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THE ORIGIN OF Fe \(\text{ii}\) AND [Fe \(\text{ii}\)] EMISSION LINES IN THE 4000–10000 Å RANGE IN THE BD WEIGELT BLOBS OF \(\eta\) CARINAE\(^1\)

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ABSTRACT

We present numerical simulations that reproduce the salient features of the amazingly strong [Fe \(\text{ii}\)] and Fe \(\text{ii}\) emission spectra in \(\eta\) Carinae. For our studies we have used spectra obtained during the 1998 epoch observations with the Hubble Space Telescope (HST). The spectrum of the B and D Weigelt blobs dominates in [Fe \(\text{ii}\)] and Fe \(\text{ii}\) emission lines. The same observations show no Fe \(\text{i}\) or Fe \(\text{iii}\). We have compared our measurements of the strongest (≥200) [Fe \(\text{ii}\)] and Fe \(\text{ii}\) lines and blends in the spectrum with theoretical predictions. Our predictions are based on non-LTE modeling of the Fe \(\text{ii}\) atom, which includes the lowest 371 energy levels (all levels up to 11.6 eV). We have investigated the dependence of the spectrum on electron density, pumping by the blackbody-like stellar continuum, and intense Ly\(\alpha\) emission. We find that radiative pumping is essential in explaining the observed spectrum. We have identified the main pumping routes responsible for the observed Fe \(\text{ii}\) emission. Comparison between the model and observations reveals details of the radiation field. Pumping by the blackbody-like stellar radiation field from \(\eta\) Carinae explains the numerous strong [Fe \(\text{ii}\)] and Fe \(\text{ii}\) lines in the range of 4000–6500 Å. The strongest Fe \(\text{ii}\) lines in a range of 8000–10000 Å are pumped by intense Ly\(\alpha\) radiation.

Subject headings: H \(\text{ii}\) regions — ISM: individual (\(\eta\) Carinae) — line: formation — line: identification

On-line material: machine-readable table

1. INTRODUCTION

In 1986, the object known as \(\eta\) Carinae was resolved by speckle interferometry into four components within 0\(^\prime\).3 (Weigelt & Ebersberger 1986). Davidson et al. (1995, 1997) found that the brightest component A is the central star and that the other objects B, C, and D are slow-moving ejecta with extremely unusual emission-line spectra. The size of each blob is comparable to that of our solar system (Davidson & Humphreys 1997). More detailed spectra of the central source, \(\eta\) Carinae, and the BD blobs have been obtained with the Space Telescope Imaging Spectrograph on the Hubble Space Telescope (HST/STIS) (e.g., Gull, Ishibashi, & Davidson 1999). The radial velocity measurements suggest that they are slow-moving gaseous ejecta with radial velocity \(\approx -45\) km s\(^{-1}\). The angular distance from the BD blobs to the central stellar source is \(\approx 0\).1. Thus, it is large enough to be spatially resolved in HST/STIS long-slit spectra. The spectrum of the BD blobs differs from that of the central source. The stellar spectrum is rich in broad H \(\text{i}\) and Fe \(\text{ii}\) permitted lines (Hiller et al. 2001). In contrast, narrow forbidden Fe \(\text{ii}\) lines dominate in the BD spectrum.

We schematically depict in Figure 1 the position of the STIS long slit relative to that of the BD blobs and the central source (A) in the 1998 observations. We have delineated the outline of the region sampled by the STIS aperture denoted by the bold rectangle. The arrow shows direction and size (9\(^\prime\).0) to the periphery of the \(\eta\) Carinae Nebula, where the angular scale from the central star to the BD position is much smaller (\(\approx 0\).71).

Space International Ultraviolet Explorer (IUE) and HST as well as ground-based observations of the BD blobs show narrow (FWHM \(\approx 50\) km s\(^{-1}\)) and strong Fe \(\text{ii}\) and [Fe \(\text{ii}\)]
emission lines. Zethson (2001) has identified more than 2500 lines in the BD spectra recorded during 1998, 1999, and 2000 HST/STIS observations. These STIS data sets span a 1640–
10400 Å spectrum range, with most of the strong [Fe ii] and Fe ii lines found from 4000 to 10000 Å. Moreover, since there is no Fe i or Fe iii emission in the spectrum, the iron must be almost entirely in the form Fe ii.

The most peculiar features are the unusually strong Fe ii λ2507, 2509 lines identified by Johansson & Jordan (1984). Johansson & Hamann (1993) suggested that a Bowen mechanism produces the unexpected large intensities of these UV lines. Recently, Johansson & Letokhov (2001) considered an excitation mechanism for the 2507 and 2509 Å features based on photoexcitation and photoionization of Fe ii by intense Lyα radiation. In their qualitative model, Lyα is generated by the absorption of blackbody radiation from the central star and then trapped within the optically thick environment of the BD blobs.

Previous attempts (Davidson et al. 1997; Hamann et al. 1999) to derive physical conditions in BD blobs were based on an Osterbrock (1989) type line ratio analysis. They deduce electron densities of $n_e = 10^5 - 10^{10}$ cm$^{-3}$, depending on the ionic diagnostics used. Only a self-consistent quantitative model can provide a more accurate determination of electron density and other physical conditions in the BD blobs.

In this paper, in spite of considerable uncertainties in the atomic data and complexity of physical processes, we present the first quantitative model for the Fe ii and [Fe ii] emission in the BD knots. We focus on explaining the Fe ii and [Fe ii] emission spectrum in the wavelength range from 4000 to 10000 Å. Specifically we address the following questions: Why do we see so many strong Fe ii and [Fe ii] lines? What is temperature and density of the emitting region? What excitation process dominates: collisional excitation, pumping by the incident stellar continuum, or fluorescence by Lyα? Finally, can we explain the Fe ii emission by a single density model?

In addition to a complex physical environment, η Carinae also displays dramatic spectral variability. Based on spectroscopic observations of η Carinae with the 1.6 m Laboratorio Nacional de Astrofisica (LNA, Brazil) telescope between 1989 and 1996, Damineli, Conti, & Lopes (1997) have proposed that η Carinae is a stellar binary with a 5.52 yr period. Independent confirmations on the periodicity have come from X-ray observations by Ishibashi et al. (1999) and Pittard & Corcoran (2002). The optical spectrum of the Weigelt blobs mimics this periodicity with the appearance and disappearance of high excitation lines.

The 1998 epoch HST/STIS spectra we use here exhibit only low-ionization species. These spectra display no lines of higher ionization like [Ar iii], [S iii], [Fe iii], [Fe iv], and [Ne iii] seen in other epochs (e.g., Zethson 2001). As we will show, the general behavior of the Fe ii emission lines at optical wavelengths (4000–6500 Å) is mainly due to pumping via the stellar continuum radiation field, while most of strong Fe ii emission lines at near-IR wavelengths (8000–10000 Å) are due to the effects of the stellar continuum and intense Lyα radiation. We further employ a previously developed atomic model for Fe ii (Verner et al. 1999) to probe the physical conditions in the emitting region where Fe ii emission originates.

