly between the values given in the tables in Appendix 6. If, instead of using this procedure, the designer chooses a curve which is determined without interpolation so as to be on the safe side, that is to say, if he chooses the next higher values of the opening factor and the fire load, then this will probably not involve errors which are too great. In order to afford a basis for the choice of the type of enclosed space, the temperature-time curves which correspond to an opening factor  $A \cdot \sqrt{H}/A_t = 0.04$  $m^{1/2}$  and to a fire load of 60 Mcal  $\cdot m^{-2}$  of bounding surface area are represented in Fig. 23 for the Type A to the Type F enclosed spaces. For comparison, this graph also reproduces the standard ISO temperature-time curve and the curve for an opening factor of  $0.04 \text{ m}^{1/2}$  published in the Swedish Building Regulations 1967.

A comparison between the temperature-time curves which correspond to the different types of enclosed spaces in Fig. 23 shows that the maximum difference in the maximum temperature amounts to about 400°C. The Type C enclosed space, which is provided with lightweight concrete bounding structures, exhibits markedly higher temperatures than the other types of enclosed spaces comprised in the present calculations. The lowest maximum temperature was obtained in the case of the Type F enclosed space ( $\varepsilon_{res} = 0.60$ ), which is equipped with bounding structures made of sheet steel, 2 mm in thickness. However, it is seen from Fig. 23 that the resultant emissivity for radiation between the flames and a sheet steel surface produces a substantial effect on the magnitude of the maximum temperature. If  $\varepsilon_{res}$  is supposed to change from 0.6 to 0.1, then the corresponding difference in the maximum temperature is slightly over 200°C, other conditions being equal. Therefore, it is important that the heat transfer conditions, and particularly the resultant emissivity, should be accurately determined in the calculation of the temperature-time curve for an enclosed space of this type.



Fig. 23. Temperature-time curves for the Type A to the Type F enclosed spaces. Opening factor  $A|\overline{H}/A_t=0.04 \text{ m}^{1/2}$ . Fire load  $q=60 \text{ Mcal} \cdot \text{m}^{-2}$  of bounding surface area. Furthermore, this graph also shows the respective temperature-time curves calculated by means of Eq. (1.1) and determined in conformity with the Swedish Building Regulations 1967 (SBR 67) for an opening factor  $A|\overline{H}/A_t=0.04 \text{ m}^{1/2}$ .

With the exception of the Type F, it is the Type B enclosed space, which is bounded by concrete walls, that exhibits the lowest maximum temperature. This is due to the relatively high thermal conductivity and the great heat capacity of the concrete. On the other hand, since the large quantity of heat that is stored in the concrete is partly transferred back to the enclosed space during the cooling period, comparatively high temperatures are obtained in the course of the cooling phase.

The temperature-time curves published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures" relate to "enclosed spaces bounded by wall, roof or ceiling, and floor structures which are made of brickwork, concrete, or lightweight concrete as a material that is predominant in thermal respects". As has previously been mentioned, these curves had been calculated on the basis of those characteristics of the bounding structures which were used to describe the Type A enclosed space. Moreover, for guidance, the comments on the Draft Specification "Aluminium Structures" also comprise temperature-time curves for an enclosed space which is bounded by walls made of mineral wool. In addition, it is shown how the temperature-time curve is influenced by a concrete wall, 20 cm in thickness, which is situated in the interior of the enclosed space. By examining Fig. 23, it will readily be understood that a further differentiation of the above-mentioned Swedish standard temperature-time curves according to the thermal characteristics of the bounding structures would be desirable.

In order that the temperature-time curves for the cooling phase which have been determined in the present publication might be compared with the corresponding Swedish standard curves, Fig. 24 shows the temperature-time curves for the Type A enclosed space calculated on the basis of an opening factor  $A \cdot \sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$ , together with the Swedish standard temperaturetime curves for the cooling phase determined on the assumption that the rate of temperature decrease is  $10^{\circ}\text{C} \cdot \text{min}^{-1}$ . The ascending branches of the curves are identical because the curves for the Type A enclosed space are based on the same assumptions as the standard curves. The linear temperature-time curves for the cooling phase, which start from the ascending branch at the time T, calculated by means of Eq. (1.2), are represented by dash lines in Fig. 24. As is seen from this graph, the calculated curves result in a rate of cooling which is higher or lower than the standard rate of temperature decrease,  $10^{\circ}\text{C} \cdot \text{min}^{-1}$ , according as the duration of the process of fire development is shorter or longer, respectively.

Thus, if the duration of the fire is comparatively short, then an application of the temperature-time curves which have been computed in the present publication gives considerably more favourable results, i.e. lower temperatures, than the standard rules which are at present in force in Sweden. For instance, the temperature-time curve for the cooling phase in fires of short duration is a decisive factor in determining the temperatures of unprotected steel structures, as has already been shown in an example which was adduced under



Fig. 24. Temperature-time curves for the Type A enclosed space. Opening factor  $A \gamma \overline{H} / A_t = 0.04 \text{ m}^{1/2}$ . (Full-line curves.) Curves for the cooling phase corresponding to a rate of decrease in temperature of  $10^{\circ}\text{C} \cdot \text{min}^{-1}$  stipulated in the relevant Swedish regulations. (Dash-line curves.)

the heading "Introduction". However, if the duration of the fire is relatively long, then the calculated temperature-time curves are less favourable, i.e. they give higher temperatures, than the curves stipulated in the Swedish regulations. All the same, since a given structure that is exposed to fire has already been subjected to the action of high temperatures for a long time before the beginning of the cooling phase, the difference between these curves in the latter case will have a comparatively slight effect on practical design.

## 8. Summary

The point of departure of the present investigation was the fact that the results obtained in recent years from research in the field of structural fire engineering had made it possible to carry out reliable calculations of the load bearing and separating capacity in the design of structural components exposed to fire. Such design calculations must be based on the knowledge of the temperature-time curve which covers the whole process of fire development. However, further progress towards realistic structural fire engineering design was impeded by the circumstance that it was not possible to make a theoretical determination of the temperature-time curve for the cooling phase of the process of fire development under known external conditions. Up to now, the research in this field has evolved methods for calculating the variation in the temperature with the time during the flame phase of the process of fire development on the basis of known external conditions, whereas the cooling phase has not been dealt with in this connection.

In consequence of this gap in our knowledge, the methods for determining the temperature-time curves for the flame phase and the cooling phase stipulated in the Swedish Building Regulations 1967 are widely different in degree of accuracy. The determination of the temperature-time curve for the flame phase shall be based on the fire load which characterizes the case under consideration, as well as on the shape and the dimensions of the openings in the enclosed space. For the cooling phase, on the other hand, it is stipulated only that the rate of temperature decrease shall be set equal to  $10^{\circ}C \cdot min^{-1}$ , irrespective of the actual conditions which characterize the case in question. This undifferentiated characterization of the cooling phase is particularly unfavourable to structures which possess a low thermal inertia, e.g. noninsulated or slightly insulated load-bearing steel structures. It was therefore considered to be urgently required to undertake an investigation in order to find out whether a theoretical determination of the temperature-time curve for the cooling phase would be possible.

The theoretical calculations in the present publication are founded on a basic equation of heat balance in an enclosed space which has been deduced by Kawagoe and Sekine, as well as by Ödeen. This equation states that the quantity of heat,  $I_c$ , which is released per unit time during the process of combustion is at any instant equal to the sum of the quantities of heat which

are withdrawn per unit time in different ways from the enclosed space. Heat is ordinarily abstracted from the enclosed space by heat transfer through the structures which bound the enclosed space (term  $I_W$  in the equation of heat balance), by radiation through the openings in the enclosed space (term  $I_R$ ), and by the replacement of combustion gases by cold air (term  $I_L$ ).

In order to extend the range of application of the equation of heat balance so that it might cover the whole process of fire development, it was necessary to solve two fundamental problems. In the first place, the quantity of heat released per unit time had to be determined as a function of the time for the entire process of fire development. In the second place, the expression for  $I_L$  which had been deduced previously, and which was applicable to the flame phase only, had to be extended and supplemented.

The study of the last-mentioned problem resulted in an expression for  $I_L$  which was based on the magnitude of the heat transfer by convection through the openings in the enclosed space. The rates of gas and air flow involved in this process were calculated in two steps, viz., first by determining the velocity distribution of gas flow in a vertical opening by which two masses of gas differing in density are separated from each other, and second, by satisfying the condition that the net exchange of gases between the enclosed space and its surroundings shall be equal to the difference between the quantity of gas produced and the quantity of air consumed in the process of combustion. After that, it was possible to determine  $I_L$  directly as the difference in heat content between the outgoing gases and the incoming air. It was found that  $I_L$  was approximately proportional to the temperature of the combustion gases and to the air flow factor  $A \cdot \sqrt{H}$ .

Since no physical basis is available which could enable the quantity of energy liberated per unit time during the process of fire development to be determined as a function of the time, a study of the literature was carried out with a view to an analysis of full-scale fire tests. For the tests where the external conditions were stated in a sufficiently precise manner, comparative calculations of temperature-time curves were made by means of a computer.

A tentatively chosen time graph of the quantity of energy liberated per unit time was used for this purpose. The time graph in question was varied until the agreement between the observed and calculated temperature-time curves became as close as possible. The only requirement to be fulfilled in this connection was that the total quantity of energy released during the whole process of fire development should be equal to that which was available in the fuel from the outset. When all those tests which were suited for this study had been examined, it was possible to systematize the results of the study in such a way that the time graph of the quantity of energy released per unit time during the process of fire development might generally be assumed to be known. This procedure was primarily justified by the consideration that an error, if any, could only be involved in the time graph of the quantity of energy liberated per unit time, since the total magnitude of this quantity is determined by the fire load, i.e. by the quantity of energy which is available from the beginning.

The computer programme which was used for the calculation of the temperature-time curves has a far-advanced general validity. One of the features of this programme is that it affords a possibility of taking into account various factors, viz., first, those thermal properties of the materials entering into the structures bounding the enclosed space which are dependent on the temperature, second, the variations in the dimensions of the openings during the process of fire development, third, the moisture content of the bounding structures, and fourth, the effects of heat-absorbing structures in the interior of the enclosed space. This programme can be used for enclosed spaces which are bounded by structures of up to three different types at the same time, and one of these structures may be built of up to three different materials.

Moreover, a modified programme has been prepared for enclosed spaces provided with plasterboard panel walls, which are assumed under certain definite conditions to disintegrate during the fire.

The time graphs of the quantity of energy liberated per unit time which had been obtained by means of the method outlined in the above were used to calculate the time graphs of the temperature of the combustion gases during the process of fire development. The latter time graphs were computed on the assumption of different values of the fire load and the opening factor for seven types of enclosed spaces which differ in respect of the bounding structures. The results of these calculations are represented in graphs as well as in tables.

In carrying out the comparative theoretical analyses, it was possible to discuss to a limited extent the effects produced on the temperature-time curve of the process of fire development by some quantities which do not directly enter into the equation of heat balance, with the result that their effects on this process must be determined in each individual case. In addition to the size and the shape of the openings, the factors which may be expected to be of importance in this connection comprise, among others, the porosity of the fuel and its distribution in the enclosed space, the moisture content of the fuel, and the magnitude of the fire load. For future research in this field, it may be urgently recommended to make a study of the effects produced by these and other parameters on the quantity of energy released per unit time in the process of combustion, and hence also on the temperature-time curve, which is dependent on this quantity.

# Acknowledgements

The Authors wish to express their heartfelt gratitude to Professor Ove Pettersson, Head of the Division of Structural Mechanics and Concrete Construction, Lund Institute of Technology, Lund, Sweden, who has provided the impetus to the investigation described in the present publication, and whose ever-active support has greatly contributed to the accomplishment of this task. Our thanks are furthermore due to Mr. Ulf Hårdner, Miss Yvonne Fransson, Miss Kerstin Svensson, and Miss Lena Öberg, for their valuable collaboration in editing the manuscript, and to Mr. Ilya Cyon, for the translation of the manuscript into English.

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Appendix

## **APPENDIX 1**

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.



Test series A. See chapter 5.



Percentages of the total bounding surface area: Concrete, 20 cm in thickness, 34.8 per cent. Lightweight concrete, 12.5 cm in thickness, 42.2 per cent. Concrete, 3 cm in thickness+lightweight concrete, 10 cm in thickness, 18.3 per cent. Window area 4.7 per cent. Opening factor 0.06 m<sup>1/2</sup> (t > 0.1 h). Duration of the fire 0.17 h. Fire load 15.1 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.





Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 34.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 42.2 per cent.

Concrete, 3 cm in thickness+lightweight concrete, 10 cm in thickness, 18.3 per cent.

Window area 4.7 per cent.

Opening factor 0.06  $m^{1/2}$  (t > 0.1 h).

Duration of the fire 0.21 h.

Fire load 19 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.



#### Test A3

Percentages of the total bounding surface area: Concrete, 20 cm in thickness, 38.6 per cent. Lightweight concrete, 12.5 cm in thickness, 60.0 per cent. Window area 1.4 per cent. Opening factor 0.0356 m<sup>1/2</sup> (t > 0.2 h). Duration of the fire 0.36 h. Fire load 19.6 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.



#### Test A4

Percentages of the total bounding surface area: Concrete, 20 cm in thickness, 38.3 per cent. Lightweight concrete, 12.5 cm in thickness, 59.4 per cent. Window area 2.3 per cent. Opening factor 0.0486  $m^{1/2}$  (t > 0.2 h). Duration of the fire 0.325 h. Fire load 22.4 Mcal  $\cdot m^{-2}$  of bounding surface area.



#### Test A5

Percentages of the total bounding surface area: Concrete, 20 cm in thickness, 38.3 per cent. Lightweight concrete, 12.5 cm in thickness, 59.4 per cent. Window area 2.3 per cent. Opening factor 0.0548  $m^{1/2}$ . Duration of the fire 0.24 h. Fire load 20 Mcal  $\cdot m^{-2}$  of bounding surface area.



#### Test A6

Percentages of the total bounding surface area: Concrete, 20 cm in thickness, 37.8 per cent. Lightweight concrete, 12.5 cm in thickness, 58.6 per cent. Window area 3.6 per cent. Opening factor 0.068 m<sup>1/2</sup>. Duration of the fire 0.23 h.

Fire load 23 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.



Living-room.

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 40.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 47.3 per cent.

Concrete, 3 cm in thickness+lightweight concrete 10 cm in thickness, 8.3 per cent.

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Window area 3.6 per cent.

Opening factor 0.04 m<sup>1/2</sup>.

Duration of the fire 0.32 h.

Fire load 32 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.



#### Test A7b

Bedroom.

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 40.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 47.3 per cent.

Concrete, 3 cm in thickness+lightweight concrete 10 cm in thickness, 8.3 per cent. Window area 3.6 per cent.

Opening factor  $0.04 \text{ m}^{1/2}$ .

Duration of the fire 0.32 h.

4

Fire load 32 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.

## APPENDIX 2

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series B. See chapter 5.





Opening factor 0.0467 m<sup>1 /2</sup>. Duration of the fire 0.48 h. Fire load 33.3 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.







