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Fast 4D imaging of fluid flow in rock by high-speed neutron tomography

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¹¹ Key Points:

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12	•	Full 3D fluid front velocity map can be obtained from high-speed (1 minute/tomography)
13		neutron tomography during water invasion into air-filled samples
14	•	Quantitative measurements are validated by comparing experimental results to
15		1D analytical solution of pressure-driven flow
16	•	During imbibition, compactant shear bands result in locally higher fluid-flow ve-
17		locity due to decreased pore size in the bands

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18 Abstract

High-speed neutron tomographies (one minute acquisition) have been acquired during 19 water invasion into air-filled samples of both intact and deformed (ex-situ) Vosges sand-20 stone. 3D volume images have been processed to detect and track the evolution of the 21 waterfront and to calculate full-field measurement of its speed of advance. The flow pro-22 cess correlates well with known rock properties, and is especially sensitive to the distri-23 bution of the altered properties associated with observed localised deformation, which 24 is independently characterised by Digital Volume Correlation (DVC) of x-ray tomogra-25 phies acquired before and after the mechanical test. The successful results presented herein 26 open the possibility of in-situ analysis of the local evolution of hydraulic properties of 27 rocks due to mechanical deformation. 28

²⁹ 1 Introduction

Deformation in rocks is typically not homogeneous and is often localised into fea-30 tures such as shear or compaction bands. Such deformation can have significant influ-31 ence on the hydraulic properties of the rock because it locally alters the rock structure 32 in ways that are expected to change the local flow properties. Standard experimental 33 approaches to assess deformation and changes in fluid flow in rocks are based on mea-34 surements of fluxes and pressures taken at the boundaries of a tested sample. These mea-35 surements, however, cannot provide information on internal heterogeneities, which likely 36 37 exert a strong control on the mechanical and fluid flow behaviour. Therefore, understanding how rock deformation influences the hydraulic properties of rocks requires local ob-38 servations that can provide a direct link between the mechanical deformation-induced 30 changes and the altered fluid flow responses. This necessitates 3D observations of flow 40 and measurements over time, since the fluid flow is a dynamic process that can only be 41 assessed as it occurs, and deformation features can have a complex 3D arrangement. Thus, 42 time-resolved 3D (i.e., 4D) measurements and observations are required. 43

X-ray imaging has become a common method to characterize the internal struc-44 ture of bulk objects (e.g., Maire and Withers (2014)), including geomaterials, such as 45 rocks and soils (e.g., Cnudde and Boone (2013)). Such 3D imaging has been extended 46 to look at the evolution of localised deformation in rocks using time-lapse imaging and 47 Digital Volume Correlation (DVC) (e.g., Charalampidou et al. (2013)). Water imbibi-48 tion in rock specimens with compaction bands has also been studied using medical x-49 ray tomography (David et al., 2008; Pons et al., 2011) and pore-scale fluid flow processes 50 have been imaged in small rock samples using 4D x-ray micro-tomography (Berg et al., 51 2014; Youssef et al., 2013; Pak et al., 2015). X-rays, however, have a strong interaction 52 with the rock material and a much smaller interaction with the pore fluids of interest. 53 Neutron imaging provides an alternative method that is much better adapted to the imag-54 ing of hydrous fluids in rocks. Neutrons have a different interaction mechanism with mat-55 ter to x-rays. In particular, neutrons are absorbed or scattered by the nucleus of the atoms 56 while x-rays interact with the electrons shells. Therefore, x-rays absorption is mainly pro-57 portional to material density whilst neutrons absorption is linked to the capacity of the 58 nucleus to host an additional neutron. This leads to a strong interaction of neutrons with 59 hydrogen, whose nucleus is formed by only a proton and it is keen to receive a neutron, 60 and thus with the hydrogen-rich fluids of interest in rocks. In contrast, deuterium does 61 not have such a capture site (it is already filled), and thus the D_2O form of water is sub-62 stantially transparent to neutrons, much like the solid components of the rock and the 63 metals used in the experimental configuration. (e.g., Perfect et al. (2014); Kaestner et 64 al. (2016)). Therefore, much higher contrasts of fluid, compared to the rock, can be achieved 65 with neutron imaging than with x-ray imaging, which enables much lower saturations 66 to be detected. Furthermore, with a higher contrast, a sufficient signal to noise ratio can 67 be obtained with a shorter acquisition time leading, potentially, to faster imaging. 68