Using the unique spatial and wavelength resolution capabilities of HST/STIS, we have performed a detailed study of the Fe ii spectrum of the BD knots. This paper presents the first sophisticated numerical model of Fe ii emission as a part of fully self-consistent photoionization plasma calculations.

2. THE OBSERVED Fe ii SPECTRUM

We have analyzed the highest quality HST/STIS spectra of the compact BD ejecta of the η Carinae recorded 1998 March 19 during the 1998 broad spectroscopic minimum (Gull et al. 1999). The spectra were taken with the STIS $52''$ × $0''1$ aperture oriented at a position angle P.A. = $-27.9$ (Fig. 1). Spectral coverage is complete from 1640 to 10600 Å using ~30 grating settings with the CCD and G230MB, G430M, and G750M gratings. The spectral resolving power of the CCD is $R = \lambda/\Delta\lambda \sim 5000-10,000$. The spatial sampling is $0.0507''$ for the STIS/CCD, which corresponds to $\sim 0.05''$ at 6500 Å and broadens to $\sim 0.14''$ at 10000 Å. While the data resolve the B and D blobs up to 5000 Å, we cannot easily separate the individual spectrum of B from D at longer wavelength, as they are $\sim 0.1''$ apart. Therefore, we have extracted a single spectrum for $0.095$ to $0.40''$ offset from η Carinae A (see Fig. 1).

The diffracted and scattered light from the central source (component A) of η Carinae primarily contributes to a slowly varying continuum. The Balmer P Cygni profiles are
extended from $-1000$ to $+1000$ km s$^{-1}$, with broad wind absorptions extending to $-500$ km s$^{-1}$. Exposure times for each grating setting were chosen to obtain reasonable dynamic range without pixel saturation through the entire wavelength sampled by STIS.

Data reduction and analysis were performed using the IDL CALSTIS software package.\textsuperscript{2} It was necessary to disentangle spectra of objects with small angular separation. Thus, the standard STScI pipeline software products were inadequate. The CALSTIS tools available at NASA Goddard Space Flight Center were used to correct for the temporal variations of the optical elements and to linearize the spectral dispersion.

At wavelengths below 3500 Å, both the first-order and echelle spectral modes of the STIS (T. R. Gull et al. 2002, in preparation) demonstrate the presence of complex circumstellar and interstellar absorption along the line of sight to the Weigelt BD blobs. The spectra are dominated by absorption from the ground and low-lying levels of Fe $\Pi$ and other species. These line absorption systems contribute far less longward of 3500 Å. Because of the contamination of these absorptions, we have concentrated on an analysis of the strongest Fe $\Pi$ and [Fe $\Pi$] emission lines at wavelengths longward of 4000 Å.

3. Fe $\Pi$ LINE IDENTIFICATIONS AND MEASUREMENTS

Preliminary spectroscopic line identification suggested that most of lines/blends in BD blobs are due to singly ionized iron (Zethson 2001). The strongest [Fe $\Pi$] and Fe $\Pi$ lines are the most prominent candidates to reveal the physical conditions in the gas. Consequently, we have limited the measurements to lines with high signal-to-noise ratio. Our final line identifications are heavily influenced by the results of the model calculations.

Using the energy level data from S. Johansson (2002, private communication), we have identified the relevant Fe $\Pi$ and [Fe $\Pi$] transitions in the 1998 HST spectra. We have adopted atomic data (Quinet, Le Dourneuf, & Zeippen 1996; Zhang & Pradhan 1995; Bautista, Peng, & Pradhan 1996; Kurucz 1981; Nahar 1995; Fuhr, Martin, & Weise 1988) in our modeling. We have compiled a list of the strongest Fe $\Pi$ and [Fe $\Pi$] lines that are expected to appear in the spectrum of the BD blobs. Based upon this list, we identified more than 200 strong Fe $\Pi$ and [Fe $\Pi$] lines. For non-blended Fe $\Pi$ lines the heliocentric velocity range is $v_0 = -44 \pm 4$ km s$^{-1}$. A number of the observed emission features have several possible identifications due to uncertainties in atomic data.

The adopted Fe $\Pi$ ion model includes the lowest 371 energy levels (energies up to 11.6 eV). All the strongest observed lines are produced by transitions within the lowest 344 levels (see Table 1). The upper energy level in the model was chosen to be higher than the highest observed primary cascade produced by Ly$\alpha$ pumping.

Table 1 includes only the strongest [Fe $\Pi$] and Fe $\Pi$ lines observed in 2000–10000 Å range. Column (1) gives the line identification number (multiple identifications [possible blends] have the same number; column (2) the lists the observed wavelength in the BD rest frame, in Å; column (3) the FWHM; column (4) the observed intensity; column (5) the reddening-corrected intensity; column (6) the multiplet identification, including also information on parity; column (7) the order number of lower level; column (8) the order number of upper level; column (9) the laboratory wavelength, in Å; column (10) the energy of lower level, in cm$^{-1}$; column (11) energy of upper level, in cm$^{-1}$; column (12) the statistical weight of lower level; column (13) the statistical weight at upper level; column (14) the transition probability, in s$^{-1}$; and column (15) the measured line velocity.

We have attempted to use Fe $\Pi$ line fluxes to determine the extinction correction. However, the emission from many of the lowest states of Fe $\Pi$ is affected by collisional deexcitation. In addition, the strong emission observed from Fe $\Pi$ upper states excited by Ly$\alpha$ pumping shows mixing effects that significantly alter the observed intensities. Finally, there is no Fe $\Pi$ upper level, isolated from mixing, with primary cascades that extend across the wide observed wavelength range. Both collisional deexcitation and mixing negate any attempts to use the Fe $\Pi$ lines to estimate extinction.

Since we found no appropriate Fe $\Pi$ transitions to derive an extinction correction, we used the flux ratios of [S $\Pi$] $\lambda\lambda 4069, 4079$, and 10288. These lines all arise from the same upper level and can provide reasonable extinction estimates. We have adopted the transition probabilities of Pradhan & Peng (1995) in defining our estimates of the reddening. The observed flux ratio of these lines indicates that they have low extinction. The wide wavelength coverage of our data includes [S $\Pi$] $\lambda 10322$, but STIS sensitivity drops precipitously at this spectral range. Thus, from the measured [S $\Pi$] intensities, we derived the corrected intrinsic intensities ($I_k$) for other lines $I_k = k_3 \times I_{obs, k}$, where $k_3 = 8.3 \times 10^{-5} \lambda + 0.63$. In § 4 we show that emission of [S $\Pi$] and Fe $\Pi$ lines arise in the same emitting region, which demonstrates that the use of [S $\Pi$] to derive the extinction is relevant.