Opening factor 0.0467 m<sup>1/2</sup>. Duration of the fire 1.18 h. Fire load 83.5 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.

## APPENDIX 3

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series C. See chapter 5.



#### Test Cl

Quantity of combustible material 270 kg. Rate of air supply by fan 1.0 m<sup>3</sup>/s. Opening factor 0.0234 m<sup>1/2</sup>. Hydraulic radius —.





Quantity of combustible material 675 kg. Rate of air supply by fan 2.0 m<sup>3</sup>/s. Opening factor 0.0601 m<sup>1/2</sup>. Hydraulic radius 1.0 cm.





Quantity of combustible material 675 kg. Rate of air supply by fan 1.0 m<sup>3</sup>/s. Opening factor 0.0434 m<sup>1/2</sup>. Hydraulic radius 1.0 cm.









Quantity of combustible material 675 kg. Rate of air supply by fan 1.5 m<sup>3</sup>/s. Opening factor  $0.0551 \text{ m}^{1/2}$ . Hydraulic radius 1.0 cm.







#### Test C7

Quantity of combustible material 675 kg. Rate of air supply by fan 1.0 m<sup>3</sup>/s. Opening factor 0.0367 m<sup>1/2</sup>. Hydraulic radius 1.7 cm.



Test C8 Quantity of combustible material 675 kg. Rate of air supply by fan 0.7 m<sup>3</sup>/s. Opening factor 0.015 m<sup>1/2</sup>. Hydraulic radius 1.7 cm.



## Test C9

Quantity of combustible material 675 kg. Rate of air supply by fan 1.0 m<sup>3</sup>/s. Opening factor 0.051 m<sup>1/2</sup>. Hydraulic radius 0.6 cm.









Quantity of combustible material 405 kg. Rate of air supply by fan 0.7 m<sup>3</sup>/s. Opening factor  $0.02 \text{ m}^{1/2}$ . Hydraulic radius 1.0 cm.



Test C12 Quantity of combustible material 405 kg. Rate of air supply by fan 0.7 m<sup>3</sup>/s. Opening factor 0.012 m<sup>1/2</sup>. Hydraulic radius 2.4 cm.





Quantity of combustible material 135 kg. Rate of air supply by fan  $0.7 \text{ m}^3$ /s. Opening factor  $0.005 \text{ m}^{1/2}$ . Hydraulic radius 1.0 cm.







#### Test C15

Quantity of combustible material 1350 kg. Rate of air supply by fan 2.0 m<sup>3</sup>/s. Opening factor 0.0601 m<sup>1/2</sup>. Hydraulic radius 1.4 cm.



Test C16

Quantity of combustible material 405 kg. Rate of air supply by fan 0.7 m<sup>3</sup>/s. Opening factor 0.06 m<sup>1/2</sup>. Hydraulic radius 0.4 cm.

## **APPENDIX 4**

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases. Test series D. See chapter 5.



#### Test DI

Percentages of the total bounding surface area: Lightweight concrete, 20 cm in thickness, 97.8 per cent. Window area 2.2 per cent. Opening factor 0.028  $m^{1/2}$ . Duration of the fire 0.5 h. Fire load 20.2 Mcal  $\cdot m^{-2}$  of bounding surface area.



#### Test D2

Percentages of the total bounding surface area: Lightweight concrete, 20 cm in thickness, 98.7 per cent. Window area 1.3 per cent. Opening factor 0.0147  $m^{1/2}$ . Duration of the fire 0.51 h. Fire load 11.2 Mcal  $\cdot m^{-2}$  of bounding surface area.



#### Test D3

Percentages of the total bounding surface area: Lightweight concrete, 20 cm in thickness, 99.15 per cent. Window area 0.85 per cent. Opening factor 0.0092 m<sup>1/2</sup>. Duration of the fire 0.47 h. Fire load 6.5 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.



#### Test D4

Percentages of the total bounding surface area: Lightweight concrete, 94.7 per cent. Window area 5.3 per cent. Opening factor 0.075  $m^{1/2}$ . Duration of the fire 0.58 h. Fire load 65 Mcal  $\cdot m^{-2}$  of bounding surface area.

# APPENDIX 5

Calculated time graphs of temperature of combustion gases for seven types of enclosed spaces differing in opening factor and in bounding structures. See chapter 7.



Type A enclosed space. Opening factor 0.01  $m^{1/2}$ .





Opening factor 0.02 m<sup>1/2</sup>.



















Type A enclosed space. Opening factor  $0.12 \text{ m}^{1/2}$ .
















Type B enclosed space. Opening factor  $0.06 \text{ m}^{1/2}$ .

















Type C enclosed space. Opening factor  $0.02 \text{ m}^{1/2}$ .

















Type D enclosed space. Opening factor  $0.02 \text{ m}^{1/2}$ .

























Type D enclosed space. Opening factor 0.12  $m^{1/2}$ .





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Type E enclosed space. Opening factor  $0.02 \text{ m}^{1/2}$ .

















#### F1

Type F enclosed space. Resultant emissivity 0.1. Opening factor 0.01  $m^{1/2}$ .





Type F enclosed space. Resultant emissivity 0.1. Opening factor  $0.02 \text{ m}^{1/2}$ .



FЗ

Type F enclosed space. Resultant emissivity 0.1. Opening factor 0.04  $m^{1/2}$ .





Type F enclosed space. Resultant emissivity 0.1. Opening factor 0.08  $m^{1/2}$ .



F5

Type F. enclosed space. Resultant emissivity 0.1. Opening factor 0.12  $m^{1/2}$ .







### F7

Type F enclosed space. Resultant emissivity 0.35. Opening factor  $0.02 \text{ m}^{1/2}$ .



#### F8

Type F enclosed space. Resultant emissivity 0.35. Opening factor 0.04  $m^{1/2}$ .



### F9

Type F enclosed space. Resultant emissivity 0.35. Opening factor 0.08  $m^{1/2}$ .



Type F enclosed space. Resultant emissivity 0.35. Opening factor  $0.12 \text{ m}^{1/2}$ .



FII

Type F enclosed space. Resultant emissivity 0.60. Opening factor 0.01  $m^{1/2}$ .





Type F enclosed space. Resultant emissivity 0.60. Opening factor  $0.02 \text{ m}^{1/2}$ .



#### F13

Type F enclosed space. Resultant emissivity 0.60. Opening factor  $0.04 \text{ m}^{1/2}$ .





Type F enclosed space. Resultant emissivity 0.60. Opening factor 0.08  $m^{1/2}$ .



F15

Type F enclosed space. Resultant emissivity 0.60. Opening factor  $0.12 \text{ m}^{1/2}$ .





Type G enclosed space. Opening factor 0.01  $m^{1/2}$ . O=Plasterboard panel falls down.





O=Plasterboard panel falls down.



Type G enclosed space. Opening factor 0.04  $m^{1/2}$ . O=Plasterboard panel falls down.

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Type G enclosed space. Opening factor 0.08  $m^{1/2}$ . O=Plasterboard panel falls down.





### APPENDIX 6

Calculated time graphs of temperature of combustion gases represented in tabular form for seven types of enclosed spaces differing in opening factor and in bounding structures. See chapter 7.

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# <u>A1</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_t = 0.01 \text{ m}^{1/2}$ 

т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q.	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	Т	e m	p	е	r a	. t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.20 3.40 3.60 3.20 3.40 3.60 3.80 4.20 4.60 4.80 5.00 5.20 5.40 5.60	272 395 228 196 150 98 97 94 91 88 82 79 75 69 61 54 46 45 43 42 41 40 39 38 37 36 35 34 33 32 32	272 3990 368 317 248 248 248 248 241 0 998 377 665 31 976 432 40 998 3776 3555 3 333 35554	272 3901 4085 4095 402 222 1878 6429 62 11 11 11 11 11 10 885 1964 0631 965 555 55 50 98 76 45 4322 1 165432 20 98 76 45 4322 1 1655 55 55 55 55 55 55 55 55 55 55 55 55	272 3990 4010 4121 437 4412 3355 422 218 299 1900 1800 150 137 20 83 80 77 52 08 66 53 20 98 7 70 86 65 320 58 7	272 3260 405 5 434 2 444 465 1 434 2 444 465 1 439 2 33 3 3 15 9 282 2 43 2 20 7 5 2 185 2 104 0 9 3 0 7 5 2 8 8 8 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 5 2 0 8 7 7 100 6 3 0 7 7 100	2728 3606 441 444 446 4768 866 441 444 466 444 44 466 37 33 33 32 22 22 21 188 40 59 40 59 40 59 40 63 99 22 22 21 188 40 59 40 59 40 63 99 29 20 63 20 10 63 20 10 63 20 10 63 20 10 10 50 10 10 50 10 10 50 10 50 10 50 10 50 50 50 50 50 50 50 50 50 50 50 50 50	2728065555429306801111442555555544823171663198765443210 380655554293068011114425555555544823171663198765443210	2728 3606 4055 554 4493 7688 011111 555 555 555 555 555 555 555 555

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor  $A\cdot\sqrt{H}/A_t$  = 0.02  $m^{1/2}$ 

т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0
Time	T	е п	р	е	r a	t	u	r e
h			2					
$\begin{array}{c} n \\ 0.05 \\ 0.10 \\ 0.15 \\ 0.20 \\ 0.25 \\ 0.30 \\ 0.35 \\ 0.40 \\ 0.45 \\ 0.50 \\ 0.60 \\ 0.65 \\ 0.70 \\ 0.80 \\ 0.90 \\ 1.00 \\ 1.10 \\ 1.20 \\ 1.30 \\ 1.40 \\ 1.50 \\ 1.60 \\ 1.70 \\ 1.80 \\ 1.90 \\ 2.00 \\ 2.20 \\ 2.40 \\ 2.60 \\ 2.80 \\ 3.00 \\ 3.20 \\ 3.40 \\ 3.60 \\ 3.80 \\ 4.00 \\ 4.20 \\ 4.40 \\ 4.60 \\ 4.80 \\ 5.00 \\ 5.20 \\ 5.40 \\ 5.60 \\ 5.80 \\ 6.00 \end{array}$	396 568 322 277 210 134 131 126 121 107 102 98 88 77 66 54 51 43 42 40 38 35 34 33 231 30 30	396 5559 330 48 222 11 11 120 88 77 77 76 53 96 31 97 64 432 10 98 76 63 55 51 97 64 432 10 98 76 63 33 62 197 64 432 10 98 76 63 33 33 66 55 54 55 55 54 56 55 55 55 55 55 55 55 55 55 55 55 55	396862591276774902454317398638306318642198765544444444444444444444444444444444444	3968 5572 566 663 554 419 995 595 2022 214 994 8518 773 718 664 615 586 615 586 615 548 419 559 50 222 214 948 5336 94 95 517 777 718 664 615 586 557 777 718 664 615 586 557 777 718 666 615 586 557 777 718 666 615 586 557 7777 718 666 615 557 587 5777 718 666 615 577 718 718 718 718 718 718 718 718 718 7	$\begin{array}{c} 397\\ 4511\\ 581\\ 578\\ 593\\ 612\\ 666\\ 642\\ 452\\ 517\\ 329\\ 402\\ 316\\ 504\\ 432\\ 222\\ 206\\ 718\\ 128\\ 128\\ 128\\ 128\\ 128\\ 128\\ 128\\ 1$	$\begin{array}{c} 397\\ 551\\ 573\\ 593\\ 579\\ 620\\ 22\\ 22\\ 22\\ 230\\ 231\\ 22\\ 230\\ 231\\ 212\\ 103\\ 231\\ 22\\ 230\\ 231\\ 212\\ 230\\ 231\\ 212\\ 230\\ 231\\ 212\\ 230\\ 231\\ 212\\ 230\\ 231\\ 212\\ 230\\ 231\\ 212\\ 230\\ 231\\ 212\\ 230\\ 231\\ 212\\ 230\\ 231\\ 231\\ 212\\ 230\\ 231\\ 231\\ 231\\ 231\\ 231\\ 231\\ 231\\ 231$	$\begin{array}{c} 3967\\ 55818\\ 5793\\ 6666\\ 6666\\ 6671\\ 777\\ 777\\ 777\\ 776\\ 6666\\ 6698\\ 552\\ 515\\ 008\\ 8936\\ 1292\\ 261\\ 594\\ 9348\\ 8936\\ 158\\ 2219\\ 1849\\ 138\\ 30936\\ 159\\ 1849\\ 138\\ 138\\ 138\\ 1849\\ 158\\ 138\\ 138\\ 1849\\ 158\\ 138\\ 1849\\ 158\\ 138\\ 1849\\ 158\\ 138\\ 1849\\ 158\\ 138\\ 1849\\ 158\\ 186\\ 186\\ 186\\ 186\\ 186\\ 186\\ 186\\ 18$	398 4511 587 595 66222221 8667 129268 9007 878 84033922221 866774 3860743 277 76662218 3367774 3807431222222222222222222222222222222222222

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# <u>A2</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_t = 0.04 m^{1/2}$ 

	v							
т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time	т	e m	P	e	r a	t	u	r e
h								
0.05	504	504	504	504	504	504	504	504
0.10	745	745	745 75-	745	621 681	621 681	621 681	621 681
0.15 0.20	422 360	747 696	747 767	747 767	661 777	601 777	601 777	777
0.25	268	587	784	784	776	776	776	776
0.30	164	472	734	799	793	793	793	793
0.35	162	437	665	814	808	808	808	808
0.40	155 148	389	593	828	822	822	822	822
0.45	148 142	337 281	513 481	841 779	836 848	836 848	836 848	836 848
0.50 0.60	142		397	682	874	874	874	874
0.65	120	259 246	352	626	882	882	882	882
0.70	114	232	307	565	839	894	894	894
0.80	100	204	285	527	785	912	912	912
0.90	86	178	260 235	483 437	720 645	862 827	928 942	928 942
1.00 1.10	71 54	149 118	235 208	437 388	589	787	942	955
1.20	51	85	183	337	555	740	967	967
1.30	49	82	156	316	518	688	942	977
1.40	46	77	128	296	480	632	931	987
1.50 1.60	45 43	74 70	98 94	276 255	441 400	602 571	919 895	996 1004
1.70	43 41	68	94 89	235	358	540	870	981
1.80	40	65	85	214	343	507	843	973
1.90	39	62	82	194	328	474	813	963
2,00	38	60	79	174	313	440	781	953
2.20 2.40	36 35	. 56 52	73 69	131 104	288 263	369 339	718 655	923 890
2.40	33	50	64	96	238	311	587	853
2.80	32	47	61	90	214	286	516	813
3.00	31	45	57	84	190	261	442	769
3.20		43	55	80 76	166 141	236 211	388 362	727 682
3.40 3.60		42 40	52 50	72.	141	187	338	635
3.80		39	48	69	99	163	316	587
4.00		38	47	66	94	137	296	537
4.20		37	45	64	89	110	277	485
4.40		36	44 20	61	85	104	258 240	431 388
4.60 4.80		35 34	42 41	59 57	81 78	99 94	240 221	300 360
5.00		33	40	55	75	90	204	332
5,20			39	53	73	87	186	305
5.40		di se	38	52	70	83	168	280
5.60			37 37	50 49	68 66	81 78	150 131	255 229
5.80 6.00			37 36	49 48	66 64	10 75	131	203
0.00			00	40	<b>V</b> 7			~~J