Previous work (e.g., Masschaele, Dierick, Cnudde, et al. (2004), Masschaele, Di-69 erick, Van Hoorebeke, et al. (2004), Hall (2013), Tötzke et al. (2017), Tudisco, Hall, et 70 al. (2017)) demonstrated the potential of neutron imaging to follow internal fluid flow 71 in geomaterials. In Masschaele, Dierick, Cnudde, et al. (2004) and Masschaele, Dierick, 72 Van Hoorebeke, et al. (2004) an example of 3D visualisation of fluid advanced is presented. 73 However, the dimensions of the sample and the resolution were limited and no informa-74 tion on acquisition time is provided. The authors express the need for a quantitative anal-75 ysis of the water front motion. In deformed rock samples, this analysis was restricted to 76 2D (radiography) imaging due to limitations with the possible acquisition speed. Here, 77 the restriction to only 2D imaging of dynamic flow processes is overcome with advances 78 in high-speed neutron tomography, to provide full 3D observations of flow processes that 79 operate over a finite time interval. These measurements are analysed to provide full-field 80 analysis of fluid–front velocity distributions, which is linked to hydraulic conductivity 81 and the boundary conditions, for; (i) an intact rock sample subjected to pressure-driven 82 flow; (ii) a sample exhibiting localised deformation subjected to a water-imbibition pro-83 cess. In the second case, the fluid flow analysis is compared with 3D measures of defor-84 mation from DVC. 85

⁸⁶ 2 Experimental approach

The neutron tomographies employed in this work were acquired with a high-speed 87 imaging set-up at the CONRADII beamline at Helmholtz Zentrum Berlin (HZB) (Kardjilov 88 et al., 2011), as described in the following. High-speed neutron imaging (1 minute per 89 tomography) is possible at CONRADII using a white beam at the experimental posi-90 tion closest to the neutron guide, where the neutron flux is maximum (i.e., 2×10^8 n/cm²), 91 which enables the exposure time to be minimized. However, at this position, the distance, 92 L, between the pinhole aperture (diameter, D), which defines the source, and the sam-93 ple is small. Since the resolution for neutron imaging is defined by the ratio L/D, this 94 implies a reduced spatial resolution. For the experiments presented here, a pinhole di-95 ameter of 30 mm was used, which gives an L/D ratio of 167. In addition to the L/D ra-96 tio, the distance, l, between the sample and the detector also influences the image res-97 olution according to the equation $d = l \cdot D/L$, where d is the maximum blur (e.g., Banhart 98 (2008)). Therefore, samples should be placed as close as possible to the detector; here, 99 due to the sample diameter, this distance was 20 mm, giving a theoretical resolution of 100 120 µm. 101

The neutron tomography acquisitions involved 300 radiographic projections acquired 102 during rotation of the sample over 180° . Acquisitions were made with consecutive, con-103 tinuous positive and negative rotations over 180° , as full rotation over 360° was not pos-104 sible due to the possibility of collision between the detector and the water inlet tube (see 105 Fig. 1). To minimise the exposure time for each projection the camera was used with 106 a pixel binning of 4x4 pixels. This allowed fast exposures (0.2 s per projection) with suf-107 ficient signal:noise, which, in turn, permitted continuous rotation, without deadtime for 108 stopping for each projection. With these settings, the full tomography acquisition of 300 109 projections could be acquired in just 1 minute and with a pixel size of 100 μ m over a field 110 of view (FoV) of about $60x73 \text{ mm}^2$, which is slightly smaller than the height of the sam-111 ples. (Note that, due to the high attenuation by the silicone used to seal the base of the 112 sample cup, the useable FoV was slightly smaller vertically). 113

Two samples have been studied in this work to illustrate the experimental approach. These were both 38 mm diameter cylindrical samples of arkosic Vosges sandstone (22% porosity) selected from the materials used in Charalampidou et al. (2011, 2013). One sample (IGSN: IEETV003) was "as-cored" and the other (IGSN: IEETV002) had been loaded under triaxial conditions at 30 MPa confining pressure such that it contained shear bands. A notch was cut at about 2/3 of the height of the sample and covering 1/4 of the circumference to encourage the expected localised deformation to occur in the middle



Figure 1. Photograph of the experimental set-up and a sketch of the inside of the cup. The water reservoir is to the left and the sample, mounted on the rotation stage, is to the right, in front of the scintillator of the detector. The electro-valve is visible below where the tube exits the reservoir. The metal post below the sample is attached to the rotation stage (out of sight, below).

region of the sample, similarly to Charalampidou et al. (2011). This deformed sample 121 had been imaged before and after the triaxial testing using both x-ray and neutron to-122 mography. DVC was performed on the x-ray images using the code TomoWarp2 (Tudisco, 123 Andò, et al., 2017) to provide 3D volumetric and shear strain maps throughout the sam-124 ple that are used here to compare to the flow measurements. This analysis indicated that 125 two main shear bands had developed extending from the notch, one reaching the top of 126 the sample and the second extending diagonally downward for about 25 mm. Additional, 127 mainly smaller (shorter), bands occur in complex patterns that are associated with the 128 larger bands, or they may be located away from it. Although the orientations of these 129 smaller bands are mainly similar to the main bands, there are differences that lead to 130 a reluctance to describe the entire array as having variations only in two axes. 131