Since the lines of singly ionized iron are narrow and symmetrical in the BD spectrum, this allowed us to easily fit them with Gaussian profiles and determine the FWHM. The source of errors in the emission-line measurements is more than just statistical. Even for the strongest Fe $\Pi$ lines, where the signal-to-noise ratio is large, significant uncertainties can result from estimating the underlying stellar continuum. Blending of emission lines can also be a contributor. The measured FWHM of the observed [Fe $\Pi$] and Fe $\Pi$ lines are in a range from 30 to 158 km s$^{-1}$ (Fig. 2). For most of the single lines the FWHM is 50 ± 10 km s$^{-1}$. Lines with larger FWHM are typically blends. Blends are common because of the richness of the Fe $\Pi$ ionic spectrum as well as that of other ions. For example, the 3259 Å feature with a FWHM = 157.6 km s$^{-1}$ is an example of strong blending. Another example of line overlap is the 9204 Å feature with a FWHM = 100 km s$^{-1}$ (Fig. 3). The feature at 9204 Å is resolved and has two peaks at 9203.84 and 9205.32 Å. However, detailed Gaussian fitting suggests that a third Fe $\Pi$ line contributes. The model predictions using CLOUDY indicate that all three lines should be present and are produced by Ly$\alpha$ pumping. More details on these and other line identifications are included in Table 1.

We have selected two spectral regions, one with the weakest ([Fe $\Pi$] $\lambda 5557, 844$, Fig. 4) and one with the strongest ([Fe $\Pi$] $\lambda 7157$, Fig. 5) measured lines as examples of narrow single Fe $\Pi$ lines. Another example that demonstrates unusual physical conditions in the BD blobs is the intensity

\footnote{CALSTIS Reference Guide (D. Lindler 1999) is available at: http://hires.gsfc.nasa.gov/stis/docs/calstis/calstis.4.html.}
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Note.—Table 1 is published in its entirety in the electronic edition of the Astrophysical Journal. A portion is shown here for guidance regarding its form and content.
ratio of [N ii] λ5755 to [Fe ii] λ7157. In an H ii region like the Orion Nebula, this ratio is ∼20. The depletion of iron in the Orion Nebula is one of factors that make [Fe ii] lines weak. η Carinae is known as a nitrogen-rich object, but [Fe ii] λ7157 is ∼5 stronger than [N ii] λ5755 in the BD blobs spectrum. Studies of η Carinae physical conditions (e.g., kinematics, abundances) can provide quantitative information on nebular enrichment and test models of stellar evolution and nucleosynthesis.

Table 1 includes the reddening-corrected [Fe ii] and Fe ii intensities from 4000 to 10000 Å. After extinction correction, the strongest Fe ii line is still the [Fe ii] λ7156 ($a^4F_{9/2} - a^2G_{9/2}$, Fig. 5). Other strong, but low-excitation, lines are [Fe ii] λ5160 ($a^4F_{9/2} - a^4H_{13/2}$), and [Fe ii] λ4360 ($a^6D_{7/2} - a^6S_{5/2}$). The strongest high-excitation Fe ii lines originate in conditions that must include Lyα excitation. The strongest measured line produced by Lyα pumping, and not affected by blending, is Fe ii at λ8492 ($e^6D_{5/2} - e^6F^*_6$). However, all the observed lines arise through transitions within the 371 lowest levels.

The UV spectral region from 2000 to 4000 Å of the BD blobs is very difficult to analyze because of complexities in defining the stellar continuum and intervening absorption systems. In Figure 6 the UV emission spectrum of the BD blobs is strongly contaminated by absorption from the ground state and fine-structure levels of the resonance Fe ii UV1 multiplet (λλ2593–2620). This shows that the BD emission is veiled by high column density circumstellar gas along the line of sight. Although we include the measurements of the strongest UV emission Fe ii lines in Table 1, the measurements shortward of 4000 Å should be treated with special caution.

4. BASIC PARAMETERS FOR THE MODELING

We need to define the initial parameters for the modeling. As a first step we have searched for the traditional nebular line diagnostics in the spectrum. The FWHM of [S ii] λ6731 is twice as large as that of [S ii] λ6716. This means that the [S ii] λ6731 line must be a blend. The [N ii] λλ3063/5755 line ratio is not reliable because of large UV continuum uncertainties. There is no [O ii] emission, and the [O i] line at λ6300 is very weak, while the [O i] λλ6363 is not detected. Also, no carbon lines are identified. We find that none of the normal emission-line diagnostics are reliable indicators. Since Fe ii and [Fe ii] lines dominate the spectrum of the BD...
blobs, we cannot avoid them in our modeling. Moreover, failure to use them will certainly lead to the wrong conclusions. Intensities of the strongest and weakest Fe ii lines provide initial clues to the dominant excitation mechanisms and to the range of the electron densities in the Fe\textsuperscript{+}-emitting plasma. Specifically, the simultaneous presence of strong permitted and forbidden Fe ii lines suggests that the electron density range is between 10\textsuperscript{4} and 10\textsuperscript{7} cm\textsuperscript{-3} (Verner et al. 2000).

The range of physical parameters for numerical modeling is defined using the analysis of Davidson & Humphreys (1997), Davidson (1999), and Hillier et al. (2001). Based on these works, we assume that the total luminosity of central star is \( \sim 10\textsuperscript{40} \text{ergs s}^{-1} \) and its effective blackbody temperature is in the range \( T_{\text{BB}} = (1-3) \times 10^4 \text{ K} \). The distance from the central star (ionizing source) to the BD blobs is taken to be \( R_{\text{BD}} = 10^{16} \text{ cm} \). The BD blobs have estimated particle densities from \( 10^3 \) to \( 10^{10} \text{ cm}^{-3} \).

Long-term observations of \( \eta \) Carinae (Damineli 1998) and photoionization modeling of the central source (Hillier et al. 2001) support that photoionization and photoexcitation determine the physical conditions in the BD blobs. We have calculated a number of non-LTE models using our atomic model for Fe ii and the photoionization code CLOUDY (Ferland 1996; Ferland et al. 1998).\(^3\) To improve upon the Fe ii Werner et al. (1999) model, we have upgraded the Fe ii energy levels, and include the effects of mixing for the most important levels that are closely spaced in energy (Johansson 1984). To make Ly\(\alpha\) pumping effectively work at every pumping channel, we have dramatically increased the number of frequency points defining the Ly\(\alpha\) radiation. It is now possible using a new option in CLOUDY to obtain finer energy spacing to adequately describe the radiation field.

The effects of dust are highly variable along different lines of sight toward \( \eta \) Carinae and to its surrounding nebulosity, the Homunculus (Davidson et al. 1995; Hillier et al. 2001). The dust scattering is much stronger in the direction to the central source than in the direction to the BD blobs. Although we have corrected the Fe ii intensities for dust extinction before initiating our modeling, we do not investigate the details. Davidson et al. (1995) previously noted the possibility of gray dust in the BD direction. If it were included in models, it would cause significant Fe ii line enhancement in the near-IR range relative to that at shorter wavelengths, which is not observed. We have not included dust in our subsequent modeling.