A3

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor  $A\cdot\sqrt{H}/A_{\rm t}$  = 0.06 m<sup>1/2</sup>

		•							
Т		0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q		9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time		т	e m	р	е	r a	t	u	r e
h				-					
0.05		575	575	575	575	575	575	575	575
0.10		858	858	858	858	704	704	704	704
0.15		493	861	861	861	784	784	784	784
0.20		404	802	879	879	882	882	882	882
0.25		296	679	898	898	889	889	889	890
0.30		175	538	838	914	908	908	908	908
0.35		174	490	761	928	923	923	923	923
0.40		166	430	669	942	936	936	937	937
0.45	20	159	369	572	954	949	949	949	949
0.50		151	303	532 3	877	961	961	961	961
0.60		136	277	433	762	982	982	982	982
0.65		128	262	402	694	992	992	992	992
0.70		120	247	326	620	939	1001	1001	1001
0.80		104	215	300	574	872	1018	1018	1018
0.90		89	185	272	520	795	954	1032	1032
1.00		71	152	243	466	705	909	1044	1044
1.10		51	116	213	409	637	858	1054	1054
1.20		48	80	184	343	593	803	1064	1064
1.30		45	76	155	327	550	742	1029	1072
1.40		43	72	123	303	505	675	1013	1080
1.50		41	68	89	281	460	640	996	1087
1.60		40	65	86	259	413	603	966	1093
1.70		38	62	81	236	364	567	935	1062
1.80		37	59	78	213	348	529	902	1049
1.90		36	56	74	191	332	491	866	1036
2.00		35	54	71	169	317	452	830	1022
2.20		33	50	66	121	289	371	756	984
2.40		32	47	61	93	263	340	683	943
2.60		31	44	57	85 -	236	310	605	900
2,80		30	42	54	79	210	283	524	854
3.00		29	40	51	74	185	257	440	805
3.20			39	48	70	159	230	381	756
3.40			37	46	67	131	204	355	705
3.60			36	կկ	63	103	178	331	652
3.80			35	42	60	86	152	308	597
4.00			34	41	57	81	123	288	541
4.20			33	40	55	77	95	269	483
4.40			32	38	53	73	90	249	423
4.60			31	37	51	70	85	230	377
4.80			31	36	49	67	81	211	348
5.00			30	35	47	64	77	193	319
5.20				35	46	62	74	174	292
5.40				34	կկ	59	71	155	265
5.60				33	43	57	69	135	238
5.80				32	42	55	66	114	210
6.00				32	41	54	6h	94	185

<u>A5</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.08 m<sup>1/2</sup>

	v							
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
Q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time	T	e m	ą	e	r a	. t	u	r e
h								
0.05	622	622	622	622	622	622	622	622
0.10	935	935	935	935	766	766	766	767
0.15	532	937	937	937	853	853	853	853
0.20	432	869	955	955	959	959	959	959
0.25	314	734	973	973	965	965	965	965
0.30	181	575	903	987	981	981	981	982
0.35	180	521	818	1001	995	995	995	995
0.40	171	454	720	1013	1008	1008	1008	1008
0.45	163	386	611	1024	1020	1020	1020	1020
0.50	155	314	561	937	1031	1031	1031	1031
0.60 0.65	139	285	454 396	807	1050	1050	1050	1050 1058
0.70	131 122	269 253	396	732 651	1058 996	1058 1066	1058 1066	1056
0.80	106	219	306	598	990	1081	1000	1081
0.90	89	186	275	539	833	1005	1092	1092
1.00	70	151	245	479	735	953	1102	1102
1.10	47	113	214	417	659	897	1111	1111
1.20	դդ	73	183	352	612	836	1119	1119
1.30	42	70	151	328	564	769	1077	1126
1.40	40	66	117	304	516	695	1058	1132
1.50	39	62	81	281	466	657	1038	1138
1.60	37 36	59	78	257	415	010	1004	1143
1.70	36	56	74	234	363	577	969	1105
1.80 1.90	35 34	53 51	71 68	210 187	347 331	537 496	932 893	1090 1074
2.00	33	49	65	163	316	490 454	853	1058
2.20	31	46	59	112	287	368	774	1016
2.40	30	43	55	83	260	337	695	971
2.60	29	40	51	77	233	307	611	923
2.80	28	38	48	71	206	279	52 <sup>1</sup> 4	873
3.00	27	. 37	45	67	180	252	434	820
3.20		35	43	62	153	224	373	769
3.40		34	41	59	124	198	347	714
3.60		33	40	55	94	171	323	658
3.80 4.00		32 31	38	53	77	143	302 282	599 539
4.00		51	37 36	50 48	72 68	113 84	262	539 478
4.20			35	40	65	79	242	415
4.60			34	45	62	75	222	368
4.80			33	43	59	71	203	338
5.00			32	42	56	68	185	309
5.20				40	54	65	166	282
5.40				39	52	62	146	254
5.60				38	50	60	125	227
5.80				37	48	57	104	199
6.00				36	47	55	82	172

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_t = 0.12 \text{ m}^{1/2}$ 

	ΰ.							
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	5.0
đ	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time	T	e m	р	e	r s	ı t	u	r e
h								
0.05	670	670	670	670	670	670	670	670
0.10	1027	1027	1027	1027	847	847	847	847
0.15	581	1033	1033	1033	933	933	933	933
0.20	465	951	1049	1049	1051	1051	1051	1051
0.25	333	799	1063	1063	1057	1057	1057	1057
0.30	186	620	981	1076	1071	1071	1071	1071
0.35	185	556	882	1088	1083	1083	1083	1083
0.40	176	480	774	1098	1094	1094	1094	1094
0.45	168	404	650	1107	1103	1103	1103	1103
0.50	159	324	593	1004	1112	1112	1112	1112
0.60	142	292	472	856	1127	1127	1127	1127
0.65	133	275	407	774	1133	1133	1133	1133
0.70	124	257	341	681	1060	1139	1139	1139
0.80	106	221	309	622	971	1150	1150	1150
0.90	88	186	276	556	873	1062	1159	1159
1.00	67	149	244	490	765	1001	1166	<b>116</b> 6
1.10	41	107	211	422	680	937	1173	1173
1.20	39	63	178	351	628	868	1178	1178
1.30	37	60	145	327	575	794	1128	1183
1.40	36	56	108	301	523	713	1106	1188
1.50	34	<i>,</i> 53	69	277	469	672	1082	1192
1.60	33	50	67	253	414	629	1043	1195
1.70	32	47	63 60	228	358	585 542	1003 962	1151
1.80	31	45 43		203 180	343	542 497	962 919	1133 1114
1.90 2.00	30 29	43 42	57 54	155	327 311	491	919 874	1096
2.00	29	39	54 49	100	283	492 361	789	1090
2.40		37	49	70	255	330	704	998
2.60		35	43	64	227	300	613	946
2.80		33	40	59	200	272	519	891
3.00		32	38	54	173	245	423	834
3.20		24	37	51	144	216	361	780
3.40			35	48	113	190	336	721
3.60			34	46	83	162	313	660
3.80			33	43	64	132	293	598
4.00			32	42	59	102	273	534
4.20				40	55	69	253	469
4.40				39	52	65	233	403
4.60				37	50	61	213	355
4.80				36	48	58	194	325
5.00				35	46	55	175	297
5.20				34	44	52	155	269
5.40				33	42	50	134	241
5.60				33	41	48	112	213
5.80				32	40	46	91	186
6.00				31	38	45	67	158

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# <u>A6</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.01 \text{ m}^{1/2}$ 

	U							
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time	т	e m	р	e	r a	t	u	r e
								r e 2376 2376 33520 3772 8804 400 100 398 4454 4488 824 43870 97727 1146665224 1852207 2271146665224 1852207
6.00		32	38	52	70	84	127	195

<u>B1</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor  $A \cdot \sqrt{E}/A_t = 0.02 \text{ m}^{1/2}$ 

T	0.1	0.2	0.3	0,5	0.75	1.0	1.5	2.0	
đ	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0	
Time h	Т	е п	р	e	r a	t	u		e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.60 5.80 6.00	353 498 288 291 110 105 60 7 65 8 8 54 45 42 40 38 7 65 83 33 33 21	353 497 592 203 201 11111119777666666575519766443210988877655544 4432109888776553344	3487232232222222222222222222222222222222	3548722600945339037998852852962347277777686653109776 111111098887742086653109776	354 411 511 500 5238 555 5459 555 5459 555 5459 555 5459 555 555	354 4511 506 528 555 560 575 526 637 90 379 308 579 212 58 100 109 126 134 1124 118 100 59 59 99 99 90	354 451 5528 555555555555555556666666666665555442 5528 55555555555555566666666666555547267 1110 11111111111111111111111111111111	$\begin{array}{c} 356\\ 414\\ 5516\\ 5528\\ 555\\ 564\\ 346\\ 677\\ 128\\ 716\\ 638\\ 639\\ 260\\ 338\\ 319\\ 260\\ 241\\ 221\\ 241\\ 225\\ 392\\ 639\\ 266\\ 557\\ 559\\ 466\\ 392\\ 260\\ 338\\ 319\\ 260\\ 241\\ 221\\ 221\\ 221\\ 221\\ 221\\ 221\\ 221$	

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## <u>B2</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A.VE/A\_t = 0.04  $\rm m^{1/2}$ 

90.0 u	120.0
u	
	r e
4578 606 6938 725 7469 802 837 806 8377 806 8377 806 8377 806 8377 806 8377 806 8377 806 8377 8079 8079 8079 8078 8078 8077 8078 8079 8079	457 558 6938 7252 749 769 8062 8355 8937 919 9312 9522 9424 9257 7031 6161 51724 3856 3799 9292 9425 9425 9316 51743 3569 3040 2551 206
	558 606 693 7249 7769 8022 831 8556 8907 9197 888 8563 8768 8563 87748 8563 87748 8563 87748 8563 87748 8563 87748 8566 9197 9273 800 2733 2037 2755 2037 2037 2037 2037 2037 2037 2037 2037

<u>B3</u>

	•				•			
Т	0.1	0.2	0.3	0.5	0.75	1.O	1.5	2.0
đ	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	р	e	r a	t	u	r
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.40 2.40 2.40 2.40 2.40 2.60 2.80 3.20 3.40 3.20 3.40 3.60 3.80 4.20 4.40 4.50 5.00	528 761 440 378 280 159 152 145 130 125 145 130 225 145 130 25 49 47 44 39 38 37 36 35 42 30 30	528 7666 0 5 6 6 2 9 7 5 1 8 6 5 2 9 7 5 5 1 8 6 4 4 2 1 9 8 7 6 5 3 4 4 3 3 2 2 2 2 2 2 1 1 1 1 7 7 1 8 5 2 9 7 5 5 1 8 6 4 4 2 1 9 8 7 6 5 3 4 3 3 2 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3	5216608702716477503570221962877763963197543109992	528 766 798 838 875 826 855 500 959 754 333 222 2218 12 98 827 7307 41 97 5531 5531	528 632 703 806 805 850 850 935 940 575 486 401 552 486 401 320 882 258 200 851 887 52 95 136 90 51 887 72 97 697	5282 7036599 889999999988 877765306559597559979988 87776530661551581444 333599483999999888777655551581444 333599483994799506823	528 632 703 806 805 935 949 979 909 1021 999 960 975 960 878 807 595 438 807 767 2515 437 207 778 809 1021 999 960 878 887 667 2515 437 287 287 212 4 212 4 212 4 212 4 212 212	$\begin{array}{c} 528\\ 632\\ 703\\ 8065\\ 8250\\ 9359\\ 949\\ 9755\\ 9091\\ 1032\\ 10519\\ 10325\\ 1067\\ 7522\\ 9963\\ 886\\ 7492\\ 641\\ 884\\ 871\\ 963\\ 887\\ 7492\\ 641\\ 3320\\ 320\\ 8836\\ 8122\\ 1051\\ 10252\\ 9963\\ 988\\ 7694\\ 1037\\ 5348\\ 320\\ 1037\\ 10025\\ 1$

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76

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267

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Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space Opening Factor  $A\cdot\sqrt{H}/A_t$  = 0.06  $m^{1/2}$ 

5.00 5.20 5.40

5.60

5.80 6.00

<u>B4</u>

e

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A  $\sqrt{H}/A_{\rm t}$  = 0.08  $\rm m^{1/2}$ 

-	U U							
т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time	т	e m	р	e	r a	t	u	r e
h								,
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.40 2.40 2.60 2.40 3.60 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.60 5.80 6.00	576 836 480 409 301 177 175 167 159 151 135 127 103 87 459 475 420 398 36 353 33	576 836 845 547 94 376 837 437 2247 183 932 247 183 932 247 183 932 247 183 932 247 183 932 247 183 932 247 17 660 555 51 853 1 938 653 4 33 34 34	5884177424246001111108317307284196532098765543333222111188777662841965320987655433332221111887776628419653209876554333333333333	5768887999999977165826886439974188645706666555319765444444444444444444444444444444444444	284 258	$\begin{array}{c} 576\\ 693\\ 773\\ 886\\ 910\\ 935\\ 973\\ 9908\\ 1035\\ 9978\\ 1035\\ 9918\\ 10352\\ 9918\\ 10352\\ 9918\\ 10352\\ 9918\\ 10352\\ 9931\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9933\\ 10352\\ 9333\\ 10352\\ 9333\\ 10352\\ 9333\\ 10352\\ 9333\\ 10352$	$\begin{array}{c} 576\\ 694\\ 773\\ 871\\ 886\\ 910\\ 935\\ 973\\ 9018\\ 1036\\ 1056\\ 1077\\ 1087\\ 10956\\ 1095\\ 1005\\ $	576 694 773 871 886 910 935 973 9018 1026 1057 1097 11124 1087 10975 1012 1026 10421 1058 10421 1058 5376 5935 415 9311 284 72303 17
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<u>B5</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.12 \text{ m}^{1/2}$ 

	v							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time	т	e m	q	е	r e	ı t	u	r e
h								
h 0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.20 4.40 5.00 5.20 5.40	T 624 940 548 1447 125 183 174 166 157 140 131 122 105 87 43 37 34 33 31 31	e m 624 940 952 892 762 600 541 397 289 275 2195 149 108 66 53 548 46 44 42 37 653 322 31 30 31 30	p 40206253625397530960207396420876543221 11177666073964208765433221 1117766607396420876543221 11177666073964420876543221 11177666073964420876543221	6240 9950 1002518 10083 9840 100518 87671283 19940 10752 2221 107683 100518 87671283 1075208 4644210 38775208 46575208 46575208 46775208 47775208 47775208 47775208 47775208 47775208 47775208 47775208 47775200 47775200 47775200 47775200 47775200 47775200 47775200 47775200 47775200 47775200 47775200 47775200 47775200 4777520000000000000000000000000000000000	6247 867 980 9959 1040 10666 1010 1079 1113 9559 2079 8592 209 17466 1113 1045 1040 1040 1040 1040 1040 1040 1040	$\begin{array}{c} 624\\ 777\\ 860\\ 999\\ 1019\\ 1060\\ 1079\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1120\\ 1046\\ 787\\ 702\\ 105\\ 105\\ 105\\ 105\\ 105\\ 105\\ 105\\ 105$	u 624 777 867 980 995 1019 1060 1076 1089 1106 1113 1120 1133 1143 1159 1165 1093 1070 1031 9951 864 780 696 8517 3628 3314 274 254 215 196 158 138	r e 624 777 867 980 995 1019 1040 1060 1076 1089 1106 1113 1120 1133 1143 1152 1159 1165 1171 1176 1181 1185 1241 1205 1086 1039 938 884 827 773 5656 594 237 299 244
4.80 5.00 5.20			32 32	41 40 39 38 37	55 52 50 49 47	67 63 60 58 55	215 196 178 158 138	357 327 299 272 244
5.60 5.80 6.00				36 35 34	45 44 43	53 51 49	116 96 73	216 189 162