The intact sample had parts of the cylindrical sides removed to create two, diametrically-132 opposite flat faces on which two notches had been cut on the surfaces at different heights. 133 This particular shape enables further advanced measurements (i.e., ultrasonic tomog-134 raphy) to fully characterise the sample before and after deformation. For the neutron 135 experiments (to ensure a good fluid sealing and optimal imaging), specially machined 136 teflon inserts were placed on the flat surfaces to recover a cylindrical shape. Teflon tape 137 was wrapped around the sample before the inserts were positioned to prevent fluid from 138 being able to flow out of the sides of the sample. The sample, plus teflon inserts, was 139 confined in a heat-shrunk Fluorinated Ethylene Propylene (FEP) membrane to seal the 140 ensemble. The complete sample assembly was mounted and sealed with silicone onto a 141 specially-designed end-cup, which allows water to contact the entire sample base. This 142 cup was connected by flexible tubing to a small reservoir on a table whose height could 143 be remotely adjusted (see Figure 1). An electro-valve, controlled from outside the ex-144 perimental area, allowed the fluid supply from the reservoir to the end-cup to be turned 145 on or off. 146

During the first test, on the intact sample, the level of the water was kept at a constant height, above the top of the sample, by moving the reservoir upwards to compensate for the fluid leaving the reservoir and entering the sample. In this way, an almostconstant water pressure was applied at the bottom of the sample. For the initial stage of the second test (sample with shear bands), the level of the water was kept constant and level with the top of the sample to accelerate the advance of the fluid until the wa-



Figure 2. (a) Two vertical cross-section slices through the high-speed neutron tomography images of the intact sample at different times in the fluid flow process. (b) The same slices reconstructed after subtracting the initial (dry) images, as described in the text. (c) Corresponding slices through the binarisation of (b). Note that the first (dry) tomography was made when water was already in the field of view, so the region of the sample where water was already present appears as a dark area at the bottom of the image in (b) and (c).

ter front reached about 1/3 of the sample height. Subsequently, the reservoir was low-153 ered to maintain the water level a few mm above the sample base, which allowed imbibition-154 dominated flow to be studied in the area of the sample that presented localised defor-155 mation (middle-top part). The flow experiments lasted for about 2.5 hours for the in-156 tact sample and 5.5 hours for the deformed one, which resulted in 149 and 384 tomo-157 graphies, respectively. A technical problem occurred in the middle of the second test such 158 that no images were acquired for about 40 minutes, which resulted in a gap in the imag-159 ing of the imbibition process. 160

¹⁶¹ 3 Image analysis

The tomographic reconstructions were performed using an in-house Python code 162 based on the ASTRA tomography toolbox (van Aarle et al., 2015; W. J. Palenstijn et 163 al., 2017) assuming a parallel beam and using the 3D SIRT GPU-based algorithm (W. Palen-164 stijn et al., 2011). Figure 2(a) presents two vertical slices extracted from the centre 3D 165 images of the intact sample at two different times. These images clearly show the ad-166 vance of the fluid front and the imaged texture of the rock in the dry region above the 167 advancing front. In the water-filled area, the neutron attenuation is high and the tex-168 tural information is obscured. Moreover, a beam hardening effect is evident in the sat-169 urated region. To facilitate the thresholding of the images described above, the tomo-170 graphic reconstruction was repeated using radiographic projections normalised with the 171 corresponding projections from the initial tomographic scan of the dry sample. This pro-172 cedure allowed only the water to be reconstructed, as shown in Figure 2(b). It can be 173 seen that this process removes the rock sample texture from the images, which provides 174 further support that the sample texture visible in Figure 2(a) is real and not noise. 175

The focus of these experiments is to demonstrate the capability to achieve a quantification of the advance of the fluid front, which is clearly visible in the images in Figure 2. The contrast between the fluid-filled and the dry regions of the samples is much greater than the noise level or the effect of beam hardening, especially after removing



Figure 3. Speed field determination steps: binary image of (a) the initial front (Front 1) and (b) the manually-adjusted next front (Front 2), (c) negative of Front 2, distance map from (d) Front 1 and (e) from Front 2 masked to show only the area of interest, (f) sum of the distance maps.

the tomographic images of the dry sample (Figure 2(b)). By binarisation of the 3D images it is possible to separate the dry and wet volumes to provide a series of binary 3D images that represent a time series of the water front progression. From each consecutive pair of binary image series, a 3D map of the speed of advance of the waterfront has been calculated, as described in the following.