5. THE PHYSICAL CONDITIONS IN THE Fe ii EMITTING REGION OF THE BD BLOBS

In recent years, great strides have been made in interpreting the observed \( \eta \) Carinae spectroscopic events (Damineli 1997; Davidson & Humphreys 1997) and a large amount of data have been obtained using HST/STIS (Gull et al. 1999, 2001). Still, there is no quantitative model that explains why the BD spectrum in 1998 observations is dominated by lines almost solely from singly ionized iron-group elements.

It is not a typical situation when iron emits only in Fe ii and when Fe ii are also the strongest lines in the spectrum. For example, in H ii regions like the Orion Nebula, [Fe iii], as well as [Fe ii] emission (Baldwin et al. 2000), is present. Why did we not see any sign of Fe i or Fe iii lines in the 1998 observations? Why is the Fe ii spectrum so strong and rich in the BD blobs? Is it an excitation or abundance effect? How dense is the material of the Fe\textsuperscript{+} emitting region? Why do we see so many lines from metals (like Ni ii, Cr ii, etc.)? In this section we present our photoionization model of the emitting region in the BD blobs. Our approach predicts the ionization structure of all elements, electron density, and temperature. It also determines the dominant coolant in the emitting plasma.

We have done the photoionization calculations over the range of possible electron densities to investigate how the ionization structure changes. We further adopt a blackbody temperature \( T_{\text{BB}} = 15,000 \text{ K} \) (see § 7 for more details). Figure 7 presents model results for the expected range of electron densities for \( \eta \) Carinae in the vicinity of the BD blobs. Each plotted value denotes the electron density and corresponding distance from central source to the BD position where the Fe\textsuperscript{+}/Fe ionization fraction reaches 97%. The dashed line depicts \( R_{\text{BD}} = 10^{16} \text{ cm} \) for the distance of the BD blobs from the central source. The distance \( R_{\text{BD}} \sim 10^{16.3} \text{ cm} \) (Davidson & Humphreys 1997) is even more consistent with an electron density of \( 10^6 \text{ cm}^{-3} \) for the Fe\textsuperscript{+} emitting region.

If electron density ranges from \( 10^4 \) to \( 10^7 \text{ cm}^{-3} \), at least 97% of the atomic iron exists in Fe\textsuperscript{+}. If the emitting gas has \( n_e \) within this range, it could easily explain the unusual Fe\textsuperscript{+}/Fe ionization in the BD blobs during the 1998 epoch. As seen from Figure 7, the model with \( n_e = 10^5 \text{ cm}^{-3} \) predicts the Fe\textsuperscript{+2}/Fe\textsuperscript{+} transition at a larger distance \( (R_{\text{BD}} > 3 \times 10^{17} \text{ cm}) \) than assumed for the distance of BD blobs from the central ionizing source. If the electron density is near its upper allowed limit \( (10^7 \text{ cm}^{-3}) \), then the range of distance where Fe\textsuperscript{+} ionization fraction dominates over Fe\textsuperscript{+2} is much closer to the central source than the expected BD position. Assuming constant density, we find \( n_e = 10^6 \text{ cm}^{-3} \) best provides the required dominant Fe\textsuperscript{+} fraction for acceptable distances of the BD blobs.

The model predicts an extremely sharp ionization transition between Fe\textsuperscript{+} and Fe\textsuperscript{+2} ionization fractions, with little

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\(^3\) See http://www.pa.uky.edu/~gary/cloudy.
overlap at any electron density. A sharp ionization transition is typical for collisionally ionized plasma (Verner & Yakovlev 1990). Figure 8 shows the Fe$^+$ and Fe$^{2+}$ ionization structure calculated for a typical model of the BD blobs with electron density $10^6$ cm$^{-3}$. With an electron density range of $10^5$ to $10^7$ cm$^{-3}$, the predicted total hydrogen column densities are large, varying from $3 \times 10^{22}$ to $4.5 \times 10^{23}$ cm$^{-2}$. The corresponding logarithm of ionization parameter, log $U(\text{H}) = \log[\Phi/(n_H c)]$ in our model varies from 0 to $-2$, where $\Phi$ is the total flux of hydrogen ionizing photons at the illuminated face, $n_H$ is the hydrogen number density, and $c$ is the speed of light.

The Fe$^+$ ion is the most important coolant in the BD blobs and carries up to 85% of the total cooling (Fig. 9). Figures 9 and 10 show changes in Fe $\pi$ cooling (relative to total cooling) and changes in the electron temperature with increasing cloud depth. As the distance from the central source increases, the temperature slightly decreases, and the

Fe$^{2+}$ recombines to yield Fe$^+$. When the ionization fraction of Fe$^+$ reaches 0.1, Fe $\pi$ and [Fe $\pi$] start to dominate the emission. Then the temperature drops even faster, further increasing Fe$^+$. This Fe$^+$ cooling also causes all other ions to recombine to lower ionization stages. Our modeling of the BD blobs indicates that the Fe $\pi$ emission originates at electron temperatures, $5000 \text{ K} \leq T_e \leq 7500 \text{ K}$ at distance $\sim 5.76 \times 10^{16}$ cm from central source. The derived size of Fe$^+$ emitting zone (Fig. 8) is $\sim 10$ times smaller ($\approx 10^{15}$ cm) than the distance to the BD blobs and agrees with the value derived by Davidson & Humphreys (1997).

The ionization transition for H$^+$/H$^0$ is not as sharp as that for Fe$^{2+}$/Fe$^+$ in the BD blobs. Within the Fe $\pi$ emitting region, hydrogen is significantly ionized. At the position of the Fe$^{2+}$/Fe$^+$ transition, the ionization fraction of hydrogen is $\sim 0.5$. Then at lower temperature, the H$^+$ gradually recombines to yield H$^0$ (Fig. 11). Thus, the Fe $\pi$ emission originates in a region where hydrogen is partially ionized. This also explains why the Fe $\pi$ and H $\pi$ lines are represented by very different spectral line shapes in BD blobs, since their emitting regions do not coincide.
6. ARE NLTE EFFECTS IMPORTANT?

Is it possible to apply a simple local thermodynamic equilibrium (LTE) approach to explain Fe II emission spectrum in BD blobs? We have previously verified that our model goes to LTE conditions (Verner et al. 1999) in four limiting cases: collision-dominated, radiative dominated, large optical depths, and thermal equilibrium. These tests also show that the Fe II multiple energy structure makes the ion very sensitive to collisions, continuum pumping, and line pumping. The lower and upper levels of the Fe II ion approach LTE at different electron densities. For example, the level populations of the entire atom approaches LTE in the collision-dominated case at densities \( n_e > 10^{16} \text{ cm}^{-3} \). Meanwhile, the lowest Fe II levels, at energies \( E_i/C_2 < 3 \text{ eV} \), reach LTE at \( n_e > 10^8 \text{ cm}^{-3} \). For comparison, the BD blob’s spectrum clearly shows the presence of Fe II lines originating from levels up to 11 eV. Thus, in order to adequately populate these higher energy levels, we must have either high collisional excitation (high \( n_e) or high radiation excitation (strong radiation field).