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### <u>B6</u>
<u>C1</u>

Opening Factor	$A \cdot V H / A_t = 0$	0.01 m <sup>-/-</sup>						
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	Т	е т	p	e	r a	t	u	r e
Time	T 502 778 430 347 156 154 148 142 137 125 113 101 88 75 60 57 52 9 47 46 42 139 37 35 34 32 32 32 34 32 32 32 32 32 32 32 32 32 32	502 77694 413 319 22171 123 905 17708 506 506 5207 45310 310 408 506 5207 45310 408 307 507 507 507 507 507 507 507 507 507 5		e 502 778 769 779 780 780 780 780 780 780 780 780 780 780	r 8 502 630 696 7777 782 788 795 801 808 825 782 788 795 801 808 825 782 788 668 604 557 766 802 782 782 782 788 801 808 825 782 782 782 788 801 808 825 782 782 782 788 801 808 825 782 782 782 782 782 788 801 808 825 782 782 782 782 782 788 801 808 825 727 803 802 825 727 803 802 825 727 803 802 825 727 803 802 825 727 803 825 727 803 825 727 803 825 727 803 825 727 82 827 82 827 827 827 82 827 827 82	t 502 630 697 778 808 825 808 825 808 825 836 825 836 825 836 825 828 837 758 837 758 838 825 836 825 836 825 836 826 837 758 838 826 837 758 838 826 837 758 838 828 836 758 837 758 838 828 836 758 837 758 838 828 839 758 837 758 838 828 839 758 837 758 838 828 839 758 839 758 838 839 758 8395 758 839 758 758 757 758 757 758 757 758 757 758 757 757	u 563067778285518888888888888888888888888888888	r e
4.80 5.00 5.20 5.40 5.60		35 34 33	41 40 39 38 37	57 55 53 51 50	81 77 74 71 69	99 94 90 86 82	225 208 191 173 156	360 333 307 283 259
5.80 6.00			36 35	48 47	66 64	79 76	137 118	234 209

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor  $A \cdot \sqrt{H} / A_{+} = 0.01 \text{ m}^{1/2}$ 

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor  $\Lambda\cdot\sqrt{H}/A_{\rm t}$  = 0.02  $m^{1/2}$ 

	0								
т	0.1		0.3	0.5	0.75	1.0	1.5	2.0	
đ	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0	
Time	Т	e n	n p	e	r e	a t	u	r e	!
h									
0.05	634	634	634	634	634	634	634	634	
0.10 0.15	992 538	992 977	992 977	992 977	805 884	805 884	805 884	805 884	
0.20	425	889	979	979	986	986	986	986	
0.25	303	734	982	982	981	981	981	981	
0.30	175	563 501	901	987	985 989	985	985	985	
0.35 0.40	174 167	504 439	805 701	991 997	909 995	989 995	989 995	989 995	
0.45	160	373	590	999	997	997	997	997	
0.50	153	304	542	906	999	999	999	999	
0.60	138	279	441	774	1003	1003 1010	1003 1010	1003 1010	
0.65 0.70	130 122	265 250	387 331	703 627	996 946	1010	1010	1016	
0.80	107	219	304	581	880	1032	1010	1032	
0.90	91	188	276	529	803	967	1045	1045	
1.00	73	155	248	474	715	922	1057	1057	
1.10 1.20	53 50	120 82	218 189	417 356	648 605	872 817	1069 1079	1069 1079	
1.30	47	78	159	333	562	756	1044	1088	
1.40	45	73	126	309	517	689	1028	1095	
1.50 1.60	42 41	69 66	93 88	286 263	470 422	654 616	1011 980	1102	
1.70	41 39	62	83	203 240	422 373	579	960 949	1108 1076	
1.80	37	59	78	217	356	540	915	1062	
1.90	36	56	74	194	339	501	879	1048	
2.00 2.20	35 33	53 49	71 64	172 123	322 293	462 379	841 767	1033 993	
2.40	31	49	59	93	266	345	693	995 951	
2.60	30	42	54	84	239	314	614	907	
2,80	29 28	40	50	77	212	286	532	860	
3.00 3.20	28	38 36	47 44	72 67	186 160	259 232	447 386	811 762	
3.40		34	42	62	131	205	357	711	
3.60		33	40	58	105	179	332	658	
3.80		32	38	55	85	151	309	602	
4.00 4.20		31	37 36	52 49	79 74	122 93	288 268	546 488	
4.40			34	47	70	87	249	428	
4.60			33	45	66	81	229	380	
4.80 5.00			32 32	43 bo	62	77	209	349	
5.20			22	42 40	59 56	73 69	191 172	318 290	
5.40				39	54	66	152	263	
5.60				38	51	63	132	236	
5.80 6.00				37	49 48	60 57	111	208	
0.00				36	40	57	90	182	

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<u>C2</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor A· $\sqrt{H}/A_t = 0.04 m^{1/2}$ 

-	U U							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time	T	e :	m p	e	r a	, t	u	r e
Ь								
0.05	726	726	726	726	726	726	726	726
0.10 0.15	1126 618	1126 1110	1126 1110	1126 1110	916 1002	916 1002	916 1002	916 1002
0.20	478	1005	1111	1111	1122	1122	1122	1122
0.25	334	829	1112	1115	1111	1111	1111	1111
0.30 0.35	183 182	631 559	1013 902	1114 1115	1113 1115	1113 1115	1113 1115	1113 1115
0.40	174	480	902 784	1116	1118	1118	1118	1118
0.45	166	401	651	1115	1117	1117	1117	1117
0.50	158	320	593	1005	1115	1115	1115	1115
0.60 0.65	141 132	290 273	469 404	852 770	1126 1131	1126 1132	1126 1132	1126 1132
0.70	123	256	338	679	1058	1137	1137	1137
0.80	106	221	308	621	970	1148	1148	1148
0.90	88	186	276	557	874	1061	1158	1158
1.00	67 41	149 107	244 211	491 424	767 684	1002 938	1166 1172	1166 1172
1.20	39	63	178	353	631	870	1178	1178
1.30	37	60	144	328	579	797	1128	1182
1.40 1.50	35 34	55 52	107 68	302 278	526 473	717 676	1105 1081	1184 1187
1.60	34 32	52 48	66	254	413	631	1001	1189
1.70	31	46	62	229	361	589	1003	1146
1.80	30	43	58	204	345	545	962	1128
1.90 2.00	29 28	41 39	54 51	180 155	328 312	500 456	919 875	1110 1091
2.20	20	36	46	100	283	363	790	1045
2.40		34	42	69	255	331	703	996
2.60 2.80		32 30	39 37	61 55	227	300 272	614 521	945 890
3.00		29	35	22 51	199 172	212	425	834
3.20			33	47	143	216	362	780
3.40			31	<u>4</u> 4	112	189	336	722
3.60 3.80			30 29	41 39	81 61	161 131	312 291	661 599
4.00			28	37	56	100	271	535
4.20				36	51	66	251	471
4.40 4.60				34 33	48 45	62 57	231 211	404 355
4.80				32	49	53	193	324
5.00				31	41	50	173	295
5.20					39	47	153	267
5.40 5.60					38 36	45 43	132 110	239 210
5.80					35	43 41	89	184
6.00					34	39	64	155

<u>C3</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor A+ $\sqrt{H}/A_t$  = 0.06 m<sup>1/2</sup>

		v							
Т		0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ		9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h		Т	e m	р	e	r a	t	u	r e
0.05 0.10		766 1183	766 1183	766 1183	766 1183	766 962	766 962	766 962	766 962
0.15		650	1165	1165	1165	1052	1052	1052	1052
0.20 0.25		500 346	1953 867	1165 1166	1165 1166	1178 1166	1178 1166	1178 1166	1178 1166
0.30		185	659	1060	1168	1167	1167	1167	1167
0.35 0.40		185 176	581 498	941 815	1167 1166	1167 1167	1167 1167	1167 1167	1167 1167
0.45		168	490 411	674	1166	1166	.1166	1166	1166
0.50 0.60		159 142	323 292	611 477	1048 884	1166	1166	1164	1166
0.65		142	292 275	477 408	004 795	1177 1181	$1177 \\ 1181$	1177 1181	1177 1181
0.70 0.80		123	257	336	696	1100	1186	1186	1186
0.90		105 85	220 183	307 273	634 564	1003 898	1194 1096	1194 1201	1194 1201
1.00 1.10		64	144	240	494 422	782	1031	1207	1207
1.20		35 33	101 53	206 173	422 348	693 637	961 887	1211 1214	1211 1214
1.30 1.40		32	49	137	323	582	809	1158	1216
1.40		30 29	45 42	98 56	297 273	526 469	723 681	1132 1107	1218 1220
1,60		28	40	54	248	412	635	1064	1222
1.70 1.80		27 26	38 36	50 47	222 197	- 353 338	589 543	1022 978	1174 1153
1.90		26	34	44	173	321	497	933	1134
2.00 2.20		25	33	42 38	147 90	306 278	450 354	885 796	1114 1064
2.40			4	35	55	250	323	707	1012
2.60 2.80	· .			33 31	49 կկ	221 194	294 266	612 514	957 900
3.00				29	41	166	238	414	841
3.20 3.40					38 36	136 105	209 182	351 326	784 723
3.60					34	71	153	304	660
3.80 4.00					32 31	49 44	122 91	284 264	595 529
4.20					10	41	54	244	462
4.40 4.60						39 37	49 45	224 205	394 344
4.80						35	42	186	314
5.00 5.20						34	40 38	166 146	286 258
5.40							36	124	229
5.60 5.80							35 34	102 79	201 174
6.00							33	51	144

### <u>C4</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor  $A^{*}\sqrt{H}/A_{\rm t}$  = 0.08  $\rm m^{1/2}$ 

	Ũ							
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	Т	e m	р	e	r e	a t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	788 1215 668 512 352 186 186 186 160 142 132 122 104 84 61 32 30 29 28 27 26 25 25 24 23	788 1215 1196 1079 889 675 593 504 416 325 275 257 219 181 141 97 46 42 39 37 35 33 32 31 30	788 1215 1196 1196 1086 963 832 6821 481 410 337 239 205 173 239 205 173 239 205 173 239 205 173 239 205 173 239 205 173 239 205 27	$\begin{array}{c} 788\\ 1215\\ 1196\\ 1196\\ 1196\\ 1196\\ 1196\\ 1196\\ 1196\\ 1073\\ 901\\ 705\\ 649\\ 426\\ 224\\ 219\\ 194\\ 226\\ 244\\ 219\\ 194\\ 238\\ 36\\ 33\\ 29\\ 28\end{array}$	$\begin{array}{c} 788\\ 987\\ 1079\\ 1210\\ 1196\\ 1197\\ 1196\\ 1209\\ 1209\\ 1209\\ 1209\\ 1021\\ 910\\ 790\\ 639\\ 582\\ 525\\ 467\\ 408\\ 348\\ 334\\ 318\\ 303\\ 275\\ 247\\ 218\\ 191\\ 163\\ 132\\ 100\\ 66\\ 42\\ 38\\ 36\\ 34\\ 32\\ 31\\ 30\end{array}$	$\begin{array}{c} 788\\ 987\\ 1079\\ 1210\\ 1196\\ 1197\\ 1196\\ 1209\\ 1213\\ 1209\\ 1213\\ 1219\\ 1$	$\begin{array}{c} 788\\ 987\\ 1079\\ 1210\\ 1196\\ 1197\\ 1196\\ 1297\\ 1206\\ 1297\\ 1$	788 987 1079 1210 1196 1197 1197 1206 1209 1213 1219 1225 1229 1213 1225 1229 1233 1235 1237 1238 1240 1189 1169 1169 1169 1128 1265 907 846 724 659 592 525 457 388 338 308 253 224 196 128 128 253 224 196 128 253 224 253 224 196 128 253 224 196 128 253 253 253 253 253 253 254 254 253 254 255 25

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.12 \text{ m}^{1/2}$ 

-	-	b								
т		0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
đ		18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0	
Time		т	e	m p	e	r	a t	u	r	е
h										
h 0.05 0.10 0.15 0.20 0.35 0.45 0.60 0.65 0.60 0.65 0.60 0.65 0.60 0.65 0.60 0.60		T 812 1250 688 525 359 186 186 177 169 160 141 132 122 103 82 58 28 27 26 25	e 812 1250 1230 912 691 606 513 421 326 217 179 138 938 35 331 30 28 275 26 217 26 217 256 217 256 217 256 227 256 217 256 227 226 226	m p 812 1250 1230 1230 1230 1230 1230 1230 1230 1230 230 230 236 202 167 130 90 40 38 35 333 32 31	e 812 1250 1239 1209 1209 1209 1209 1209 1209 1200 266 241 137 78 38 35 32 30	r 812 1014 1109 1245 1230 1230 1230 1230 1230 1231 1239 1231 1239 1231 1239 1231 1239 1231 1239 1233 1239 1233 1239 1233 1239 2123 2158 1598 960 342 329 213 1289 329 213 329 213 329 329 329 329 329 329 329 329 329 329 329 329 329 329 329 329 329 329 32 329 32 329 32 329 32 329 32 329 329 32 329	812 1014 1109 1245 1230 1230 1230 1230 1230 1231 1239 1242 1247 1136 986 905 820 728 683 635 587 91 442 343 314 258 230 202 174 145 113 80 37 34 32 31	812 1014 1109 1245 1230 1230 1230 1230 1231 1232 1253 1255 1255 1255 1255 1255	r 812 1014 1245 1230 1230 1230 1230 1230 1230 1231 1237 1254 1257 1258 1257 1258 1257 1258 1259 1260 1163 1142 1087 1031 972 8499 724 658 590 521 452 303 275	e
5.00 5.20							29	159 137		
5.20 5.40								137 115	247 219	
5.60								93	191	
5.80 6.00								68 36	163 132	
								- ·	-	

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# <u>C6</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.01 \text{ m}^{1/2}$ 