Figure 3 illustrates the speed map determination approach for an artificial 2D case 185 with a front manually moved up 20 pixels in the vertical direction and flattened on one 186 side so that the maximum displacement there is 48 pixels. The first step in the analy-187 sis procedure involves calculation of the euclidean distance of each pixel in the image from 188 the nearest black pixel to give distance maps. This is performed for the first binary im-189 age in each consecutive image pair and is repeated for the second image in the pair, but 190 with a negative of the binarised image to retrieve the distance in the correct direction 191 from the second front (i.e., downward). To reduce artefacts at the boundaries of the sam-192 ple, the outside of the sample in the images was set to zero (see Figure 3(c)). For the 193 real data, the shape of the sample is determined from the last tomography image, when 194 the sample is full of water. The two distance maps are subsequently masked to retain 195 only the image of the sample region between the two fronts. These masked distance maps 196 are summed together to provide a new image that represents, for each point in the im-197 age, the shortest-path distance that water crossing the point travels to reach the second 198 front from the first. Knowing the time elapsed between the acquisition of the two im-199 ages, this distance can be directly converted to a speed of the front advance for each point 200 between the fronts in the pair of images. Figure 3(f) shows that the calculated distance, 201 in the artifical example, is constant and equal to 20 px where the front has moved uni-202 formly, whereas the distance increases where the front is distorted. Curvature of the front 203 at the boundaries results in an inclined path being shorter than the vertical one. Whilst 204 this does not correctly represent the imposed constant advance of the front and indicates 205 caution should be applied when interpreting the boundary areas, this might be consid-206 ered as a reasonable representation of a real flow where there is no reason for the wa-207 ter to advance only in the vertical direction. 208

The described procedure provides the "instantaneous" flow speed field of the advance of the water-front between two tomographies. Repeating the process for each consecutive pair of images and assembling the results provides a full speed map volume covering the entire imaged sample. The resolution of the speed map is determined by the frequency of the front reconstructions (each tomography) and it is limited by the noise, which affects the smoothness of the surface. To obtain reasonable results, the movement between two fronts has to be larger than the front roughness.



Figure 4. (a) A series of waterfronts in the central vertical slice, (b) []correspondingvertical slice of the water speed map for the intact sample, and (c) an analytical solution, as in equation (1), fitted with experimental data.

216 4 Results

Figure 4 shows, for the intact sample, the positions of the water-fronts (at 2 minute 217 intervals) in a vertical central slice through the imaged volume and the corresponding 218 slice through the 3D speed map (the former determined by binarisation of the tomog-219 raphy images during the water advance, and the latter calculated using the method de-220 scribed above). The speed map shows a decrease in the speed of the front advance with 221 increasing height of the front. This is consistent with a decreasing pressure gradient with 222 increased distance of the front from the constant-pressure condition at the sample base. 223 Furthermore, the flow-speed field does not show any significant local variations, which 224 is as expected from the assumption of a nominally intact and quasi-homogeneous sam-225 ple. Moreover, closer inspection of the fluid front in the tomography images indicates 226 there is no significant gradient in saturation at the front. Therefore, the role of capil-227 larity effects can be neglected in the analysis: i.e., the flow is predominantly pressure driven. 228

The second test differs from the one described above in two respects: it contains 229 an array of experimentally-created deformation bands; and the flow experiment is con-230 ducted so that, in the top part of the sample, imbibition of water is the only operative 231 driving force. Figure 5 shows horizontal and vertical slices through: (a) the neutron to-232 mography image at the end of the flow experiment; (b) the water-front speed volume; 233 (c-d) the maximum shear- and volumetric-strain volumes derived from DVC applied to 234 x-ray tomographies acquired before and after triaxial deformation. Strain localisation 235 bands, and their related flow effects, are clearly visible in all the presented images, in-236 cluding the smaller scale features of the localised deformation in the upper part of the 237 sample. Figure 6 presents the 3D rendering of the water speed map and the maximum 238 shear strain, in which low values are displayed as transparent to allow the visualisation 239 of the deformation band. The images show the variability of the deformation along the 240 sample in both shape and intensity. Such variability would be lost in a 2D analysis, where 241 only an average projected value in the plane of the deformation can be measured. 242

The higher attenuation of neutrons seen in the region of the deformation bands in Figure 5(a) indicates higher water saturation, and the water-front speed map (Figure 5(b)) shows a higher value inside the bands. This experiment was conducted so that imbibition of water is the only operative driving force, therefore the higher speed in the localised deformation bands can be interpreted as being related to higher air-water cap-



Figure 5. horizontal and vertical slices through: (a) the neutron tomography image at the end of the flow experiment; (b) the water speed volume; (c) the maximum shear-strain volume; (d) volumetric-strain volume (positive values indicate compression). The strain fields were derived from DVC applied to x-ray tomographies acquired before and after triaxial deformation. To better highlight the variations in the water speed map, a Gaussian filter with radius 5 pixels was applied twice to enhance the signal:noise) this also generates the dark borders visible in (b)).