To evaluate the conditions in BD blobs, we have used departure coefficients, \( b_i = n_i/n_{i,\text{LTE}} \), where \( n_i \) is the population of \( i \)th level and \( n_{i,\text{LTE}} \) is that given by LTE from Boltzmann distribution. Figure 13 shows departure coefficients of each level in our model. The full model calculations for an electron density \( n_e = 10^{16} \text{ cm}^{-3} \) confirm that most of the high Fe II energy levels are far from LTE, while the lowest 27 level populations \( (E_i/C_2 < 3 \text{ eV}) \) are near LTE, with \( b_i \) approaching unity. The high optical depth in the BD blobs reduces the radiative rates and allows the lowest levels to approach LTE conditions. The absorption goes effectively from the first odd parity level (i.e., 64 in our model, \( \sim 4.8 \text{ eV} \)) to these lowest levels. Consequently, the absorption in Fe II can occur up to 3970.50 Å in BD conditions. The most significant departure from LTE in level population is at 64th level. This is first odd parity level, so the first permitted transition is possible from this level.

Figure 13 also illustrates that upper Fe II levels are effectively pumped by blackbody radiation. The departure coefficients of the levels pumped by Ly\( \alpha \) (in the range \( \sim 89890–89925 \text{ cm}^{-1} \), or 11–11.31 eV) are smaller than those pumped by continuum radiation. (Note that 1 eV = 8065.53851 cm\(^{-1}\)).

7. EFFECT OF BLACKBODY PUMPING ON Fe II EMISSION IN BD BLOBS

\( \eta \) Carinae is extremely luminous, but the exact nature of the central object is not yet resolved. If one represents the radiation field by a blackbody, then spectroscopic studies suggest a wide range for the blackbody temperature.
Davidson et al. (1995) advocated that the UV-visible energy distribution corresponds to a range of blackbody temperature from 20,000 to 30,000 K. Based on HST/GHRS spectra analysis at wavelengths less than 1765 Å, Ebbets, Walborn, & Parker (1997) concluded that the stellar wind is similar to a B-type supergiant with a temperature range from 13,000 to 30,000 K.

In our calculations we have included continuum pumping from the central source of η Carinae, approximated with the shape of a diluted blackbody radiation field. The initial range of the blackbody temperature is varied from 10,000 to 30,000 K. We find that the line fluxes and wavelength distribution of predicted observable Fe II transitions change dramatically depending on the electron density and blackbody temperature. The best fit between the model and the observed spectrum is achieved at blackbody temperature $T_{BB} \sim 15,000 \pm 2000$ K.

Hillier et al. (2001) compared the spectrum of the η Carinae central source with the P Cygni star HDE 316285. One of the most remarkable spectroscopic similarities between η Carinae and HDE 316285 is the Fe II stellar spectrum. The only major difference is the Fe II λ6489 feature, which is much weaker in η Carinae. The spectral variability, stellar properties, and chemical composition of HDE 316285 is typical of luminous blue variables (LBVs). Our derived blackbody temperature $T_{BB} \sim 15,000$ K for the radiation field of the BD blobs is similar to effective temperature of HDE 316285 (Hillier et al. 1998). At this temperature the radiation field effectively pumps a large number of Fe II levels from 912 to 10000 A.

In a previous study of continuum pumping, Verner et al. (2000) identified important routes for exciting the observed [Fe II] lines in the Orion Nebula. Lower effective temperature and higher densities make continuum pumping even more effective in the BD blobs. Collisions and continuum pumping more effectively populate the lower levels at higher densities to make them “launching pads” for continuum pumping. A lower temperature provides a flatter spectral energy distribution for the continuum, which is more effective in pumping over a wider range of wavelengths. The detailed Fe II analysis shows that the strong pumping routes arise from the $a^2D$ and $a^2G$ terms. The excited electrons in these levels then cascade down through intermediate levels, emitting visible photons.

We have evaluated changes in the Fe II emission spectrum by varying the electron density and blackbody temperature. We suggest three spectral ranges with various degree of sensitivity to electron density and blackbody temperature:

1. The 4000–6000 Å range contains most of the prominent Fe II and [Fe II] features. The upper levels for observed transitions are clustered near 2 eV (forbidden lines) and near 6 eV (permitted lines). The unique fit for these lines is achieved at $T_{BB} = 15,000$ K and $n_e = 10^6$ cm$^{-3}$. The evidence to support the presence of a strong blackbody radiation field is the number of strong permitted lines (e.g., $a^6S_\Delta-6P_o$ multiplet) reproduced only in models with continuum pumping.

2. The 6000–6500 Å range contains $\sim 15$ permitted lines. Their intensities are smaller than that of lines predominantly in the 4000–6000 Å range. These lines arise due to transitions between levels 4–6 eV above ground (multiplets $a^2G-z^2F_o$, $b^4D-z^4P_o$, $z^4D^c-e^4D$, and $a^6S-z^6D^o$). Their predicted intensities are slightly underestimated at $T_{BB} = 15,000$ K and $n_e = 10^6$ cm$^{-3}$. The model predicts these lines become unobservable at $n_e < 10^6$ cm$^{-3}$ at any blackbody temperature. Their predicted intensities are too strong compared to observations at $n_e \geq 10^6$ cm$^{-3}$ at blackbody temperatures $\leq 10,000$ K. Since there are no collisional excitation rates yet available for these transitions, we estimate them using the $g$-approximation (Mewe 1972).

3. The 7000–9300 Å range contains a few forbidden lines (multiplets $a^4F-a^4P$ and $a^4F-a^4G$). Their upper levels are $\leq 2$ eV above ground. For densities for 1 and 2 above, the [Fe II] $\lambda7156$ line is predicted to have an intensity half that observed. Possibly, regions with electron densities $n_e < 10^5$ cm$^{-3}$ contribute to this line. Models with blackbody temperatures $\leq 10,000$ K do not predict this line well at any density. A change of $T_{BB} = 15,000-20,000$ K does not appreciably affect the predicted intensity.

We have used the “default CLOUDY” version to predict Fe II spectrum (“model A3”). This version is included the standard set of Fe II atomic data and excitation conditions. Excitation conditions are determined by collisions, continuum pumping, and pumping by nebulae Lyα radiation. We have found the best single density model (Fig. 14) at $n_e = 10^6$ cm$^{-3}$ and $T_{BB} = 15,000$ K. This model also predicts that iron is mainly in the form of Fe I.