-	-	° t							
Т		0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ		1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time		Т	e m	р	е	r a	t	u	r e
ß									
h 0.05 0.10 0.25 0.20 0.35 0.40 0.45 0.60 0.65 0.70 0.80 0.90 1.00 1.20 1.20 1.50 1.50		315 453 224 110 224 100 97 95 246 57 45 421 40 57 45 421 40	315 453 445 361 291 247 217 186 174 128 124 128 111 94 76 73 69 67 64	315 453 445 456 430 303 303 257 206 182 168 159 123 100 86	3153 4457 44576 44576 44574 44574 41853 3192 2216 2216 2216 2216 2216 2216 2216 2	315 376 411 456 460 472 484 494 509 525 539 481 449 412 384 350 330 3309 288	3176 3171 446 78 4900 555 552 94 448 400 555 552 495 3806 444 400 555 555 552 495 3806 444 400 555 555 555 495 3806 1 4006 1	315 3761 4560 4844 5025 5339 5566 591 591 591 591 591 591	315 376 456 460 48 490 509 523 9336 514 44 624 634
1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.40 3.60 4.20 4.20 4.60 4.80 5.00 5.40 5.60 5.60 5.80 6.00		39 38 37 36 35 34 32 32 31	62 60 7 5 5 5 5 5 5 5 5 4 4 4 4 3 2 1 0 9 8 8 7 7 6 5 5 5 4 4 4 4 3 8 8 7 7 6 5 5 5 4 4 4 4 3 8 8 7 7 6 5 5 5 5 5 5 4 7 6 6 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	83077507418643109876543221	182 170 159 146 121 104 97 97 83 80 77 74 72 66 65 63 60 59 58 56	265 257 248 240 225 210 195 181 166 151 135 119 109 103 99 95 92 89 86 84 82 80 78 76	375 358 341 323 284 267 250 234 219 204 189 173 158 142 125 119 105 102 99 96 93 91	571 558 544 529 498 467 432 393 352 304 289 275 261 249 236 221 199 187 175 163 151 138	634 636 635 633 623 610 593 574 552 531 507 482 456 428 398 367 342 323 305 287 270 253 287 270 253 236 218

### <u>D1</u>

Openin	ng Factor A	$\sqrt{H/A_t} = 0$	0.02 m <sup>1/2</sup>						
T		0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ		3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0
Time h		Т	e m	р	e	r a	t	u	r e
$\begin{array}{c} n\\ 0.05\\ 0.10\\ 0.15\\ 0.20\\ 0.30\\ 0.45\\ 0.50\\ 0.45\\ 0.50\\ 0.65\\ 0.70\\ 0.65\\ 0.70\\ 0.80\\ 1.20\\ 1.20\\ 1.20\\ 1.50\\ 1.60\\ 1.50\\ 1.60\\ 2.20\\ 2.40\\ 2.60\\ 2.80\\ 3.20\\ 3.80\\ 4.20\\ 4.60\\ 4.80\end{array}$		444 631 358 307 232 145 142 131 125 113 108 103 91 79 67 53 48 46 44 43 41 40 38 36 35 34 33 22	$\begin{array}{c} 444\\ 631\\ 616\\ 574\\ 990\\ 375\\ 224\\ 7\\ 206\\ 161\\ 111\\ 84\\ 77\\ 76\\ 65\\ 60\\ 57\\ 45\\ 44\\ 41\\ 40\\ 398\\ 37\\ 37\\ 37\\ 37\\ 37\\ 37\\ 37\\ 37\\ 37\\ 37$	44316599972844692534210749950663077666685555594765	4316594417655445248869258137939159517429753	444 520 566 645 645 672 686 717 752 719 673 492 400 752 9623 400 752 9623 400 752 208 207 166 145 228 208 207 166 145 228 208 207 166 145 289 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 207 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 208 209 209 208 209 209 208 209 209 209 209 209 209 209 209 209 209	445 5665 6682 717 752 783 740 6509 541 593 5109 753 2097 1452 2310 1452 25 1452 1452 1452 1452 1452 1452 1	4420 5666 6458 667774 77778 8818 8888 8888 887777538 8818 887776503 8884 77531136 922 22425	$\begin{array}{c} 446\\ 526\\ 644\\ 6572\\ 686\\ 701\\ 740\\ 753\\ 801\\ 818\\ 847\\ 801\\ 822\\ 877\\ 801\\ 882\\ 877\\ 882\\ 877\\ 481\\ 860\\ 838\\ 812\\ 748\\ 812\\ 877\\ 481\\ 638\\ 812\\ 748\\ 812\\ 877\\ 481\\ 638\\ 812\\ 748\\ 812\\ 877\\ 481\\ 638\\ 812\\ 877\\ 481\\ 836\\ 812\\ 836\\ 812\\ 838\\ 838\\ 812\\ 838\\ 838\\ 838\\ 838\\ 838\\ 838\\ 838\\ 83$
5.00 5.20 5.40 5.60 5.80			36 35 35 34 34 34	44 43 42 41 40	61 59 58 56 55 54	83 80 78 76 74	100 97 93 90 87	209 193 177 161 145	335 310 287 264 241
6.00			33	40	54	72	85	127	217

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor  $A \cdot \sqrt{H/A} = 0.02 \text{ m}^{1/2}$ 

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#### <u>D2</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor  $A\cdot\sqrt{H}/A_{\rm t}$  = 0.04  $m^{1/2}$ 

	υ υ							
т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time	т	e m	р	e	r a	t	u	r e
h								
0.05	547	547	547	5 <sup>1</sup> 47	547	547	547	547
0.10	797	797	797	797	660	660	660	660
0.15	456 387	802 752	802 825	802 825	734 826	734 826	734 826	734 826
0.20 0.25	286	638	851	851	837	837	837	837
0.30	172	510	799	870	862	862	862	862
0.35	169	470	729	889	881	881	881	881
0.40	161	416	643	906	898	898	898	898
0,45	154 147	358	552	921 851	915	915 930	915	915 930
0.50 0.60	132	295 270	515 420	740	930 955	930 955	930 955	950 955
0.65	124	255	370	674	966	966	966	966
0.70	117	241	319	603	91 <sup>4</sup>	974	974	974
0.80	102	210	293	559	849	990	990	990
0.90	87	181	266 238	507	774 688	930 887	1006 1019	1006 1019
1.00	70 51	150 116	230	455 400	622	839	1019	1019
1.20	48	80	182	343	580	785	1042	1042
1.30	45	77	154	321	539	725	1009	1051
1.40	43	72	123	298	495	661	994	1060
1.50 1.60	42 40	69 65	91 87	277 255	452 406	627 592	978 949	1068 1075
1.70	40 39	62	.83	234	360	592 557	949 918	1045
1.80	37	60	79 79	211	344	521	886	1032
1.90	36	57	75	190	328	484	851	1020
2.00	35	55	72	169	313	446	815	1006
2.20 2.40	34 32	51 48	67 62	123 95	286 260	368 337	743 672	969 929
2.60	31	46	58	87	234	308	597	887
2,80	30	43	55	81	209	281	519	841
3.00	29	42	52	76	185	255	437	793
3.20 3.40		40 39	50 48	72 69	159 133	229 204	380 354	745 696
3.60		39 37	46	65	105	204 179	330	644
3.80		36	կկ	62	89	153	308	591
4.00		35 34	43	60	84	126	288	537
4.20		34	42	57	80	98	269	481
4.40 4.60		34 33	40 39	55 53	76 73	93 88	250 231	423 378
4.00		32	38	51	70 70	84	231	348
5.00		32	37	50	68	81	194	320
5.20			37	48	65	77	176	293
5.40			36	47 1	63	75 70	157	267 01-0
5.60 5.80			35 34	46 44	61 59	72 70	138 118	240 214
6.00			34 34	44	57	68	98	188
			-		• •	-		

<u>D3</u>

<u>D4</u>

	τ								
Т	0.1	0,2	0.3	0,5	0.75	5 1.0	1.5	2.0	
đ	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0	
Time h	Т	е	m p	e	r	a t	u	r e	
0.05 0.10 0.15 0.20 0.25 0.35 0.40 0.45 0.50 0.65 0.70 0.900 1.10 1.20 1.300 1.20 1.40 1.500 1.670 1.900 1.200 1.400 1.9	613 903 525 429 312 181 179 170 162 154 138 129 121 104 88 69 46 44 42 40 38 37 35 34 33 32	903 911 853 727 573 519 453 384 313 284 268 251 218 185 150 112	613 99116 9988 7656 543934 444 21811 1877 766 555 5453 443 10 976 5543 333 222 2111 1877 766 555 5453 140 9776 5543 333 333 333 333 333 333 333 333 33	613 9914 9999 9006 10348 2966674 4332018 552852 16212 60628 5520865324 452222 2085212 60628 55208653221 444444 409833 33322222852122 6066285520866532214 44444444 409833 333222228521226 55555466532214 4444444444 409833	616 746 9379910094 10056215895913331658322222115297766665554333222221152977666655542097	6136 74612 937991009463320101008366254444433166444702229600222397718652665755 102100836636626657544566447022296002223977186562605755	614 746 831 937 949 971 909 1024 1053 1096 1053 1096 1113 10835 1096 11123 1080 1040 1040 1040 10048 850 1075 1220 240 2212 283 165 122 203 165 220 240 2212 283 165 122 203 165 123 165 123 1060 10040 10040 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1230 1005 1200 1005 1200 1005 1200 1005 1200 1005 1200 1005 1200 1005 1200 1005 1200 1005 1115 1200 1005 1005	614 746 831 937 949 971 909 1024 1053 1062 1070 1083 1095 1123 1136 1142 1147 1002 1076 970 921 870 877 765 711 654 595 536 475 337 279 252 198 171	

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A\* $\sqrt{H/A}_t$  = 0.06 m<sup>1/2</sup>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$  = 0.08 m<sup>1/2</sup>

	v							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	Т	e m	р	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.45 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.40 2.40 2.40 2.40 2.40 2.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	642 975 560 454 185 183 174 166 158 140 132 105 67 43 33 32 30	642 975 983 916 778 607 547 399 321 289 275 220 185 220 185 220 185 55 52 50 44 55 52 54 33 44 33 34	642 9783 1030 9566 1030 957 1030 1030 1030 1030 1030 1030 1030 103	642 983 1030 9943 205 8 7671 205 907 384 1502 3722 11502 75 1864 444 444 438 333 3332 2222 11502 3725 1864 444 444 438 3335 334 433 3355 443 3355 35555 3555 3555 3555 3555 3555 3555 3555 35555 35555 35555 35555 35555 35555 35555 35555 35555 35555 35555 35555 355555 355555 355555 355555 355555 3555555	642385001104582522109633826022704655739631975421	642 888 1020 10458 100 100 100 100 100 100 100 100 100 10	642 803 888 1005 1020 1041 1058 1075 1095 1109 1122 1134 1144 1153 1160 1117 1095 1071 1032 993 952 910 866 782 698 609 518 423 314 293 254 214 196 177 137 157 137 157 83 71	642 803 888 1005 1020 1058 1075 1095 11006 11039 9908 8848 7716 6577 5955 533 4694 3577 298 271 243 2158 161

<u>D5</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.12 m<sup>1/2</sup>

<u>E1</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.01 \text{ m}^{1/2}$ 

Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time	Τ	e m	ŋ	e	r a	t	u	r e
h 0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.20 2.40 2.60 2.20 2.40 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	312 282 282 1307 127 112 199 56 768 55 296 46 42 0 38 37 54 28 20 210 2127 212 20 2127 20 20 20 20 20 20 20 20 20 20 20 20 20	$\begin{array}{c} 312\\ 472\\ 465\\ 444\\ 321\\ 302\\ 275\\ 242\\ 210\\ 197\\ 188\\ 163\\ 148\\ 108\\ 894\\ 80\\ 76\\ 73\\ 708\\ 641\\ 552\\ 947\\ 531\\ 140\\ 398\\ 376\\ 554\\ 33\\ 353\\ 33\\ 35\\ 33\\ 33\\ 35\\ 33\\ 33\\ 3$	312 4682 496757308260482691272885266185520875443200988 3125524475730822221119998852626185520875443200988376	312 446 489 467 88 666507 1332222221 18118002505887768530865555209 49887768530865555209	$\begin{array}{c} 312\\ 392\\ 486\\ 488\\ 502\\ 516\\ 529\\ 534\\ 547\\ 572\\ 580\\ 562\\ 499\\ 458\\ 334\\ 308\\ 273\\ 220\\ 488\\ 172\\ 220\\ 108\\ 102\\ 96\\ 87\\ 73\\ 168\\ 66\end{array}$	$\begin{array}{c} 312\\ 3928\\ 4882\\ 515\\ 55\\ 55\\ 588\\ 68\\ 44\\ 7\\ 391\\ 990\\ 1\\ 38\\ 32\\ 22\\ 22\\ 22\\ 29\\ 58\\ 01\\ 1\\ 1\\ 1\\ 1\\ 1\\ 1\\ 1\\ 9\\ 8\\ 8\\ 7\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 8\\ 7\\ 9\\ 8\\ 8\\ 8\\ 7\\ 8\\ 8\\ 8\\ 8\\ 7\\ 8\\ 8\\ 8\\ 8\\ 8\\ 7\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\ 8\\$	$\begin{array}{c} 312\\ 392\\ 428\\ 486\\ 488\\ 502\\ 516\\ 529\\ 534\\ 757\\ 588\\ 604\\ 620\\ 632\\ 646\\ 640\\ 628\\ 616\\ 603\\ 526\\ 491\\ 456\\ 528\\ 450\\ 203\\ 272\\ 258\\ 244\\ 230\\ 714\\ 191\\ 165\\ 152\\ 138\\ 124\end{array}$	$\begin{array}{c} 312\\ 392\\ 486\\ 488\\ 502\\ 516\\ 529\\ 537\\ 588\\ 604\\ 620\\ 632\\ 366\\ 665\\ 675\\ 686\\ 668\\ 666\\ 6624\\ 629\\ 974\\ 394\\ 460\\ 74\\ 394\\ 9330\\ 900\\ 272\\ 254\\ 889\\ 460\\ 749\\ 359\\ 3309\\ 272\\ 254\\ 217\\ 199\end{array}$

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor  $A \cdot \sqrt{E}/A_t = 0.02 \text{ m}^{1/2}$ 

Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<u>q</u>	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0
Time h	T	e m	р	е	r a		u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	41 6378 3249 161 139 126 113 108 70 56 44 40 375 329 30	441 637 642 609 433 223 220 179 220 179 220 179 220 179 220 179 220 179 220 179 220 179 220 179 220 179 220 179 220 200 179 220 200 179 23 8 83 75 18 66 38 29 46 32 20 37 65 43 22 32 22 33 32 32 32 32 32 32 32 32 32	467465337739661194063061149395170506530864431098765333333333333333333333333333333333333	441 642 668 7021 775 755 54 453 3001 302 244 208 172 1307 88 76 630 855 51 54 454 454 454 455 445 445 445 445	442 598 671 6793 748 7798 6714 792 8075 608 5520 485 5520 4851 5885 208 208 430 140 100 88 49 751 866 319 208 1430 140 100 88 49 751 866 31 952 208 1430 100 100 88 155 866 31 952 208 167 167 177 18 167 177 18 167 177 18 167 177 18 177 177 18 177 18 177 177 177 1	442 598 671 6792 778 802 802 7739 803 7739 803 772 802 802 772 803 772 803 772 803 772 803 772 803 772 772 803 772 803 772 772 803 772 772 772 772 772	$\begin{array}{c} 442\\ 536\\ 671\\ 6793\\ 7131\\ 804\\ 5598\\ 8135\\ 8269\\ 8135\\ 8269\\ 8277\\ 7668\\ 837\\ 75668\\ 742\\ 337\\ 568\\ 837\\ 542\\ 222\\ 242\\ 210\\ 194\\ 171\\ 144\\ 126\\ 108\\ 108\\ 108\\ 108\\ 108\\ 108\\ 108\\ 108$	$\begin{array}{c} 443\\ 536\\ 598\\ 671\\ 6743\\ 7131\\ 7632\\ 8045\\ 8533\\ 8691\\ 910\\ 910\\ 9269\\ 901\\ 2825\\ 883\\ 8554\\ 901\\ 111\\ 6308\\ 317\\ 940\\ 333\\ 2841\\ 238\\ 215\\ 212\\ 212\\ 212\\ 212\\ 212\\ 212\\ 212$

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### <u>E2</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A· $\sqrt{H}/A_t$  = 0.04 m<sup>1/2</sup>

	U							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time h	Ţ	e m	р	e	r a	t	u	r e
$\begin{array}{c} 11\\ 0.05\\ 0.10\\ 0.25\\ 0.35\\ 0.45\\ 0.50\\ 0$	547 844 476 305 183 175 165 143 165 143 1237 91 455 186 34 33	547 844 833 7681 5506 5445 2866 2518 82 773 65 61 82 773 65 61 82 56 48 41 337 35 333 32 31	544 8317 895 555 492 224 184 52 184 53 224 184 53 224 184 53 88 76 93 93 94 63 93 96 63 93 96 63 93 87 54 93 33 22 43 32 24 15 20 88 76 95 55 93 96 55 55 93 87 68 55 55 93 22 43 32 32 32 32 32 32 32 32 32 32 32 32 32	544 8317 868 99998 71382578712863197579150595207532098731 111887765952075320983333	547 689 755 869 919 955 1955 8804 1955 8804 1955 405 5400 3332 22203 1753 125 97 805 708 61 754 208 46 542 805 125 708 542 805 125 805 708 542 805 805 805 805 805 805 805 805 805 805	$\begin{array}{c} 547\\ 689\\ 956\\ 999\\ 997\\ 1003\\ 1025\\ 755\\ 869\\ 999\\ 9956\\ 1012\\ 9917\\ 1025\\ 755\\ 800\\ 102\\ 555\\ 901\\ 102\\ 224\\ 221\\ 950\\ 112\\ 112\\ 816\\ 72\\ 630\\ 75\\ 555\\ 901\\ 102\\ 224\\ 221\\ 950\\ 112\\ 112\\ 816\\ 72\\ 630\\ 75\\ 555\\ 901\\ 102\\ 224\\ 221\\ 950\\ 104\\ 55\\ 55\\ 901\\ 102\\ 224\\ 221\\ 950\\ 104\\ 55\\ 55\\ 55\\ 901\\ 102\\ 222\\ 221\\ 950\\ 104\\ 55\\ 55\\ 55\\ 901\\ 102\\ 222\\ 221\\ 950\\ 104\\ 55\\ 55\\ 55\\ 55\\ 901\\ 102\\ 222\\ 221\\ 950\\ 104\\ 55\\ 55\\ 55\\ 55\\ 55\\ 55\\ 901\\ 102\\ 222\\ 221\\ 950\\ 104\\ 55\\ 55\\ 55\\ 55\\ 55\\ 55\\ 55\\ 55\\ 55\\ 5$	$\begin{array}{c} 547\\ 689\\ 9557\\ 99961\\ 9996\\ 10013\\ 100563\\ 00112\\ 9998\\ 8561\\ 55622\\ 2233\\ 2225\\ 237\\ 9916\\ 1012\\ 1038\\ 7663\\ 814\\ 1686\\ 1252\\ 233\\ 2225\\ 237\\ 191\\ 1623\\ 233\\ 2225\\ 237\\ 191\\ 1623\\ 233\\ 2225\\ 237\\ 191\\ 1623\\ 233\\ 2225\\ 235\\ 2222\\ 235\\ 235\\ 235\\$	547 689 759 869 899 939 971 1037 1037 1037 1037 1076 1086 1011 1027 1037 1037 1037 1037 1037 1076 1086 511 1029 835 442 947 760 3258 272 240 195 225 225
6.00				35	45	52	83	169

<u>E3</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor  $A\cdot\sqrt{n}/A_{\rm t}$  = 0.06  $\rm m^{1/2}$ 

	•							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	ą	е	r a	, t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	615 929 532 192 189 170 161 144 135 126 108 91 72 49 36 33 30 30	615 929 937 882 759 608 546 406 527 228 1514 739 64 60 563 50 46 18 34 32 32 32 32 32 32 32 32 32 32 32 32 32	619993787999875490094886531194774066294951997643332221819477406629495133333332	615 9237 958 9013 10050 10000 10050 10000 10000 10000 100000 100000000	$\begin{array}{c} 615\\ 766\\ 957\\ 10024\\ 1057\\ 10025\\ 10024\\ 1005\\$	$\begin{array}{c} 615\\ 766\\ 959\\ 9024\\ 10024\\ 10024\\ 10025\\ 10024\\ 10025\\ 10$	$\begin{array}{c} 615\\ 766\\ 854\\ 959\\ 976\\ 10024\\ 1058\\ 1070\\ 1085\\ 1097\\ 1107\\ 1112\\ 1129\\ 1134\\ 1097\\ 1017\\ 1011\\ 9734\\ 893\\ 850\\ 766\\ 415\\ 329\\ 268\\ 248\\ 229\\ 102\\ 154\\ 133\\ 112\\ 92\\ 69\end{array}$	615 766 854 960 976 1002 1024 1058 1070 1085 1091 1097 1107 1115 1122 1134 1139 1143 1147 1151 1129 1134 1143 1060 1015 967 917 864 8095 698 640 581 947 347 347 329 458 394 347 329 458 394 347 329 183 156

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# <u>E4</u>

<u>E5</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A· $\sqrt{H}/A_t$  = 0.08 m<sup>1/2</sup>

		0							
	T	0.1	0,2	0.3	0,5	0.75	1.0	1.5	2.0
	q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
	Time	т	e m	р	e	r a	t	u	r e
	h								
	0.05	663	663	663	663	663	663	663	663
	0.10	999 577	999	999	999	831	831	831	831
	0.15 0.20	577 468	1005 942	1005 1026	1005 1026	913 1027	913 1027	913 1027	913 1027
	0.25	335	805	1025	1055	1043	1043	1043	1043
	0.30	194	638	988	1076	1067	1067	1067	1067
	0.35	191	568	897	1093	1086	1086	1086	1086
	0.40	181	494	792 669	1107	1101	1101	1101	1101
	0.45	172	417	669	1118	1113	1113	1113	1113
,	0.50	163	33 <u>4</u>	608 480	1023	1122	1122	1122	1122
	0.60 0.65	145 135	295 277	480 4 <u>1</u> 4	872 788	1133 1138	1133 1138	1133 1138	1133 1138
	0.70	126	259	345	693	1068	1143	1143	1143
	0.80	108	223	309	624	979	1151	1151	1151
	0.90	90	187	276	557	879	1069	1158	1158
	L.00	69	150	244	489	767	1005	1164	1164
	1.10	44	109	210	419 al D	677	938	1169	1169
	1.20	41	65	179 145	348	623 571	867 792	1173 1125	1173 1177
	1.30 1.40	39 37	61 56	108	322 297	517 517	708	1102	1180
	1.50	36	53	70	273	463	664	1076	1183
	L.60	33	49	66	249	408	621	1037	1186
	L.70	31	47	62	224	352	577	996	1143
	1.80	30	45	58	200	335	532	954	1124
	L.90	29 28	43	54	176	319 304	489 444	911 866	1105 1086
	2.00 2.20	20	41 37	52 47	152 98	277	353	778	1038
	2.40		34	43	67	250	321	692	987
	2.60		32	39	60	222	293	601	934
	2.80		31	37	54	195	266	507	879
	3.00		30	35	50	169	238	411	821
	3.20				46	140	210 184	349 324	765 706
	3.40 3.60				42 40	109 79	157	324 303	706 645
	3.80				38	58	127	283	582
	+.00				36	54	97	264	519
	+ <b>.</b> 20				45	49	64	244	455
1	4.40				34	46	58	225	389
	•.60				33	43	55	206	341
	⊧•80 5•00				32 31	41 39	50 47	187 168	312 285
	5.20				LC	39 37	44 44	148	258
	5.40					36	42	128	231
	<b>i.</b> 60					35	40	106	203
1	.80					34	39	85	177
é	i.00					33	37	60	148

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor  $A \cdot \sqrt{H} / A_t = 0.12 \text{ m}^{1/2}$ 

	τ										
т	0.1	0.2		0.3	0.5	Ο,	75	1.0	1.5		2.0
q	18.0	35.0		54.0	90.0	135.	.0	180.0	270.0	36	0.0
Time	т	е	m	р	е	r	a	t	u	r	е
h 0.10 0.225 0.225 0.225 0.225 0.225 0.225 0.225 0.225 0.225 0.225 0.225 0.225 0.255	709 1086 610 494 348 192 183 173 164 5 135 126 87 66 38 34 32 29 28 25	$\begin{array}{c} 709\\ 1086\\ 1087\\ 1013\\ 860\\ 670\\ 594\\ 336\\ 297\\ 278\\ 260\\ 222\\ 184\\ 145\\ 102\\ 555\\ 46\\ 43\\ 41\\ 39\\ 37\\ 36\\ 34 \end{array}$		$\begin{array}{c} 709\\ 1086\\ 1087\\ 1108\\ 1130\\ 944\\ 827\\ 6927\\ 481\\ 3074\\ 241\\ 2074\\ 138\\ 54\\ 100\\ 58\\ 50\\ 47\\ 442\\ 35\\ 33\\ 30\end{array}$	$\begin{array}{c} 709\\ 1086\\ 1087\\ 1108\\ 1130\\ 1148\\ 1169\\ 1175\\ 1069\\ 903\\ 811\\ 709\\ 637\\ 565\\ 494\\ 420\\ 345\\ 2270\\ 245\\ 220\\ 195\\ 171\\ 145\\ 89\\ 48\\ 44\\ 037\\ 353\\ 332\\ 31 \end{array}$	70 89 980 112 115 116 117 117 116 117 117 119 111 100 788 637 526 40 314 327 246 215 107 111 100 788 331 327 246 131 107 44 333 334 333 343 353 343 353 345 355 355	95621562870134482603571514631443)740,54	$\begin{array}{c} 709\\ 899\\ 9106\\ 1121\\ 1155\\ 1166\\ 1178\\ 1187\\ 1190\\ 1205\\ 1037\\ 6690\\ 890\\ 813\\ 7622\\ 536\\ 943\\ 317\\ 22234\\ 206\\ 179\\ 150\\ 892\\ 74\\ 336\\ 353\\ 32\\ 31\end{array}$	709 899 985 1122 1141 1155 1166 1178 1194 1205 1209 1215 1260 1209 1215 1260 1209 1215 1209 2215 1209 2299 2200 2299 2200 1823 129 77 48	$\{ \{ \{ \{ \} \} \} \} $	73986211562870405925802455432713335393811458369247590

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# <u>E6</u>

<u>F1</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.10$ Opening Factor A·VH/A<sub>t</sub> = 0.01 m<sup>1/2</sup>

T	0.1		0.3			•5		1.0		2.0	
q	1.5		4.5		7	•5		15.0		30.0	
Time	Тe	m	ą	е	r	a	t	u	r	е	
h .											
0.05	252		252 ha 0		25 41	52		252 350		252 350	
0.10 0.15	418 297		418 456		41 49	10 56		413		413	
0.20	261		489		48			480		480	
0.25	208		506		50			498		498	
0.30 0.35	148 129		483 452		51 52			513 524		513 524	
0.40	114		409		53			535		535	
0.45	103		359		51			540		540	
0.50	94 81		329 270		53 դե			544 553		544 553	
0.65	75		239		40			557		557	
0.70	70		207		36			560		560	
0.80 0.90	59 49		180 159		33 29			566 538		566 572	
1.00	38		140		26			510		577	
1.10	28		122		23			480		581	
1.20 1.30			105 88		20 18			449 414		585 589	
1.40			70		17			377		592	
1.50			50 b o		15	57		355		595	
1.60 1.70			42 37		14 13			334 314		598 585	
1.80			33		11	.8		293		576	
1.90			31		10	)5 )2		273		568 560	
2.00 2.20			29			5		251 207		538	
2.40					Ł	<u>4</u>		186		515	
2.60 2.80					3	87 83		170 155		490 463	
3.00						31		140		436	
3.20								125		409	
3.40 3.60		·						110 96		382 354	
3.80			'					80		325	
4.00								64		294	
4.20 4.40								47 40		264 231	
4.60								37		205	
4.80								35		189	
5.00 5.20								33		174 159	
5.40										144	
5.60										129	
5.80 6.00										113 98	

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.10$ Opening Factor A· $\sqrt{H}/A_{\rm t} = 0.02 \text{ m}^{1/2}$ 

Т	0.1 3.0		0.5		]	L.O		2.0			
đ	3.0		9	••0		15.0	)	30	0.0		60.0
Time h	T	е	m	р	e	r	a	t	u	r	e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	382 612 413 354 273 181 143 120 104 97 77 63 48 32		66 66 66 65 40 35 30 20 20 20 20 20 20 20 20 20 20 20 20 20	33 51 51 51 51 51 51 51 51 51 51 51 51 51		382 645 6691 77232 4897 55489 3351 600 2318 2004 1132 1133 248 351 3060 1132 3060 1132 313 3060 1132 313 312 312 312 312 312 312 312 31		5566 7777777777777777777777777777777777	227938244 <b>2</b> 2257790100545677664039753210		$\begin{array}{c} 382\\ 585\\ 676\\ 691\\ 691\\ 7125\\ 777\\ 7555\\ 777\\ 778\\ 778\\ 788\\ 602\\ 298\\ 603\\ 559\\ 933\\ 559\\ 244\\ 413\\ 331\\ 735\\ 2232\\ 2134\\ 155\\ 11\\ 155\\ 11\\ 11\\ 125\\ 11\\ 11\\ 125\\ 11\\ 11\\ 125\\ 11\\ 11\\ 125\\ 11\\ 11\\ 11\\ 125\\ 11\\ 11\\ 125\\ 11\\ 11\\ 125\\ 11\\ 11\\ 125\\ 11\\ 125\\ 11\\ 125\\ 11\\ 125\\ 11\\ 125\\ 125$

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<u>F2</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\varepsilon_{\rm res} = 0.10$ Opening Factor A· $\sqrt{H}/A_{\rm t} = 0.04 {\rm m}^{1/2}$ 