Figure 6. 3D rendering of (a) the water speed map and (b) the maximum shear strain. Low values are displayed as transparent to allow the visualisation of the deformation band.

illary pressure, which suggests smaller pore sizes in the band than in the surrounding 248 region. This hypothesis is consistent with the strain maps from the DVC, which reveal 249 that the shear bands are predominantly compactant (although some regions of dilation 250 can also be seen, especially close to the boundary of the sample). The shear band ex-251 tending downwards from the notch shows higher water-front speeds, which suggests that 252 the capillary properties in this part of the band array are strongly affected by deforma-253 tion, even though this localisation zone is less visible in the DVC results. In contrast, 254 the band continuing from the notch towards the top of the sample appears to have a lesser 255 impact on the water speed, compared to the surrounding region. This could be due to 256 the much lower velocity of the water at this height (which can only advance as quickly 257 as more water can arrive from below), which would make the difference between the ve-258 locity in the band and in the surrounding area too small to be appreciated. Another pos-259 sible explanation is the dilatant character of this part of the shear band, which might 260 reduce the capillary effects. In the central height region, where no information is avail-261 able on flow because of the break in the image acquisition described above, it is inter-262 esting to note that the method is still able to capture the average speed and the higher 263 speed in the bands despite the long distance between the two water-fronts. 264

An additional neutron experiment is briefly described here to emphasise the rel-265 evance of the three-dimensional nature of tomography. This experiment is performed on 266 a carbonate (limestone) rock comprised of very thin layers (called laminae, <1mm thick) 267 oriented orthogonal to the cylinder axis. These laminae are composed of aggregates of 268 micro-crystals (2-5 μ m diameter), with alternating laminae exhibiting variations in inter-269 crystalline pore space such that local porosity ranges from <15% to >45% (Buckman 270 et al., 2018). The sample was pre-deformed (dry) to just past the peak stress at a con-271 fining pressure of 20 MPa. Instead of developing through-going deformation bands like the samples described above, it developed a complex array of shorter, planar and pre-273 dominantly steep features (Figure 7(a)) that exhibit a dilational (lower density) char-274 acter on x-ray tomography examination (Figure 7(d)). On subsequent inspection by SEM 275 (Scanning Electron Microscopy), they are seen to include a dilation-dominant mixture 276 but with some compactional zones (Figure 7(b),(c)). Nominally, at macro-scale, these 277 features look like partly-open fractures, with shear movement indications as well as open-278 ings like those seen at smaller scale in the SEM images. They are typically steeply in-279 clined with multiple orientations arrayed around the sample axis, and altogether, they 280 create a complex network of partially-intersecting discontinuities at a wide variety of ori-281 entations that cannot be captured in a 2D representation (Figure 7(e), (f)). They also 282 exhibit a clear relationship with the laminae, with terminations typically but not exclu-283 sively occurring at lamina boundaries. 284

The neutron experiment on this sample (performed at the Institut Laue-Langevin 285 (ILL), using the then newly-commissioned NeXT facility, which at the time could not 286 acquire high-speed neutron tomographic images due to constraints, since overcome, with 287 the control software for the rotation stage) reveals considerable complexity in the flow 288 patterns. The initially air-filled sample was subjected to slow water injection at its base, 289 using a cup analogous the one described above. The invasion of the initial fluid (D_2O) 290 was extremely non-uniform. The following description is based on neutron radiography 291 images (Figure 8), so, while the observations reveal clear patterns of flow behaviour, the 292 fixed viewpoint inevitably results in significant uncertainty except when the view is aligned 293 directly along a planar feature. In particular it is impossible to identify if a tabular zone, 294 or an irregular zone with finite length and width, as seen on the radiography is actually 295 water in porous rock or if it is a fracture oriented normal to the viewing direction. 296

Initially, the water moved up along a few of the fractures, readily identifiable in the radiography as such (Figure 8(a)) because of their suitable orientation. These fractures were known to intersect the base. As well as moving up the fractures the D_2O moved horizontally creating a set of typically approximately tabular zones tabular zones. These



Figure 7. (a) Photo of deformed sample, showing laminated character of this rock, and the expression of deformation features as they intersect the sample exterior. Sample is 38mm diameter. (b) SEM image of a small (~2mm field of view) part of the deformed sample. Image courtesy of Jim Buckman. (c) High-resolution SEM image of part of one fracture, showing the comminution of grains forming a groundmass of ultra-fine particles with nano-porosity. Note also the presence of disconnected microfractures within the band. Field of view ~20 microns. Image courtesy of Jim Buckman. (d) Vertical slice from post-deformation x-ray tomography scan, showing that the fractures are dark (less dense). Note how some of the deformation features terminate against a depositional interface, and others transect larger regions. (e) Radiograph of the sample after significant water has entered the sample. Note how the fractures are dark, indicating higher water saturations. (f) Part of a 45-minute tomography scan obtained after H₂O replaces D₂O. Note how the H₂O is mainly located within the deformation features, but has invaded the unusual set of more porous laminations in the middle of the sample (compare also with (a) and (d)).