The blackbody radiation field and Lyα pumping compete with collisional excitation to redistribute the level populations in Fe II and greatly affect the Fe II spectrum over the whole observed spectral range throughout the range of parameters $10^5$ cm$^{-3} \leq n_e \leq 10^6$ cm$^{-3}$ and $15,000$ K $\leq T_{BB} \leq 30,000$ K. We have studied the contribution of various excitation mechanisms (collisions, continuum pumping, and Lyα pumping) to the predicted Fe II spectrum. We have computed a series of models by disabling one of these excitation mechanisms. By doing so we explore the importance of this mechanism in explaining the richness of Fe II emissions.

Figure 14 depicts model A, assuming the following subsets of excitation conditions:

A1. Collisions and Lyα pumping.—This model does not include Fe II pumping by blackbody radiation (Fig. 14, upper panel). The [Fe II] $\lambda7156$ feature becomes strong for $n_e \geq 10^5$ cm$^{-3}$ and $T_{BB} \geq 15,000$ K. The predicted Fe II intensities are too weak compared to the complete model (see model A3). To explain the strength of Fe II emission, the iron abundance must be $50–80$ times larger than in the model with continuum pumping (model A2) and the model with both continuum and Lyα pumping (model A3). Moreover, lines from multiplets $a^6S_\Delta-6P_o$, $a^4G-z^4F_o$, $b^4D-z^4P_o$, $z^4D^c-e^4D$, and $a^6S-z^6D^o$ are not predicted at any density. Effects of Lyα pumping are so small that the strongest predicted lines in the near IR are due to collisions from the ground (the 17th lowest Fe II levels with energy up to 2 eV). Since the collisions at $n_e = 10^5–10^6$ cm$^{-3}$ do not sufficiently populate the upper levels of Fe II pumping channels, the effects of Lyα pumping are negligible. This model also does not reproduce the observed strong, rich spectrum in the 5000–6500 Å range.

A2. Collisions and continuum pumping.—This model does not include Fe II pumping by Lyα radiation. The strengths of Fe II lines are dramatically increased when blackbody radiation is included. Lines from $a^6S_\Delta-6P_o$, $a^4G-z^4F_o$, $b^4D-z^4P_o$, $z^4D^c-e^4D$, and $a^6S-z^6D^o$ multiplets are predicted. Comparison with model A1 shows that not only
and the distance to BD blobs from central star is 10,16 cm. See text for details of models A1, A2, and A3.

from a pumping helps explain the strength and appearance of lines permitted lines, but also low-lying forbidden transitions, are strongly affected by continuum radiation.

A3. Collisions, continua, and Lyα pumping.—This model includes all excitation processes. The presence of Lyα pumping helps explain the strength and appearance of lines from $a^4F - a^4P$ and $z^6D^--c^4P$ multiplets with upper levels below 1.7 eV and 6.2 eV, respectively. The complete model is able to predict only some Fe ii primary cascades due to Lyα pumping, the strongest being Fe ii λ9299 ($e^4D^-c^4P$). The other Fe ii intensities from primary cascades are at least 10 times smaller than seen in observations. We provide more details about Lyα pumping and show how our model can be improved in the next section.

A comparison of models A1, A2, and A3 shows that the continuum radiation field from the central source plays the dominant role in producing the strong, rich Fe ii spectrum in the 4000–6500 Å range. Its presence is also important for reproducing the observed strengths of predominantly collisionally excited and Lyα-pumped lines.

8. EFFECT OF Lyα PUMPING ON Fe ii EMISSION IN THE BD BLOBS

In this section we discuss the details of Lyα pumping fluorescence. Using existing $A$ values we analyze the branching ratios for pumping routes and primary cascades due to Lyα fluorescence. We compare the results from different models: model A (specifically A3) includes Lyα radiation based on predictions from the photoionization CLOUDY code, and model B assumes very intense Lyα radiation with the Lyα width of 10 Å. Model C is the same as model B, except with different atomic data.

8.1. Lyα Excitation Routes and Primary Cascades

Generally in Fe ii, the higher the energy level, the closer they are spaced in energy. Because of this close energy spacing, a number of upper levels can be populated via photoexcitation by Lyα emission. These levels then effectively produce cascades with strong emission lines observed in spectra. It is also possible that only one Fe ii level in an LS term gets pumped, which will cause a huge departure in intensities compared to branching ratios.

We have found that 24 Fe ii lines and blends in the ~8000–9600 Å range are primary cascades due to Lyα pumping in the BD blobs. However, only two others are strong in the UV, namely, Fe ii λ2507 and λ2509. Because of wavelength coincidence, the most “effective” term for Lyα pumping is the $a^4D^-$ term.

Table 2 presents details about the most important Fe ii pumping channels due to Lyα coincidence and the strongest Fe ii primary cascades. This information is based on radiative values used in model A3 and HST data. Table 2 has two main parts. The left side includes information about the main pumping routes and about the lower level from which electrons are radiatively pumped; the larger the $A$ value, the larger the radiative rate. Each electron that is pumped to the upper level will cascade according to the branching ratios. The strengths of the lines are in proportion to the branching ratios (if the upper level is the same). The right side of Table 2 provides information on the strongest primary cascades. Column (1) lists the Fe ii upper level in the transition pumped by Lyα fluorescence; column (2) the wavelength of Fe ii pumped channel, in Å; column (3) the Fe ii lower level in transition pumped by Lyα fluorescence; column (4) the transition probability, in s$^{-1}$, for the pumped transition; column (5) the wavelength of Fe ii primary cascade, in Å; column (6) the Fe ii lower level in primary cascade; column (7) the branching ratio (in decreasing order) for given upper Fe ii level; and column (8) gives notes. The Fe ii λ2507 and 2509 features exhibit unusual behavior that has been already discussed in previous studies (cf. Johansson & Letokhov 2001). Specifically, these two transitions have very small $A$ values for any pumping route or primary cascade, yet they are among the strongest observable Fe ii lines.

Since many Fe ii lines overlap, it is important to ascertain the strongest single Fe ii line pumped by Lyα. First, it will provide information about the wavelength of the most effective Fe ii pumping channel. Combined information about pumping channels provides an estimate of the width of Lyα. Second, its intensity can be used for reliable comparison...
between model and observations. Figure 15 shows the dereddened measured intensities versus pumped upper level from our model. Filled circles are used for intensities of the single pumped lines. Open circles are used for intensities of blended features. All primary cascades originate from levels that are in a range of \( \sim 11–11.3 \) eV. The strongest feature at \( \lambda 9204 \) is a blend of three lines pumped by \( \text{Ly}_\alpha \). The strongest measured intensity of the single line at \( \lambda 8491 \) is due to the primary cascade that corresponds to \( 322 \)th level (\( \sim 11.2 \) eV). The lower limit of \( \text{Ly}_\alpha \) at \( \sim 1212 \) Å and the upper limit of \( \text{Ly}_\alpha \) at \( \sim 1222 \) Å are determined to make pumping channels effectively work for all observed \( \text{Fe}^{\,\text{II}} \) primary cascades (see Table 1).