-	0	τ				
т		0.1	0.3	0.5	1.0	2.0
q		6.0	18.0	30.0	60.0	120.0
Time		т е	m p	e r a	t u	r e
h						
0.05		522	522	522	522	522 674
0.10		815 525	815 841	815 841	673 763	763
0.20		437	870	870	874	874
0.25		327	892	892	884	884
0.30		206 181	84 <u>1</u> 767	901 909	896 905	897 905
0.35 0.40		163	679	909 916	909 913	90) 913
0.45		150	574	921	919	919
0.50		139	521	855	924	924
0.60		121	4 <u>1</u> 4	731	936	936
0.65 0.70		112 105	359 302	661 585	939 941	939 941
0.80		89	264	525	944	944
0.90		72	234	469	877	948
1.00		54	205	413	824	951
1.10		33	178 149	356 296	771 713	953 955
1.20 1.30			149	272	651	957
1.40			88	250	583	959
1.50			54	229	548	960
1.60			<u>4</u> 4	208 188	514 479	962 931
1.70 1.80			39 36	168	419	931
1.90			34	147	407	903
2.00			° 32	124	370	888
2.20				79 44	294 268	851 811
2.40 2.60				37	260 244	765
2.80				34	221	720
3.00				32	198	672
3.20					176 152	627 581
3.40 3.60					128	533
3.80					102	483
4.00					76	432
4.20					46	379
4.40 4.60					38 35	325 284
4.80					33	261
5.00					32	238
5.20						214
5.40 5.60						192 169
5.80						146
6.00						120

<u>F3</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.10$ Opening Factor A· $\sqrt{H}/A_t = 0.08 \ m^{1/2}$ 

	v				
T	0.1	0.3	0.5	1.0	2.0
q	12.0	36.0	60.0	120.0	240.0
Time	Те	n v	era	t u	re
h		*			
Time h 0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00		36.0 m p 628 971 1012 1032 1043 972 878 771 647 582 389 2850 218 187 156 122 86 44 30 30	60.0 e r n 628 971 1012 1032 1043 1052 1058 1062 1066 973 822 740 648 581 516 451 384 315 290 267 244 221 198 176 153 128 75 36 31	t u 628 823 918 1025 1037 1047 1055 1060 1065 1068 1075 1077 1079 936 873 805 731 6508 527 486 444 401 314 287 262 236 211	r e 628 822 917 1025 1037 1047 1055 1060 1065 1066 1075 1077 1079 1081 1083 1084 1083 1084 1085 1084 1085 1085 1084 1085 1085 1085 1085 1085 1085 1086 1090 1048 1032 1014 996 955 857 806 750
3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00				211 186 161 133 105 76 37 32 30	750 697 642 584 527 468 407 346 301 276 251 226 201 176 150 122

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.10$ Opening Factor A·VH/A<sub>t</sub> = 0.12 m<sup>1/2</sup>

	° ,							_				
Т	0.1			•3		0.5	i		•0	2.0		
đ	18.0		54	.0		90.0	}	180	0.0		360.0	
Time	T	e	m	p	e	r	a	t u		r	е	
h												
0.05	677		67			677		67	7		677	
0.10	1060		106			.060		88			880	
0.15	657		108			.086		98			981	
0.20	514		110			106		109			1097	
0.25 0.30	364 205		111 103			.114 .121		111 111			1110 1117	
0.35	189		92			.125		112			1123	
0.40	175		80			128		112			1127	
0.45	164		67			130		112			1129	
0.50	154		60	4		021		113			1131	
0.60	135		46			859		113			1134	
0,65	125		39			771		113			1136	
0.70	115		32			672		113			1138	
0.80 0.90	98 #0		28			603		114	.0		1140	
1.00	78 55		25 22			534 465		104 98	ס ו		1142 1143	
1.10	27		19			393		90 91			1145 1144	
1.20			15			321		84			1145	
1.30			12			296		76	1		1146	
1.40			8			272		67			1146	
1.50				8		249		63			1147	
1.60 1.70			3	1		225 202		58 58			1148	
1.80						202 179		54 50			1100 1083	
1.90						154		45			1064	
2.00						128		41			1044	
2.20						73		32			996	
2.40						31		29	3		946	
2,60								26			894	
2.80 3.00								24			840	
3.20								21 19			781 726	
3.40								16			667	
3.60								13			607	
3.80								10	5		545	
4.00								$7^{1}$			483	
4.20								3:	3		419	
4.40											354	
4.60 4.80											307 282	
5.00											202 256	
5.20											230	
5.40											204	
5.60											179	
5.80											151	
6.00											122	

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<u>F5</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.35$ Opening Factor A·VH/A<sub>t</sub> = 0.01 m<sup>1/2</sup>

		,										
T		0.1		C	.3		0.5	0.5		.0		2.0
đ		1.5		1	1.5		7.5	ī	15	5.0		30.0
Time		т	е	m	р	е	r	a	t	u	r	е
h					-							
0.05 0.10 0.15 0.20 0.25 0.35 0.40 0.45 0.50 0.665 0.70 0.90 1.00 1.20 2.20 2.40 3.20 3.40 3.20 3.40 3.20 5.20 5.40 5.20 5.60		2417 2322 2493 2493 2493 2493 2493 2493 2493		33 44 44 44 44 46 33 32 22 22 22 22 22 22 14 14 14 11 12 14 11 12 14 11 11 12 11 11 11 11 11 11 11 11 11 11	41879755001169881330426048813331		241 3319 447 55466 771 4666 771 1508 863 368 31642 2227 1652 129 857 31		$\begin{array}{c} 33\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\ 4\\$	7177803155940335547004779213662996		2334454444444444444444444444444433333322222111111108

<u>F6</u>

<u>F7</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\varepsilon_{\rm res} = 0.35$ Opening Factor A· $\sqrt{H}/A_{\rm t} = 0.02 \ {\rm m}^{1/2}$ 

T	0,1		0	1.3		0.5	5	נ	.0		2.0
q	3.0		9	.0		15.0	)	30	0.0		60.0
Time	Т	e	m	р	e	r	a	t	u	r	e
h	-	-		7-	-	-		-		-	-
0.05 0.10 0.225 0.35 0.450 0.655 0.655 0.900 1.100 1.200 0.655 0.900 1.100 1.200	364 560 332 257 172 117 110 97 90 84 72 58 44 29		565 58 61 58 53 53 53 53 53 53 53 53 53 53 53 53 53	38453416530253675553		3640 568845 6625 66300 55130 3513 27344 3503 27344 1984 96 1537 1204 91 352		4 5 5 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6	2923339137763046533196274369246		3455114849249137024689122345638245787515891470234111121

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.35$ Opening Factor  $A \cdot \sqrt{H}/A_t = 0.04 \ m^{1/2}$ 

Т		0.1		C	).3		0.5	i	נ	L.O		2.0		
q		6.0	Ì	18	3.0		30.0	I	60	0.0		120.0		
Time		Т	е	m	р	е	r	8	t	u	r	е		
h														
0.05		473			73		473			73		473		
0.10		751 501			51 62		751 763		62	28 98		628 698		
0.15 0.20		407		763 770			770			90 94		784		
0.25		305		779		779 7			7	78		778		
0.30		191		7	734		783		7	30		780		
0.35		168			670				790			88		788
0.40	1.1	152			588					95 97		795 797		
0.45 0.50		141 132			507 հ65		737		79			799		
0.60	•	115			465 376		634			)3 )3		803		
0,65		107		316 329			578			05		805		
0.70		100			278		516		80			806		
0.80		85			246 219		471		80			809		
0.90 1.00	1	69 50			19 94		426 379		71 71			811 812		
1.10		30			58		328		6			813		
1.20		20			+1		276		62	27		815		
1.30			-		12		254		57			816		
1.40				ł	33 48		235		52 49			816 817		
1.50 1.60					40 39	÷	215 197		45			818		
1.70					34		178		4			793		
1.80					32		159		40			784		
1.90					30		139		31			773		
2.00 2.20							117 73		31 21			757 727		
2.40							39		25			696		
2.60							34		23	31		663		
2.80							31		2]			627		
3.00									18 16			590 555		
3.20 3.40	•								11			518		
3.60									12	22		479		
3.80									9	98		438		
4.00										13		394 21-0		
4.20										+2 35		349 301		
4.40 4.60										32		266		
4.60										31		245		
5.00												224		
5.20												203		
5.40												182 161		
5.60 5.80												138		
6.00												114		

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<u>F8</u>

Opening Factor A.V	$\overline{H}/A_t = 0.0$	8 m <sup>1/2</sup>		-520		res
T	0.1	0.3	0.5	1.0	2.0	
đ	12.0	36.0	60.0	120.0	240.0	
- Time h	Т е	m p	e r a	t u	re	
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	593 892 593 464 338 197 178 165 155 145 128 119 110 94 75 53 28	593 892 921 923 931 867 790 698 593 539 428 306 272 242 183 152 188 41 30	593 892 921 934 937 940 943 860 738 668 590 486 5301 279 258 214 193 171 148 124 72 33	593 776 848 923 933 936 942 949 944 949 944 949 954 886 725 564 495 8302 273 205 187 130 273 30 73 35 30	593 577843 926336994478 94447895579445779443930466492184995579959999988514393046649224996269995144383329184996269917268061439332918449626917249	

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\varepsilon_{res} = 0.35$ Opening Factor A· $\sqrt{H}/A_{\star} = 0.08 \text{ m}^{1/2}$ 

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# <u>F9</u>

<u>F10</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\varepsilon_{res} = 0.35$ Opening Factor  $A \cdot \sqrt{H}/A_t = 0.12 \text{ m}^{1/2}$ 

-	U					
т	0.1	0.3	0.5	1.0	2.0	
q	18.0	54.0	90.0	180.0	360.0	
Time	T e	m p	e r a	t u	r e	
h		Ŷ				
0.050 0.150 0.2250 0.450 0.6650 0.65500 0.65500 0.65500 0.65500	640 985 631 489 350 197 182 170 160 150 122 123 114 96 77 53 26	640 985 1000 1003 1010 934 847 746 628 569 447 382 315 281 250 218 187 155 120 83 36 30	640 985 1000 1003 1010 1013 1016 1018 1020 934 800 721 633 572 509 445 379 311 288 266 243 220 198 175 151 125 71 30	640 852 921 1006 1005 1012 1015 1017 1019 1021 1022 1023 1024 1025 946 893 837 776 631 595 558 519 480 439 397 311 286 261 236 211 186 160 132 103 73 31	640 852 921 1006 1015 1017 1021 1022 1023 1024 1025 1027 1030 1028 1029 1030 1030 1030 1031 992 948 909 868 824 726 678 628 5750 4044 300 2751 225 200 175 120 120 1020 1021 1022 1020 1023 1024 1025 1027 1030 1030 1031 2979 948 809 868 824 776 678 6285 5204 404 344 300 2751 1202 1020 175 1202 1030 1030 1031 1030 1031 225 200 175 1202 10200 10200 10200 10200 10200 10200 10200 10200 1	

Time Graphs of	Temperature of	f Combust	ion Gases.	Type F En	closed Spac	e. <sub>Eres</sub> = 0.60
Opening Factor	$A \cdot \sqrt{H}/A_t = 0.02$	1/2 1 m <sup>1/2</sup>				
T	0.1	0.3	0.5	1.0	2.0	
q	1.5	4.5	7.5	15.0	30.0	
Time h	Т е	m p	e r a	t u	r e	
	231 364 269 229 180 125 105 93 85 78 68 63 59 49 40 31	231 364 405 413 419 408 362 317 288 263 217 193 167 146 130 115 101 87 155 37 30	231 364 405 413 419 425 425 425 425 425 425 425 425 425 425	231 359 416 428 429 433 437 369 3420 228 439 433 437 369 50 228 49 238 210 125 103 185 50	231 311 359 416 419 428 429 433 434 439 444 445 4447 448 431 426 428 433 434 439 4447 448 431 426 429 223 353 467776 233	
4.20 4.40				34	209 185	
4.60 4.80					165 152	
5.00					140	

5.20 5.40 5.60 5.80 6.00

<u>F11</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.60$ Opening Factor A· $\sqrt{h}/A_{\rm t} = 0.02 \ {\rm m}^{1/2}$ 

Т	0.1	0.3	0.5	1.0	2.0
đ	3.0	9.0	15.0	30.0	60.0
Time h	Те	m p	e r a	t u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	348 525 374 310 240 159 136 121 111 104 91 85 80 68 55 41 27	348 525 547 560 531 433 348 254 218 378 254 218 374 154 133 154 134 31	348 525 547 560 564 570 534 464 385 353 322 289 254 217 200 186 172 158 143 128 112 97 63 37 31	34724127134790225651589987579482580229330 364724127134790225551865128899875794825580229330	344 56612 7134 790246 7801236 70022184 950354 3927266 925 755 5888 801236 700221884 950354 392726 1411119

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# <u>F12</u>

<u>F13</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.60$  Opening Factor A·VH/A<sub>t</sub> = 0.04 m<sup>1/2</sup>

_	_	τ										
т		0.1	0.1 0.3				0.9	5	1.0			2.0
đ		6.0	)	18	3.0		30.0	)	60	0.0	:	120.0
Time		т	e	m	р	е	r	8,	t	u	r	е
h												
0.05		450			50		450			50		450
0.10		708 470		708 703			708			97		597 650
0.15 0.20		381			05 09		703 709			50 20		650 720
0.25		288			10		710			L1		711
0.30		180			59		714			ιı		711
0.35 0.40		158			13		717			16		716
0.40		144 134			37 57		719 722			18 23		718 723
0.50		125			31		668			28		728
0.60		110			52		578			29		729
0.65		104			28		529			31		731
0.70 0.80		97 82			51 33		474 437		73 73			732 732
0.90		66			58		397		68			737
1.00		48			35		355		65	52		733
1.10		28		16			309		61			741
1.20 1.30				13 10			261 241		57 53			732 739
1.40				1	78		223		48	12		740
1.50				1	14		206		45			739
1.60					35		189		43			742
1.70 1.80				-	31		171 152		40 38			717 701
1.90							133		35			698
2.00							112		32			690
2.20 2.40							69 36		26			664
2.60							30		24 22			637 608
2,80									20	1		578
3.00									18			545
3.20 3.40									16 14			514 482
3.60									11			402 447
3.80									9	4		410
4.00										9		372
4.20 4.40										9 2		330 287
4.40										2 0		254 254
4.80									-	-		234
5.00												214
5.20 5.40												195 175
5.60												155
5.80												133
6.00												110

<u>F14</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.60$ Opening Factor A·VH/A<sub>t</sub> = 0.08 m<sup>1/2</sup>