Figure 8. Sequential false colour radiographic images of the laminite sample.

semi-tabular zones show water ingress into laminae (Figure 8(b),(c)) but it was not pos-301 sible to determine from the radiography if these zones were actually indicating D_2O oc-302 cupying the porespace of the undeformed rock or if it was occupying a fracture oriented 303 at right angles to the direction of view. But the tomographic images (Figure 7(f)) show no open fractures in this position. Significant D_2O saturation was normally achieved at 305 least the right side of the sample as viewed in Figure 8 before the process moved to a 306 progressively higher lamina and repeated. Generally filling moved from right to left, which 307 was also away from the well-connected fracture network horizontally along the laminae, 308 as well as from the base upwards. Some laminae were not invaded, while a few laminae 309 were filled by cross-layer movement from an adjacent or nearby lower lamina: again this 310 discrimination between lamina and fracture D_2O ingress required the tomographic im-311 ages. Some suitably oriented fractures can be seen in the radiography to start partic-312 ipating in the flow, and delivered D_2O to a higher laminae, with the D_2O then moving 313 downwards into a dry lamina, typically along a fracture (compare Figure 8(c) and (d)): 314 again the unsuitably oriented fractures needed the tomographic image to confirm their 315 identity. This filling process was modified as the height of the D_2O invasion increased, 316 first encountering a zone of more porous laminae, with no suitably positioned open frac-317 tures, which filled (probably radially) lamina by lamina. The early stages of this pro-318 cess are seen in Figure 8(d). Above this region of more porous laminae, the sample has 319 fewer but longer and less-well-connected fractures (Figure 7(f)). However the same pat-320 tern of water ingress along fractures and laminae was observed and confirmed by the com-321 bination of radiography and tomographic images as needed. When the invasion achieved 322 a quasi-static state (i.e., no discernible changes in saturation), the injection fluid was switched 323 to H_2O . Due to higher neutron attenuation of H_2O , the flow pattern in this semi-saturated 324 condition was readily apparent (Figure 7(c). During this phase of the experiment, the 325 flow was almost exclusively concentrated within the fractures, and water exited the top 326 of the sample very quickly. 327

Despite the high neutron flux $(8.6 \times 10^7 \text{ n/cm}^2/\text{s})$, the initial setup on the NeXT 328 facility (at the time this sample was used in the experiment) did not permit the rota-329 tion stage to operate at high-speed, so slower, higher quality tomographies (typically 45 330 minutes) were acquired. Fluid injection was stopped while individual tomographies were 331 acquired, with a limited degree of diffusion taking place during the suspended flow. For 332 the majority of the experiment, the process was only followed through radiography (with 333 an exposure time of $0.6 \, \mathrm{s}$). However, at the more interesting points, the sample was re-334 motely rotated to allow the experimental team to construct a mental picture of the 3D 335 character of the flow being observed. For both tomography and radiography a pinhole 336 of 20 mm was used, which gave an L/D of 500. The high-speed neutron tomographic method 337 described here is now operational at the ILL, and achieves tomography acquisition in less 338 than 1 minute, thanks to the uniquely high neutron flux. Subsequent experiments, us-339 ing samples similar to this one, along with other carbonate rocks, have since been un-340 dertaken, and the full analysis of that suite of experiments will be reported separately. 341 The key point, in relation to this paper, is that flow processes is intrinsically three-dimensional 342 and can exhibit major complexity depending on the nature of the rock and its loading 343 history. High-speed neutron tomography is in this respect an essential tool which enables 344 new investigations of fluid flow in complex geomaterials. 345

346 5 Discussion

The results presented in the previous section demonstrate the ability to capture the time-dependent evolution of the fluid flow front in 3D via high-speed neutron tomography, both under pressure- and capillary-driven flow, and for a nominally homogeneous sample and a heterogeneous one that contains localised deformation features. Based on these results some clear questions might be posed, such as: (i) are the speed measurements truly quantitative and reliable; (ii) how fast can the flow be and still be captured by this method?