In the previous section we have mentioned that the complete model with \( \text{Ly}_\alpha \) pumping (model A3; Fig. 14) predicts \( \text{Fe}^{\,\text{II}} \) primary cascades, yet not all of them. To understand why, we have compared observable \( \text{Fe}^{\,\text{II}} \) primary cascades with that predicted from model A.

![Figure 15](image)

**Fig. 15.** — \( \text{Fe}^{\,\text{II}} \) upper levels that are pumped by \( \text{Ly}_\alpha \) excitation and their intensities due to primary cascades. Open circles indicate intensities for blended lines. Filled circles are for single lines. The strongest measured single line originates from \( 322 \)d level in the \( \text{Fe}^{\,\text{II}} \) ion. The strongest blended feature (near 9204 Å) is due to cascades from the 322, 338, and 307 levels.

### 8.2. \( \text{Ly}_\alpha \) Fluorescence, Model A

We have found (§ 7) that model A3, when \( \text{Ly}_\alpha \) radiation is predicted from CLOUDY calculations, does not reproduce the strongest observable \( \text{Fe}^{\,\text{II}} \) primary cascades. For a range of electron densities \( 10^5–10^7 \) cm\(^{-3}\), CLOUDY predicts \( \text{Ly}_\alpha \) much narrower (\( \sim 5 \) Å) and weaker than required to explain the \( \text{Fe}^{\,\text{II}} \) primary cascades in BD blobs. With such a narrow-line profile, the \( \text{Ly}_\alpha \) intensity in model vanishes longward to \( \sim 1218 \) Å. This wavelength corresponds to the most effective pumping channel for the peculiar 2507 and 2509 Å features, and observed 8453 Å line (see also Table 2). Since the \( \text{Ly}_\alpha \) flux falls off dramatically at \( \sim 1218 \) Å, the predicted number of \( \text{Fe}^{\,\text{II}} \) primary cascades is less than observed. Figure 14 (lower panel) shows the \( \text{Fe}^{\,\text{II}} \) intensities predicted by model A.

Model A uses \( \text{Ly}_\alpha \) intensity and width predicted for the local nebular conditions. We have to introduce the intense \( \text{Ly}_\alpha \) radiation (model B) to check whether the model with constant density \( 10^8 \) cm\(^{-3}\) is able to explain observed \( \text{Fe}^{\,\text{II}} \) primary cascades.

### 8.3. \( \text{Ly}_\alpha \) Fluorescence, Model B

Keeping all parameters in the model B the same as in model A, we have intentionally modified the shape of the \( \text{Ly}_\alpha \) intensity in the stellar radiation continuum. First, the total \( \text{Ly}_\alpha \) flux has been increased to reproduce the strongest \( \text{Fe}^{\,\text{II}} \) primary cascades. Second, for simplicity, we have used a \( \text{Ly}_\alpha \) rectangular shaped profile. Evaluating pumping channels and wavelengths of primary cascades, we estimate that the \( \text{Ly}_\alpha \) width must be at least \( \sim 10 \) Å. The width of \( \text{Ly}_\alpha \) was changed to provide observed intensities consistent with observations on primary cascades of the pumped \( \text{Fe}^{\,\text{II}} \) levels. Stellar atmosphere modeling done by J. Hillier (2001; and private communication) indicates that up to 50\% of total the radiation may come from \( \text{Ly}_\alpha \). Furthermore, Johansson & Letokhov (2001) have also suggested an intense \( \text{Ly}_\alpha \) radiation field to explain the strength of the peculiar features at 2507 and 2509 Å.
Model B predicts the strongest observable line for each multiplet according to used branching ratios (e.g., 8453 Å line in Table 2). There are a few cases in which the atomic data simply are not in agreement with observations. The Fe II λλ 2509 and 2509 lines are most noteworthy (Table 2). The A value of 2509 Å feature is small in our original model. It is not clear how the line could be strong even if its branching ratio is the largest from this level. On the other hand, the upper level of 2507 Å line should be strongly pumped by Lyα due to coincidence because of a large A value. But the corresponding branching rate is too small to explain the strength of the observed line.

Figure 16 shows model B. The main difference between model A and model B is the presence of strong Fe II primary cascades pumped by Lyα. The presence of intense Lyα line (model B) does not explain the strength of famous 2507 and 2509 Å features. In model B they are too weak at any of the physical parameters: Lyα intensity, density, strength of radiation field, and $T_{BB}$. It is important to explain the origin of these lines because they are observed in other objects (e.g., in RR Telescopii; see Hartman & Johansson 2000).

8.4. Lyα Fluorescence, Model C

Because of a high level density of upper Fe II states, the level mixing causes changes in the A values. It is difficult to predict A values theoretically. In an attempt to improve the A values, we have changed them from theoretical values (Nahar 1995) to the A values of the Kurucz database (1983), where experimental A values have been inserted when available.

To solve the problem of the 2507 and 2509 Å features, we have implemented mixing for their upper levels ($^4F_{2/2}$ and $^4G_{9/2}$), which are very close in energy ($\approx 25$ cm$^{-1}$). We have adopted 50% for the mixing of these levels, as suggested by Johansson (1984). Two upper levels are treated as a combination of two different states. They are pumped with 50% of the most effective route and cascade with 50% of the most effective primary cascade.

Model C successfully predicts Fe II 2507 and 2509 Å lines as the strongest lines in the whole spectrum. However, a number of lines in near IR from 8000 Å to 10000 Å are larger than observed (Fig. 17).

8.5. Comments on Lyα Fluorescence Modeling

The HST observations of η Carinae nebula show a broad absorption in Lyα. The interstellar circumstellar absorption completely scatters any intrinsic stellar emission. Thus, there is no information on the intrinsic stellar Lyα emission profile. The simple photoionization model that calculates the nebular Lyα radiation does not predict the intense Lyα radiation that is needed to explain the presence of strong pumped Fe II lines. It is possible that Lyα line has an origin other than just the nebula condition. Thus, there may be a strong stellar Lyα emission. Indeed, the strong stellar η Cyg Balmer lines with widths of ±1000 km s$^{-1}$ imply a corresponding strong intrinsic stellar Lyα emission. In addition, η Carinae is known as a massive star with strong wind a mass-loss rate of $\approx 10^{-5} M_\odot$ yr$^{-1}$ (Davidson et al. 1995; Hillier et al. 2001). The wind speed of $\geq 300$ km s$^{-1}$ can easily produce a η Cyg profile with a strong redshifted emission of the Lyα profile, and could explain the effectiveness of Lyα pumping at $\approx 1218$ Å.

We have changed the Lyα intensity from model A to model B to reproduce the intensities of observed primary cascades. There are three main sources of uncertainties for the Lyα pumping models: the strength of Lyα, the shape of the Lyα profile, and the A values of Fe II transitions. We have demonstrated that model predictions depend on the accuracy of the A values, which determine the radiative rates. All pumped channels use theoretical transition probabilities. Experimental or better theoretical A values will most likely improve model predictions.