T	0.1			•.3		0.5			L.O		2.0
q	12.0		36	.0		60.0	)	120	0.0		240.0
Time	т	е	m	р	е	r	a	t	u	r	e
h											
0.05	568		56			550			68		568
0.10	840	40 8 <sup>3</sup>		40		884			41 1		741
0.15	560	60 852				845			91		791
0.20	440 324	854 860		854 860				56 55		856 855	
0.25 0.30	524 190		860 802		862				52		862
0.35	172		802 732		865				54		864
0.40	160		650			867			56		866
0.45	150		556		869				58		868
0.50	142		509			787			70		870
0.60	125 116	408		683 621				72 73		872 873	
0.65 0.70	108				621 552				15 74		874
0.80	92	294 263				510		8	75		875
0.90	73		23	35		460		8:	16		880
1.00	51		20			406		71	73		877
1.10	27		17			350			27		878
1.20 1.30			ן נו			290 270		61 62			880 879
1.40				32		250		56			880
1.50				39		229		- 53			885
1.60				ŝi		208		50	01		880
1.70						188		47			852
1.80						167		43			843
1.90 2.00						145 121		40 36			831 819
2.20						70		29			787
2.40						31		26	59		753
2.60								24			717
2.80								22			679
3.00								20 17			638
3.20 3.40								15			600 560
3.60								12	7		517
3.80								10			472
4.00									1		424
4.20								3	14		374
4.40											320
4.60 4.80											282 260
5.00											237
5.20											214
5.40											191
5.60										-	168
5.80 6.00											143 116
0.00											TTO

<u>F15</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res} = 0.60$ Opening Factor A·VH/A<sub>t</sub> = 0.12 m<sup>1/2</sup>

т	0.1	-	C	.3		0.5	5	נ	0		2.0
<u>a</u> .	18.0	)	54	4.0		90.0	)	180	0.0		360.0
Time	т	е	m	р	е	r	a	t	u	r	е
h											
0.05	613				613			13		613	
0.10	931	931 931				931			18		818
0.15 0.20	599 468	99 933				933 936			68 40		867 940
0.25	339	19 9 <sup>1</sup> 43				943			37		937
0.30	192		8	70		945		9	42		941
0.35	177	77 795				947			46		946
0.40 0.45		166 703				948 949			48 49		947 949
0.50	127 178	157 596 148 544				873			49 50		949 950
0.60		130 432				752			52		952
0.65	121	121 371				680			53		952
0.70		112 306				601			53		953 051
0.80 0.90	95 76	95 274				547 491			54 92		954 955
1.00	52			14		432			45		956
1.10	,			84	370			7	93		956
1.20				52		305			35		957
1.30				18 82		282 261			72 04		957 961
1.40 1.50				35		239			58		959
1.60				28		217		5	34		959
1.70						195			99		926
1.80						173 149			52 24		912 899
1.90 2.00						124			-4 85		885
2.20						70			53		850
2.40						28			30		814
2.60									56		774
2.80 3.00									32 07		732 686
3,20									33		643
3.40								1	58		598
3.60									30		550
3.80 4.00									02 71		500 448
4.20									30		392
4.40											335
4.60											293
4.80 5.00											270 246
5.20											240
5.40											197
5.60											173
5.80											146
6.00											118

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A·VH/A<sub>t</sub> = 0.01 m<sup>1/2</sup>

	0							
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time	Т	e m	ਕੁ	е	r a	t	u	r e
h								
0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.40 3.20 3.40 3.60 3.80 4.00 4.50 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.60 4.60 4.60 4.60 4.60 4.60 4.60 5.60	$\begin{array}{c} 287\\ 473\\ 290\\ 139\\ 139\\ 122\\ 108\\ 102\\ 97\\ 88\\ 79\\ 69\\ 55\\ 49\\ 44\\ 41\\ 338\\ 36\\ 34\\ 32\\ 31\\ 30\end{array}$	$\begin{array}{c} 287\\ 478\\ 456\\ 390\\ 324\\ 3058\\ 247\\ 197\\ 18\\ 160\\ 141\\ 120\\ 80\\ 75\\ 729\\ 666\\ 50\\ 572\\ 48\\ 53\\ 33\\ 33\\ 33\\ 33\\ 33\\ 33\\ 33\\ 33\\ 33$	287 1478 4900 287 423 3320 2754 2210 1274 1204 88 807 77650 55286 3108 433 2011 1204 2011 1204 88 808 55286 3108 433 333 333 333 332	28718405505515515515649519630680383963086443109887	$\begin{array}{c} 287\\ 3440\\ 498\\ 591\\ 517\\ 512\\ 548\\ 495\\ 573\\ 548\\ 411\\ 376\\ 774\\ 932\\ 210\\ 213\\ 692\\ 445\\ 739\\ 894\\ 74\\ 766\\ 630\\ 75\\ 552\\ 522\\ 213\\ 145\\ 127\\ 398\\ 89\\ 740\\ 630\\ 75\\ 552\\ 522\\ 213\\ 145\\ 127\\ 398\\ 89\\ 740\\ 630\\ 75\\ 552\\ 522\\ 552\\ 552\\ 552\\ 552\\ 552\\$	287 3449 495 555 5547 3360 423 555 555 554 453 3370 822 2180 2467 100 40 495 1766 4	$\begin{array}{c} 287\\ 3440\\ 4995\\ 5517\\ 5517\\ 5517\\ 5517\\ 5566\\ 666\\ 667\\ 668\\ 666\\ 6555\\ 427\\ 233\\ 3155\\ 81\\ 594\\ 222\\ 294\\ 940\\ 593\\ 10\\ 1259\\ 20\\ 1259\\ 20\\ 1259\\ 1$	287 361 498 495 5017 517 527 527 583 600 621 94 639 639 6317 716 88 6353 560 996 1574 3974 2098 888 8353 3574 3978 809 8288 839 179 170 510 517 517 517 517 517 517 517 517 517 517
		•			/-			

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<u>G1</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.02 \text{ m}^{1/2}$ 

1 0.	U U							
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0
Time	Т	e m	$\mathbf{p}$	е	r a	t	u	r e
h			1				1	معا
0.05	413	413 676	413 676	413 676	413 558	413 559	413	413 559
0.10 0.15	675 391	654	654	654	602	602	559 602	602
0.20	335	606	663	663	674	674	674	674
0.25	257	513 417	672 623	672 680	668 677	668 677	668 677	668 677
0.30 0.35	173 167	376	570	707	684	684	684	684
0.40	157	376 344	536	736	730	730	730	730
0.45	149	301 258	468 440	753 717	745 766	745 766	745 766	745 766
0.50 0.60	141 126	243	368	637	796	796 808	796 808	796
0.65	119	228	327	580	808	808		868
0.70	112 100	213 187	286 260	530 493	783 747	819 849	819 849	819 849
0.80 0.90	88	161	234	458	698	762	814	814
1.00	75	134	210	418	631	738	845 866	845
1.10	62 58	106 76	186 162	373 325	572 538	713 670	866 877	866 877
1.20 1.30	55	70 75	138	302	501	626	780	807
1.40	51	70	111	281 261	464	580	787	834 856
1.50 1.60	48 46	67 62	85 80	261 241	425 386	551 521	790 772	876
1.70	43	59	78	221	345	491	769	814
1.80	41	55 53	75 72	201 182	326 310	461 431	751 724	853 856
1.90 2.00	39 38	50	68	162	295	399	695	847
2.20	35	46	61	120	270	334	635	821
2.40	33 31	42 40	55 51	92 83	245 222	303 277	575 512	790 756
2,60 2,80	29	37	47	78	199	253	512 443	719
3.00	28	35	44 5-	72	176	230	374	681 645
3.20	27 26	34 32	41 39	67 61	153 128	207 185	322 297	607
3.40 3.60	26	31	37	57	103	163	277	569
3.80		30	36	53	85	139 114	258	529 488
4.00 4.20		29 28	34 33	50 47	80 73	88	241 224	400 444
4,20 4,40		28	32	45	68	81	208	400
4.60		27	31	43	63	73 67	193	364 342
4.80 5.00		27 26	30 30	ել 40	59 56	62	177 160	321
5.20			29	38	53	59	143	301
5,40			28 28	37 36	50 48	55 52	126	281 261
5.60 5.80			20	30 35	40 46	52 50	107 90	241
6.00			27	34	44	48	73	220

<u>G2</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor  $A\cdot\sqrt{H}/A_{\rm t}$  = 0.04  ${\rm m}^{1/2}$ 

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
đ	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0	
Time h	Т	e	m p	е	r	a t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.40 2.60 2.40 2.60 2.40 2.60 2.40 3.60 3.60 3.60 3.60 4.20 4.40 4.20 4.40 5.60 5.60 5.80 6.00	545 842 403 297 187 193 125 125 127 103 71 296 36 34 329	542994 3776574933077481222111066518551853964230 114379518551853964230	546 839 852 876 837 695 543 331 296 838 207 148 110 67 292 730 332 30 30 30 30 30	5428 8351 999999999999999999999999999999999999	56768534537895 <u>99977764555644332222222111075473187543333322</u> 155345378957744576441069385938491135473187543333221	56734515199999999998888177666753967766360494792138852088765345515899999988888776680256776636049479213885208876533653	546 6953459999999999999999999999999999999999	$\begin{array}{c} 546\\ 695\\ 763\\ 855\\ 995\\ 995\\ 995\\ 995\\ 995\\ 995\\ 995$	

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<u>G3</u>

<u>G4</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.06 \text{ m}^{1/2}$ 

	0							
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	9,0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	р	e	r	a t	u	r
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	5937 59472 31874 11568 31874 11568 31874 11568 3175 310 28 3175 310 28	5947 943 7594 76040 402 2667 2493 267 2493 267 2493 267 2493 267 2493 267 2493 267 2667 2493 267 267 267 267 267 267 267 267 267 267	59999987654821794286368487304074219876 9999987654821794286368487304074219876	593 947 955 9807 1031 1051 9937 8186 464 9955 4708 1388 400 1390 52 151 1753 3320 998 22 22 27 23 3320 998 22 22 27	$\begin{array}{c} 593\\ 860\\ 960\\ 979\\ 1049\\ 1008\\ 1008\\ 710\\ 1008\\ 711\\ 1008\\ 712\\ 494\\ 493\\ 3204\\ 036\\ 209\\ 152\\ 209\\ 152\\ 998\\ 40\\ 753\\ 322\\ 28\\ 209\\ 152\\ 998\\ 40\\ 753\\ 322\\ 28\\ 209\\ 152\\ 998\\ 40\\ 753\\ 322\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 29\\ 28\\ 28\\ 28\\ 29\\ 28\\ 28\\ 28\\ 28\\ 28\\ 28\\ 28\\ 28\\ 28\\ 28$	$\begin{array}{c} 593\\780\\960\\979\\1024\\1008\\1005\\1004\\1004\\1004\\993\\800\\951\\1004\\1004\\993\\800\\951\\100\\1004\\993\\800\\951\\106\\89\\242\\29\\173\\68\\99\\351\\98\\47\\65\\44\\3298\\242\\29\\173\\68\\99\\351\\98\\47\\65\\44\\45\\44\\45\\44\\3298\\242\\29\\173\\68\\99\\351\\98\\47\\65\\44\\32\\99\\351\\98\\47\\65\\44\\45\\44\\32\\98\\37\\29\\11\\12\\89\\99\\351\\98\\47\\65\\44\\45\\44\\32\\98\\37\\29\\11\\12\\89\\99\\351\\98\\47\\65\\44\\32\\99\\351\\98\\47\\65\\44\\32\\99\\351\\98\\47\\65\\44\\45\\45$	593 780 860 960 973 998 1023 1044 1059 1008 1057 1001 1002 1045 1080 1047 1074 1074	593 780 860 973 998 1023 1044 <u>1059</u> 1008 <u>1057</u> 1001 1002 1045 <u>1080</u> 1047 1074 <u>1076</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor  $A\cdot\sqrt{H}/A_{\rm t}$  = 0.08  ${\rm m}^{1/2}$ 

	v								
Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0	
Time h	Т	e I	n p	е	r a	. t	u	r e	
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.20 3.40 3.20 3.40 3.60 3.80 4.00 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00 5.80 5.80 6.00 5.80	643 1015 574 460 327 188 186 167 157 139 130 121 103 566 41 38 36 34 32	640 1015 1009 939 802 634 5491 3300 271 2150 1413 290 1255 1842 101 559 836 33 33 33	640 1015 1009 1024 985 900 797 673 610 484 350 275 129 100 56 139 100 56 139 100 51 50 43 336 331 29 336 331 29	640 1015 1009 1024 1074 1074 998 797 710 6385 527 464 397 234 108 139 407 139 133 257 40 1394 1394 133 3333 3333 3333 3333 3333 3333 3333 333	$\begin{array}{c} 641\\ 840\\ 916\\ 10068\\ 10068\\ 10065\\ 555\\ 408\\ 333\\ 252\\ 2069\\ 1001\\ 1005\\ 555\\ 408\\ 333\\ 252\\ 2069\\ 129\\ 129\\ 142\\ 97\\ 769\\ 988\\ 766\\ 665\\ 665\\ 655\\ 120\\ 120\\ 120\\ 120\\ 120\\ 120\\ 120\\ 120$	$\begin{array}{c} 640\\ 842\\ 918\\ 1027\\ 1038\\ 1090\\ 1067\\ 10$	640 842 918 1027 1038 1062 1046 1067 1092 1058 1067 1069 1097 1117	640 842 918 1027 1038 1062 1090 1046 1067 1092 1058 1067 1087 1087 1097 1117	

<u>G5</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor  $A\cdot\sqrt{H}/A_{\rm t}$  = 0.12  $\rm m^{1/2}$ 

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time h	т	e m	ă	e	r a	. t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.50 5.00 5.20 5.40 5.20 5.40 5.20 5.40 5.20 5.40 5.20 5.40 5.20 5.40 5.20 5.40 5.20 5.40 5.60 5.80 6.00	726 1099 615 486 340 188 187 177 168 159 140 131 121 103 84 63 35 33 31 29 28 27 26	$\begin{array}{c} 723 \\ 1099 \\ 1091 \\ 1009 \\ 853 \\ 669 \\ 593 \\ 508 \\ 422 \\ 233 \\ 274 \\ 255 \\ 216 \\ 140 \\ 97 \\ 41 \\ 40 \\ 375 \\ 331 \\ 309 \\ 27 \\ 25 \end{array}$	723 1099 1091 1104 1125 1048 900 796 610 477 407 610 477 407 610 130 267 130 267 130 267 130 267 130 269 167 130 269 1091	$\begin{array}{c} 723 \\ 1099 \\ 1091 \\ 1104 \\ 1125 \\ 1144 \\ 1118 \\ 1134 \\ 1153 \\ 1052 \\ 856 \\ 772 \\ 674 \\ 614 \\ 547 \\ 477 \\ 405 \\ 333 \\ 308 \\ 284 \\ 260 \\ 236 \\ 211 \\ 163 \\ 136 \\ 79 \\ 39 \\ 34 \\ 31 \\ 28 \end{array}$	726 907 988 1117 1120 1117 1120 1117 1120 1117 1120 1117 1120 1117 1120 1120	723 907 988 1104 1117 1136 <u>1157</u> 1120 1141 <u>1158</u> 1131 1147 <u>1163</u> 1148 <u>1071</u>	723 907 988 1104 1117 1136 <u>1157</u> 1120 1141 <u>1158</u> 1131 1147 <u>1163</u> 1148 <u>1169</u>	723 907 988 1104 1117 1136 <u>1157</u> 1120 1141 <u>1158</u> 1131 1147 <u>1163</u> 1148 <u>1169</u>

<u>G6</u>