To answer the first of the above questions, a comparison against an analytical so-354 lution can be made for the first test on the intact sample, as, in this case the water fronts 355 are sub-horizontal, which allows the system to be considered as 1-dimensional. On this 356 basis, Darcy's law defines a linear relationship q = -Ki, where q is the specific discharge, 357 K is the hydraulic conductivity, i is the hydraulic gradient, with $i = (h_1 - h_2)/z$ where 358 z is taken as the height of the wetted area of the experiment, and h_1 and h_2 are the hy-359 draulic head values at the bottom and top of that region. The pressure (head) at the bot-360 tom of the sample is constant and given by the water-tank level (z_w) . Setting z = 0361 at the bottom of the sample, the hydraulic head is $h_1 = z_w = 160$ mm. The pressure 362 at the water-front is zero and the hydraulic head is $h_2 = z$. The hydraulic gradient changes 363 with time because the height of the wetted region increases and the specific discharge, 364 which is equivalent to a flow speed, is given by: 365

$$q(z) = -K(1 - zw/z).$$
 (1)

Given the flow speed determined for each height in the sample, a value for the pa-366 rameter K can be determined, from which the permeability of the sample, k, can be de-367 rived by $k = K\mu(\rho q)$, where μ and ρ are the water viscosity and density, respectively. 368 The analytical equation (1) has been fitted to the experimental points obtained by av-369 eraging the water speed values over square horizontal slices in the middle of the sam-370 ple, avoiding the edge effects, see Figure 4(c). Fitting of (1) to these data provides a best-371 fit value for K of $3.3*10^{-6}$ m/s, which corresponds to an apparent permeability of 3.3*372 10^{-9} cm², which compares well with lab-measured values of 1–2 Darcy obtained with 373 other samples of this sandstone. The good fitting of the curve and the realistic value found 374 for the permeability suggests that the hypothesis of pressure driven flow is correct and 375 that the results from this imaging method can be used for quantitative analysis. 376

With respect to the second question above, accepted wisdom suggests that, in gen-377 eral, changes in a sample being imaged for tomographic reconstruction should not ex-378 ceed one pixel during the scan to allow a good reconstruction of the volume. In this case, 379 the analysis provided reasonable results even where the water-front was moving around 380 1 mm per minute (i.e., 10 pixels advance of the front during the time taken for the to-381 mography data to be acquired). The movement of the water, however, causes a distor-382 tion in the reconstructed volume since the front is lower in the first projection than in 383 the last one. For this reason, measurements of the frontal advance speed are best made 384 between even numbers of tomographies (i.e., with the same rotation direction) to have 385 the same kind of distortion in all the fronts. This issue can be overcome, in part, by analysing 386 the data as a time-series of projections and not as reconstructed tomographic volumes, 387 as has, for example, been described by Jailin et al. (2018). 388

In other tests not shown in this paper, it appears that, during imbibition, the wa-389 ter is always faster inside the deformation bands regardless of the confining pressure ap-390 plied during the prior triaxial loading (which might be considered to control the char-391 acteristics of the resulting deformation). It is, however, crucial to carry out pressure-controlled 392 flow tests with pre-saturated samples to study the effect of deformation on the local in-393 trinsic permeability. Such tests have already been performed, but with only radiographic 394 (2D) imaging of the sample during flow by Tudisco, Hall, et al. (2017), using a D_2O (heavy 395 water) saturated sample with pressure-driven H_2O to replace the D_2O . As D_2O and H_2O 396 have very similar physical properties, but quite different neutron interactions (D_2O at-397 tenuates neutrons much less than H_2O), this can be considered as a single phase flow sys-398 tem where, as demonstrated by Tudisco, Hall, et al. (2017), it is possible to distinguish 399 an advancing front (between H_2O and D_2O). Moreover, the tests of Tudisco, Hall, et al. 400 (2017) were performed in-situ (i.e., during loading), placing a triaxial cell in the beam-401 line, which enabled the study of the effects of evolving deformation and without unload-402

ing or desaturation of the sample for the imaging. The results in this current work have
shown that it is now possible to perform the 3D imaging sufficiently fast and to control
the flow to be able to follow the flow processes in full 3D.