The major factors for the efficiency of the pumping are the Lyα intensity, the Fe II transitions coincidence with Lyα wavelength, and the Fe II transition probabilities. Pumping by a blackbody radiation field works more effectively than Lyα fluorescence in BD blobs. Only 24 strong features from more than 200 Fe II lines and blends are due to Lyα fluorescence. Even in the case of strong Lyα radiation, the effects of continuum pumping are more important for the shape of the Fe II emission spectrum.

If the Lyα radiation comes from spatially extended region of the disc, the total Lyα pumping effects will be much stronger than that produced by Lyα of nebula or central source origin. The increased Lyα intensity is modeled by increasing the number of photons in order to reproduce the strength of
observed Fe II lines. The total contribution of Lyα pumping to the origin of Fe II emission is weaker than that of the blackbody radiation.

9. IMPACT OF ATOMIC DATA ACCURACY ON THE Fe II EMISSION MODELS

The modeling of the Fe⁺ ion requires an enormous amount of atomic data. In particular, energy levels, $A$ values, and collisional excitation rates are needed to predict the Fe II emission.

The atomic data are not always known with the required accuracy, and in some cases there are no data available. These limitations affect our models. All wavelengths are based on experimental data, whereas all excitation rates are theoretical, as are the $A$ values for forbidden lines. The $A$ values given for the allowed transitions vary depending on what database is used. In the IRON project (Nahar 1995), all $A$ values are theoretical, but in the Kurucz database (1981), experimental $A$ values are inserted where available.

The main discrepancies between our model and observations that can be attributed to atomic data are due to the following reasons:

1. $A$ values are not known at all for many semiforbidden transitions. For these we have inserted reasonable estimates.

2. Theoretical $A$ values for allowed transitions are calculated in the $LS$-coupling approximation (Nahar 1995). The deviation from $LS$ coupling is stronger for the upper Fe II levels (Johansson 1984). The close energy spacing of the upper levels leads to their mixing.

3. Collisional excitation rates are calculated only for transitions between the lowest 64 levels (Bautista, Peng, & Pradhan 1996). For other observed lines we used the less accurate $g$-bar approximation.

As an example of strong mixing, we discuss the famous Fe II λλ2507, 2509. They originate from upper levels that are very close to each other at 90067.347 (6$F_{9/2}$) and 90042.779 (4$G_{9/2}$) cm$^{-1}$, and they both cascade to the same lower level at 50187.813 cm$^{-1}$ (6$F_{7/2}$). The values for pumping and primary cascades derived from $LS$-coupling calculations are so small (Table 2) that when they are used in calculations, neither line is predicted to be strong. When we add mixing ($\xi$), Fe II λλ2507 and 2509 lines become the strongest lines in the whole Fe II spectrum, a factor of 3 stronger than other Fe II lines.

We have performed calculations for Lyα pumping using the different atomic data sets ($\xi$). Model B includes Nahar (1995) data, and model C includes Kurucz (1981) data for the important pumping routes and primary cascades due to Lyα. The $A$ values differ by up to a factor of 100. Comparison of models B and C (Figs. 17 and 18) demonstrates the importance of accurate $A$ values for the model. The Kurucz data used for model C predict the presence of observed pumped levels better than calculations with Nahar’s data. In addition, the model spectrum based upon the Kurucz data set predicts more Fe II emission lines between 9000 and 9700 Å pumped by Lyα, which is more consistent with observations. This suggests that the accuracy of atomic data must be improved.

The narrow and strong Fe II lines of the BD blobs allow us to compare in great detail model predictions with observations. Further calculations and/or measurements for the unknown $A$ values should improve model predictions.
The BD blob spectrum of η Carinae is uniquely rich in Fe II and [Fe II] lines. The Fe II studies reveal details of physical conditions in the emitting region, and indicate shortcomings in the atomic data for Fe II.

We conclude that Fe II and [Fe II] lines in the BD blob spectrum are due to the dominant excitation mechanisms (Fig. 18):

1. Pumping by the blackbody-like stellar radiation field from the central source of η Carinae produces the strong and rich spectrum in the 4000–6500 Å range. Most of observable [Fe II] and Fe II lines have accurate A values (Quinet et al. 1996). These lines are in a good agreement with model predictions.

2. Primary cascades after Lyα fluorescence are responsible for 24 strong Fe II lines in the 8000–10000 Å range and in particular Fe II λλ2507, 2509 in the BD blobs. Secondary cascades have much weaker effects. However, they contribute to the strengths of many other Fe II emission lines. We have postulated a large Lyα flux along with a broad emission profile to explain the observed Fe II primary cascades. In our future work we will study Lyα pumping in greater detail. The Fe II level mixing, the number of Fe II channels involved in excitation, the Lyα intensity, and the profile width and shape are the most important unknown factors.

3. Collisional excitation dominates in the formation of strong [Fe II] lines in the 7000–9000 Å range. For n_e = 10^10–10^11 cm^-3, the lowest levels of Fe II can be populated by collisions. The strength of the [Fe II] λ7156 is systematically underpredicted for a constant-density model. It is not clear if its intensity is due to lower density component (B or D), uncertainties in atomic data, or geometric effects. Small collisional strengths or an overlooked excitation route might also affect the predicted intensity.

We have used the Fe II spectrum to derive physical conditions in the BD blobs. The best fit to the observed spectrum is achieved in the 4000–6000 Å range at constant density n_e ~ 10^6 cm^-3 and T_BB ~ 15,000 K. The dominant Fe II ionization fraction and rich Fe II spectrum suggest the electron temperature range of T_e = 5000–7500 K. Our model predicts that the Fe II/Fe^+ ion has very sharp transition due to collisional ionization at ~10^6 cm^-3 and the size of blob ~10^15 cm^-3. The Fe II ionization structure, low electron temperature in plasma, the shape of blackbody radiation, and intense Lyα radiation provide conditions favorable for a rich Fe II spectrum. Blackbody radiation and Lyα pumping are the main excitation mechanisms for the Fe II ion in the BD blobs.

As our first step, we have focused on Fe II emission analysis of the spectrum of the BD blobs from 1998 observations. Extensive HST/STIS η Carinae observations at different positions and during other epochs are available. The morphology of the Homunculus Nebula observations at different positions and during other epochs are available. The morphology of the Homunculus Nebula observed in η Carinae is complex. Line profile analysis by Vioiti et al. (1989) and Davidson et al. (1995) suggests a wider range of velocities than expected from spherical geometry. In addition, an inhomogeneous dust distribution may selectively obscure the radiation field. As a result, any of the combinations considered above (e.g., collisions, blackbody radiation) may be responsible for the Fe II and [Fe II] emission spectrum at other than BD positions.

For the first time, we have used detailed non-LTE Fe II models to bracket physical parameters of Fe II-rich plasma in the BD blobs of η Carinae. Most of the [Fe II] and Fe II lines predicted by our model are in a good agreement with observations.

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