The fluid flow conditions used for these experiments, although not ideal to study 406 intrinsic permeability, provide opportunities to make observations leading to important 407 insights. An example is the situation depicted in Figure 6, where the water has prefer-408 entially moved upwards, above the general level of water and beyond the position of the 409 notch, along the inclined shear band. The water front is still advancing along the band, 410 411 while the water located within the band, but behind the front, is also continuing to move at high speed. The rock texture inside the band (based on examination of other Vosges 412 samples deformed in the same campaign, and destructively sampled for thin sections and 413 SEM studies) is believed to consist of a zone of broken grains, with smaller pore spaces 414 between the fragments as compared against the texture of the intact rock. That tex-415 ture is expected to result in high capillary pressure values within the band. The fast move-416 ment of water, as observed, is consistent with the inferred zone of higher capillary pres-417 sure. But why does the water behind the front continue to move at a rapid rate? Cap-418 illary pressure is an extremely local phenomenon, and is rooted at the pore scale, where 419 the curvature of the water/air interface (in this example) provides the fluid energy that 420 becomes expressed as an equivalent (negative or suction) pressure. Is it sensible to in-421 fer that the frontal advance pulls water along behind it? The answer is no, of course, and 422 the physics of water movement in the already partly-saturated band has to be examined 423 for another explanation. 424

It is possible to erect a working hypothesis that remains focused on capillary forces. 425 The quasi-planar (but finite-thickness) band may well have a distribution of textures across 426 its area. Some local spots could have smaller pore sizes, and thus higher effective cap-427 illary pressures. These would locally imbibe any water available from nearby regions, seek-428 ing to achieve a higher water saturation that would reduce the local energy differentials. 429 So, there are plausible mechanisms for flow within the band. But the real question is not 430 answered, because if local regions imbibe water from adjacent local regions, those source 431 regions experience de-saturation, which introduces new energy differentials that counter 432 the imbibition flow. Thus, the observed continuity of water in the partly-saturated band, 433 and the persistence of rapid flow within it, implies that the band becomes a favored flow 434 route. This reasoning leads to the hypothesis that the band also possesses a higher-permeability, 435 as well as its higher capillary pressure. Higher permeability is usually thought to anti-436 correlate with higher capillary pressure, but the detailed flow observations provide a ba-437 sis for supposing that the two characteristics exhibit a positive correlation. The same implications arise from consideration of the flow behaviour of the complexly fractured 439 sample depicted in Figure 7, which will be treated in a subsequent paper focused on that 440 sample suite. 441

The value of time-resolved 3D images of fluid flow, in rock samples with internal 442 heterogeneities, is clear. The images provide evidence of flow processes that reveal com-443 plexity in the way that invading fluids move into the rocks, and the way that fluid moves 444 in a saturated sample. In the case of the limestone sample, the experiment involved in-445 troduction of a contrasting fluid (H_2O replacing D_2O), and the resulting flow pattern 446 shows that the later fluid movement is focused in the planar deformation features that 447 were previously thought to be simple opening-mode fractures. Current work is exam-448 ining the local textures of these planar features, via destructive sampling that has en-449 abled thin sections to be made and studied (e.g., as in Figure 7 (d)). Digital-rock meth-450 ods are being applied to the textural images to derive local flow properties (Jiang et al., 451 2017, 2018). These will be assembled into whole-sample models that should provide a 452 basis for developing a full process explanation of the fluid flow processes that can be ob-453 served via high-speed neutron tomography. 454

455 6 Conclusions

This paper has presented a new method, involving high-speed neutron imaging and image analysis, for 3D monitoring of fluid advance in rock specimens, with unprecedented spatio-temporal resolution. The new analysis method allows the determination of fluidfront advance speed at each voxel position in the 3D volume.

Quantitative full-field determinations of fluid-front speeds have been made in two 460 samples of a sandstone, a nominally-intact specimen and one containing localised defor-461 mation features resulting from laboratory triaxial loading, for flow under pressure- and 462 capillary-driven conditions, respectively. Verification of quantitative speed measurements 463 has been performed by comparison to a 1D analytical model of pressure-driven flow for 464 the intact sample, for which an hypothesis of homogeneity is reasonable. This verifica-465 tion provides confidence in the local measurements of capillary-driven fluid-front advance 466 in the sample containing localised deformation features. The resultant fluid-front speed map for the deformed sample has been compared to full-field strain measurements from 468 DVC analysis of x-ray tomography images acquired before and after the laboratory tri-469 axial loading. This comparison reveals that the localisation of fluid-front speed corre-470 lates well with localisation of deformation and that the fluid-front advances faster in the 471 more deformed regions of the samples (i.e., in the shear-bands). The apparent 'pull' of 472 water in these bands likely relates to increased capillary drive relating to reduced pore-473 sizes, which is consistent with the general compactant nature of the localised shear fea-474 tures. 475

Based on the methods presented herein, it is now possible to perform 3D imaging sufficiently fast and to control the flow to be able to follow flow processes in full 3D and in-situ in a neutron imaging station. The next step is to combine this with in-situ triaxial loading and high spatial resolution time-lapse tomography to enable DVC analysis, which, combined with the 3D flow quantification of this paper, will enable studies of the evolution of both the deformation and the fluid flow, thus enabling research in hydromechanical coupling.

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Data supporting the conclusions can be obtained at https://figshare.com/projects/

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